

CHEMISTRY OF LOW-VALENT PLATINUM DIMERS

Thesis by

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Under certain circumstances there are few hours in life more agreeable than the hour dedicated to the ceremony known as afternoon tea. There are circumstances in which, whether you partake of the tea or not—some people of course never do,—the situation is in itself delightful.

Henry James

The Portrait of a Lady

ACKNOWLEDGMENTS

It's hard to believe that I've been at Caltech for five years! I won't say that "I've had the time of my life" here, but there are some people who have made this place pleasant and deserve some thanks.

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And lastly, I want to thank God for all that He has done for me in the past, at the present, and in the future.

*to my Mother, my Father, Ashes
and Tom*

ABSTRACT

Physical and chemical properties of low-valent platinum dimers, namely $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$ and $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$, have been investigated using a variety of structural and spectroscopic techniques.

Platinum(II) $d^8\text{-}d^8$ dimers have been shown to exhibit much thermal and photochemical reactivity. Chapter 2 describes studies aimed at elucidating the excited state reduction potential of $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$, Pt_2 , in organic media. By conducting excited state electron transfer studies using derivatized pyridiniums and benzophenones, the Pt_2 excited state reduction potential has been estimated to be -2 V. The Pt_2 complex undergoes partial oxidation to form Pt(II,III) linear chains. Chapter 3 describes the structural and spectroscopic techniques used to determine the translational symmetries of these $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{X}]^{4-}$ ($\text{X} = \text{Cl}, \text{Br}$), Pt_2X , chains. Pt_2Br has been found to be intermediate between $(\text{AAB})_n$ and $(\text{AABCCB})_n$, while, Pt_2Cl is of $(\text{AABCCB})_n$ translational symmetry. Investigations into the electronic transitions of Pt_2Cl and Pt_2Br were conducted using high pressure techniques and are presented in Chapter 4. The Pt_2X electronic spectrum exhibits bands attributable to the reduced Pt_2 complex and the oxidized Pt_2X_2 complex ($[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{X}_2]^{4-}$) along with an intervalence charge-transfer band characteristic of a mixed-valence solid.

Photophysical investigations of a new luminescent chromophore, $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$, a $d^9\text{-}d^9$ dimer, and its analogs are described in Chapter 5. The absorption band directly responsible for the observed emission is believed to be very weak and, as of yet, unobserved. Attempts to determine the spin multiplicity and approximate energy of this unobserved transition are described in Chapter 6. Excited-state

energy transfer studies indicate that this absorption band is a triplet transition at $\sim 13,000 \text{ cm}^{-1}$. Although, the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ excited state is non-luminescent in fluid solution, it has been shown to undergo thermal electron transfer to tetracyanoethylene and photo-induced electron transfer to methylviologen. These experiments are presented in Chapter 7. Preliminary studies, described in Chapter 8, of non-bridged d^9 - d^9 platinum(I) dimers have shown that $[\text{Pt}_2(\text{CNCH}_3)_6]^{2+}$ serves as a versatile precursor in the synthesis of new d^8 - d^8 A-frame complexes.

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CHAPTER 1

Introduction

INTRODUCTION

A photo-excited molecule is best envisioned as a new species possessing physical and chemical properties distinct from its corresponding ground state molecule. Photophysical studies of various molecular chromophores (namely monomers,^{1,2} dimers,^{3,4} and clusters^{5,6}) have resulted in elaborate and detailed descriptions of the physical properties of their excited states. During these studies, certain chromophores have been recognized as possessing great photochemical potential in the areas of solar energy conversion and small molecule activation.^{1,7} One such molecular chromophore is the d^8-d^8 dimer.

First synthesized by Roundhill and coworkers,⁸ $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$, Pt_2 , is a particularly interesting member of this class of molecules. Its structure can be described as two face-to-face square planar monomeric fragments, which are hooked together by four bridging ligands (see Figure 1.1)⁹; hence, the "idealized" point group is D_{4h} . The molecular orbital diagrams for all d^8-d^8 species have the same basic features, which are illustrated in Figure 1.2. Since all the electrons are spin paired, the ground state is a singlet. The lowest energy excited state is obtained when an electron is promoted from the $d\sigma^*$ orbital to the $p\sigma$ orbital; for Pt_2 this transition is $A_{1g} \rightarrow A_{2u}$.^{10,11} From group theory, the $(d\sigma^*)^2 \rightarrow (d\sigma^*)(p\sigma)$ transition is fully allowed along the z direction,^{10,11} thus the extinction coefficient for the ${}^1A_{1g} \rightarrow {}^1A_{2u}$ transition is expected to be $10^3 - 10^4 \text{ M}^{-1} \text{ cm}^{-1}$; the actual value being $3.45 \times 10^4 \text{ M}^{-1} \text{ cm}^{-1}$.^{10,11} The ${}^1A_{1g} \rightarrow {}^3A_{2u}$ transition is formally spin forbidden; however due to spin-orbit coupling it is observed in the absorption spectrum with an $\epsilon = 110 \text{ M}^{-1} \text{ cm}^{-1}$.^{10,11}

While these features are common to all d^8-d^8 dimers, the exact

Figure 1.1 Structure of $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$.

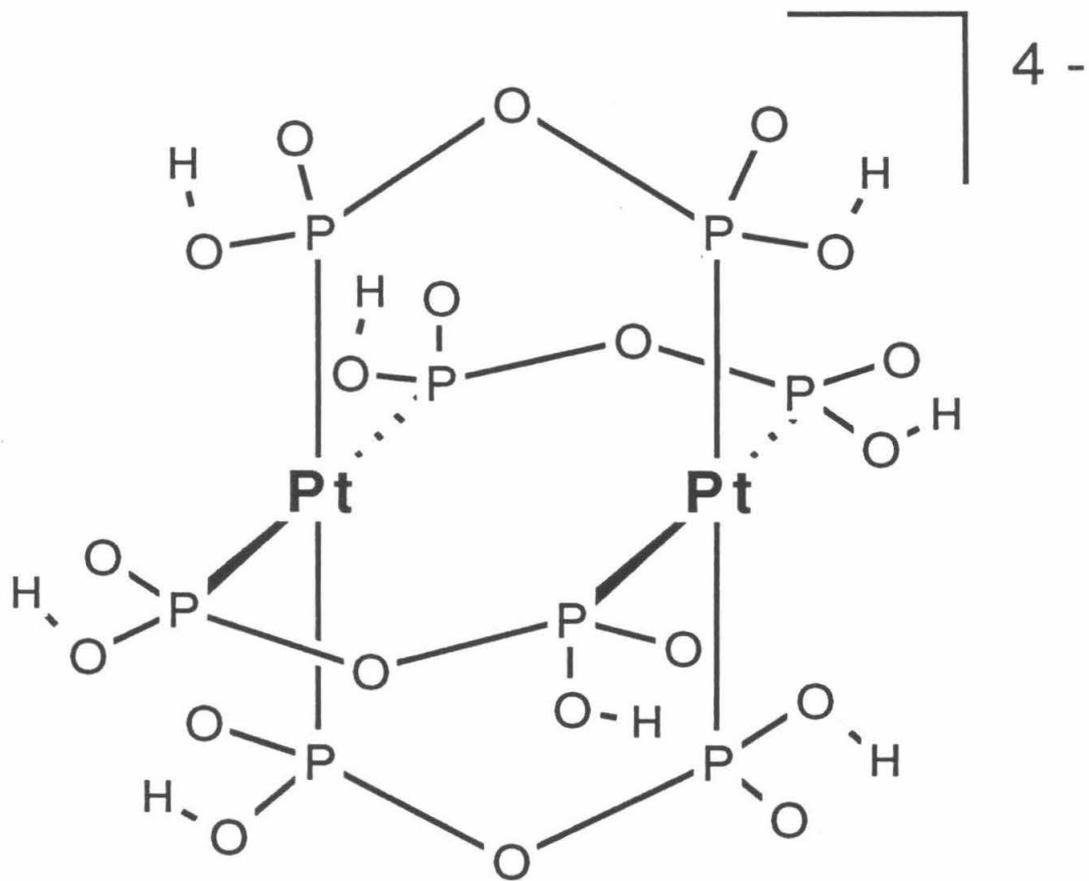
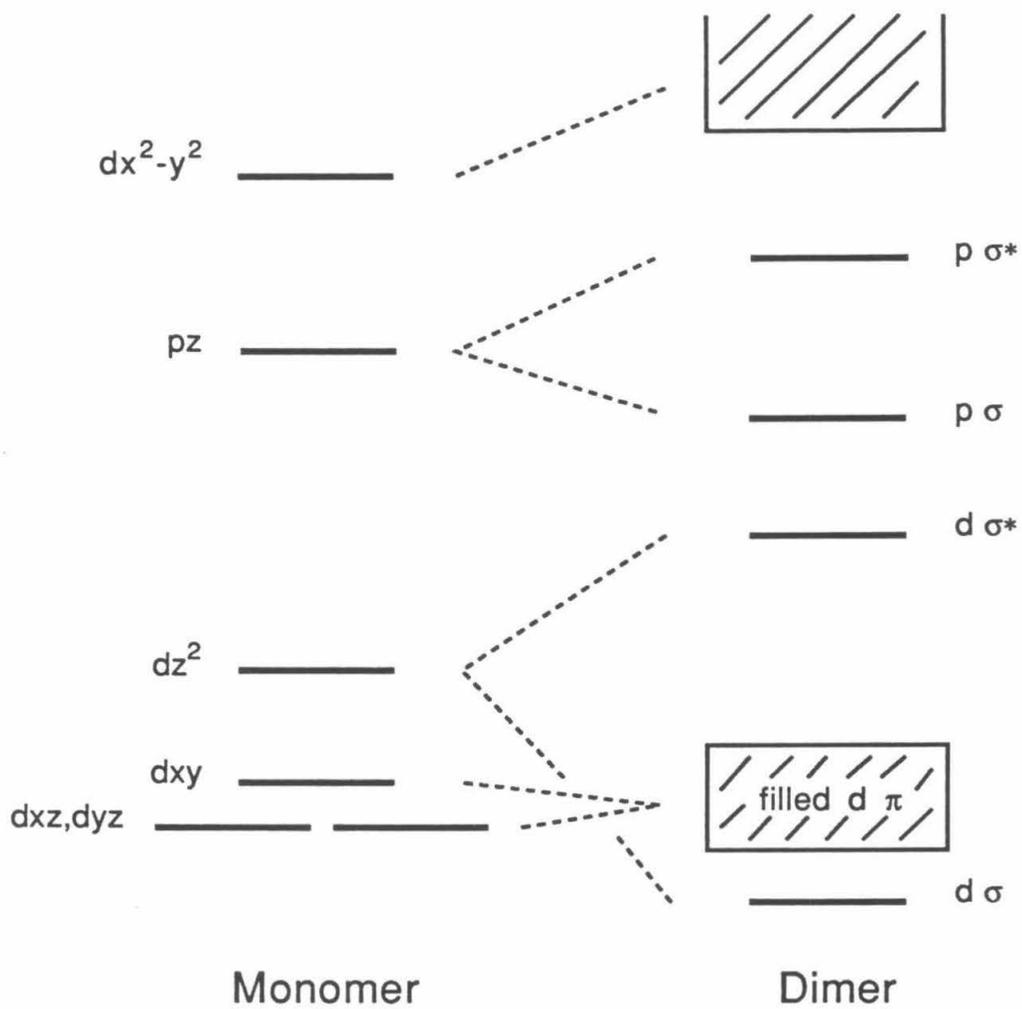


Figure 1.2

Molecular orbital diagram for a d^8-d^8 dimer built from a d^8 monomer.

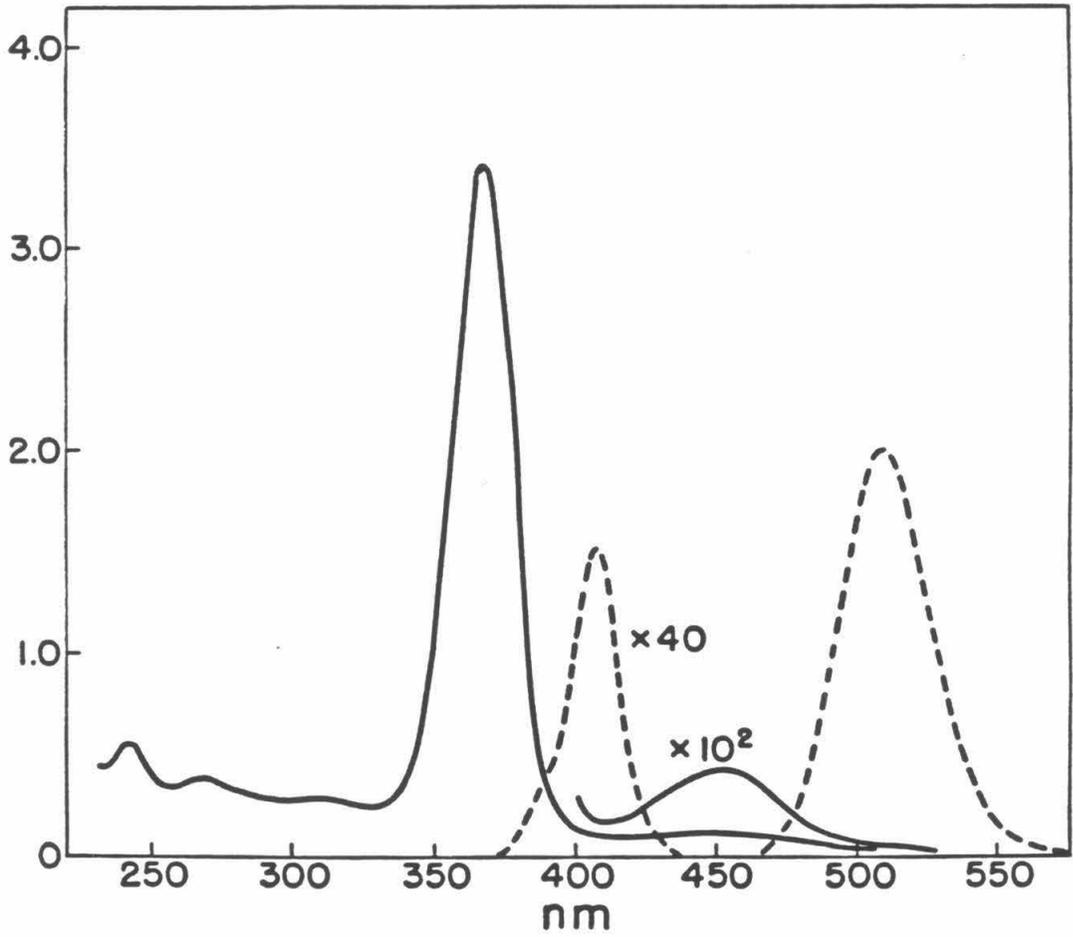


energies of the singlet-singlet and singlet-triplet transitions are a function of the metal and type of ligands, as well as constraints on the metal-metal distance due to the size of the bridging ligand. These factors cause the singlet-singlet and singlet-triplet transitions in Pt_2 to lie at relatively high energies; 27100 cm^{-1} (369 nm) and 22200 cm^{-1} (451 nm) respectively (see Figure 1.3). The molecular orbital diagram predicts that $\text{Pt}_2(\text{pop})_4$ contains no formal metal-metal bonds; however, configurational mixing of orbitals results in a ground state metal-metal interaction. The excited state is predicted to have a formal metal-metal single bond; evidence for this is observed in low temperature absorption spectroscopy, where the vibrational transitions are well resolved,¹⁰ and in time-resolved resonance Raman spectroscopy.^{12,13} The formation of a metal-metal single bond in the excited state has important implications in the photochemistry arising from this state.

The most appealing feature of all d^8-d^8 dimers is their observed luminescence. Although all d^8-d^8 dimers studied thus far luminesce to some extent, emission observed for Pt_2 is truly remarkable, with a quantum yield of emission greater than 0.50 (see Figure 1.3).¹⁴ The majority of the emission results from radiative decay from the triplet excited state (phosphorescence) with a luminescence lifetime of 10 μs at room temperature. Emission from the singlet excited state (fluorescence) is also observed; its luminescence lifetime is estimated to be 40 ps.^{10,11} Due to the shortness of the fluorescent lifetime, it is unlikely that any bimolecular photochemistry will occur directly from the singlet excited state.

These properties of Pt_2 can be used to investigate the correlation

Figure 1.3 Electronic absorption (—) and corrected emission (----) spectra (taken from reference 10b) of an aqueous solution (300 K) of $\text{K}_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]$; (left scale, $\epsilon \times 10^{-4} \text{ M}^{-1} \text{ cm}^{-1}$ values).



of one electron potentials with photoreactivity involving one or more electrons. Chapter 2 describes quenching experiments designed to investigate the photoreactivity of the triplet excited state in organic media.

The Pt_2 dimer crystallizes in the tetragonal space group $P4/mbm$ as infinite chains along the c axis when K^+ is present as the counterion.⁹ Oxidation of the Pt_2 structure by the introduction of different groups, which act as bridges between the dimers, creates d^7-d^8 systems which exhibit interesting conduction properties.^{15,16} The results of X-ray diffraction studies and Raman experiments, which offer a clearer structural view of these linear chain systems, are discussed in Chapter 3.

In addition to bond length changes in the $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4$ core resulting from oxidation by halogen, the molecular orbital description of the $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{X}]^{4-}$ chains changes. Chapter 4 investigates the effects of pressure on the electronic absorption spectra for $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)\text{X}_n]^{4-}$ ($n = 0,1,2$); these results have increased our insight into the structural and electronic factors which dominate in the molecular orbital description of these systems.¹⁷

In addition to being desirable candidates for photophysical study, d^8-d^8 dimers participate in many interesting photochemical reactions. Two types of reaction classes, established thus far, which result in net photochemistry are (i) bimolecular oxidative-addition reactions initiated by electron transfer^{18,19} and (ii) hydrogen atom abstraction reactions.²⁰⁻²⁴ In both of these reactions, steric restrictions imposed by the bridging ligands of the d^8-d^8 starting material force substrate addition to occur in the axial position, resulting in the formation of d^7-d^7 dimers as products. Because of this steric re-

striction, it is difficult to envision photochemical, concerted addition of small molecules such as CO_2 , CS_2 , and SO_2 . If the metal-metal core were made more assessible to substrates such that adjacent coordination sites were available, concerted small molecule addition may be possible.

It is with this thought in mind that the photophysical and photochemical study of the nearly planar d^9-d^9 platinum dimer $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ (dppm = bis(diphenylphosphino)methane) and its analogs was undertaken (see Figure 1.4 and Table 1.1).^{25,26} Both platinum and palladium d^9-d^9 complexes, $\text{M}_2(\mu\text{-dppm})_2\text{Cl}_2$ ($\text{M} = \text{Pt}, \text{Pd}$), have been shown to thermally add small molecules across the metal-metal bond to form the corresponding d^8-d^8 A-frame complexes (see Figure 1.5).²⁷⁻³³ The palladium dimers failed to exhibit any detectable luminescence; therefore the photochemical study of these complexes centered about the platinum dimers. Investigation into their photophysical properties is described in Chapter 5. More detailed information about the d^9-d^9 excited state reactivity is presented in Chapters 6 and 7 where energy transfer and electron transfer studies on these complexes are described. And in closing, Chapter 8 briefly discusses preliminary investigations on non-bridging d^9-d^9 dimers. Unlike the nearly planar dppm bridged complexes, these molecules have a staggered geometry.³⁴ Their thermal and photochemical reactivity is dominated by homolytic metal-metal bond cleavage reactions similar to those observed for the d^7-d^7 $\text{M}_2(\text{CO})_{10}$ ($\text{M} = \text{Mn}, \text{Re}$) complexes.³⁵⁻³⁷

Figure 1.4 Structure and ORTEP drawing (taken from reference 26) of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$.

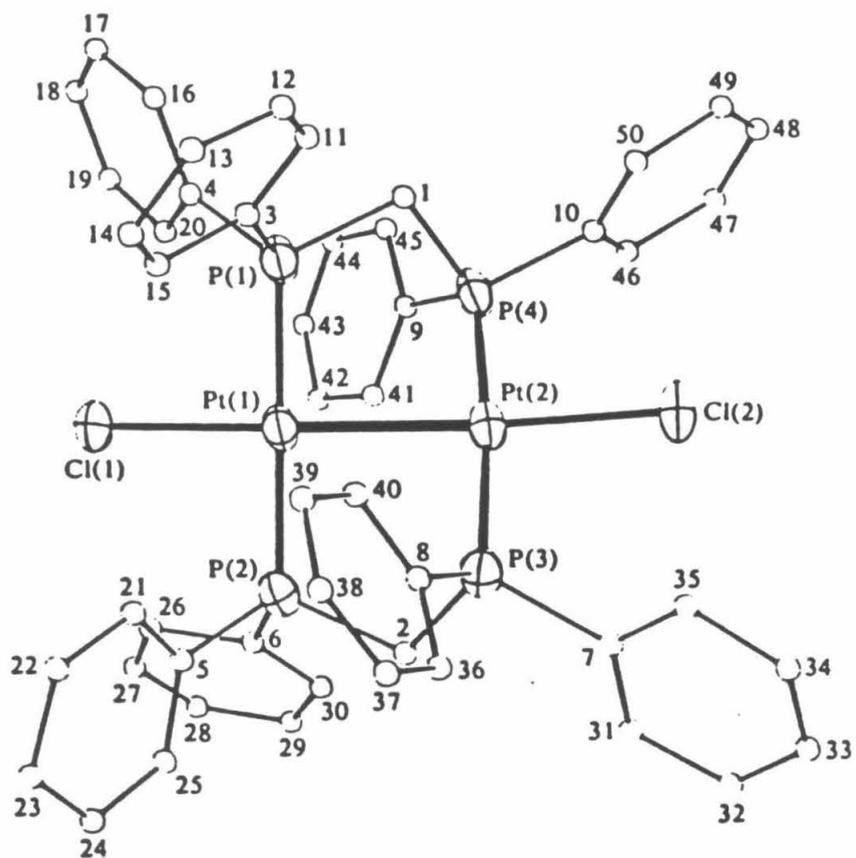
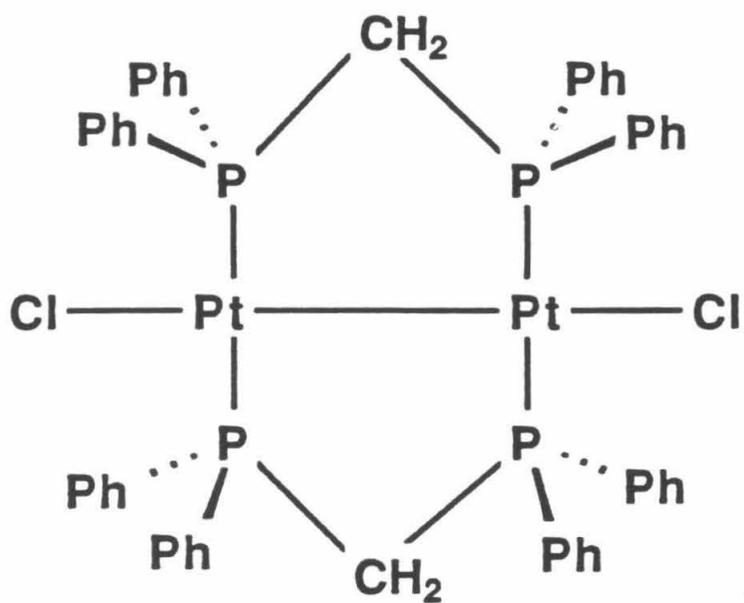


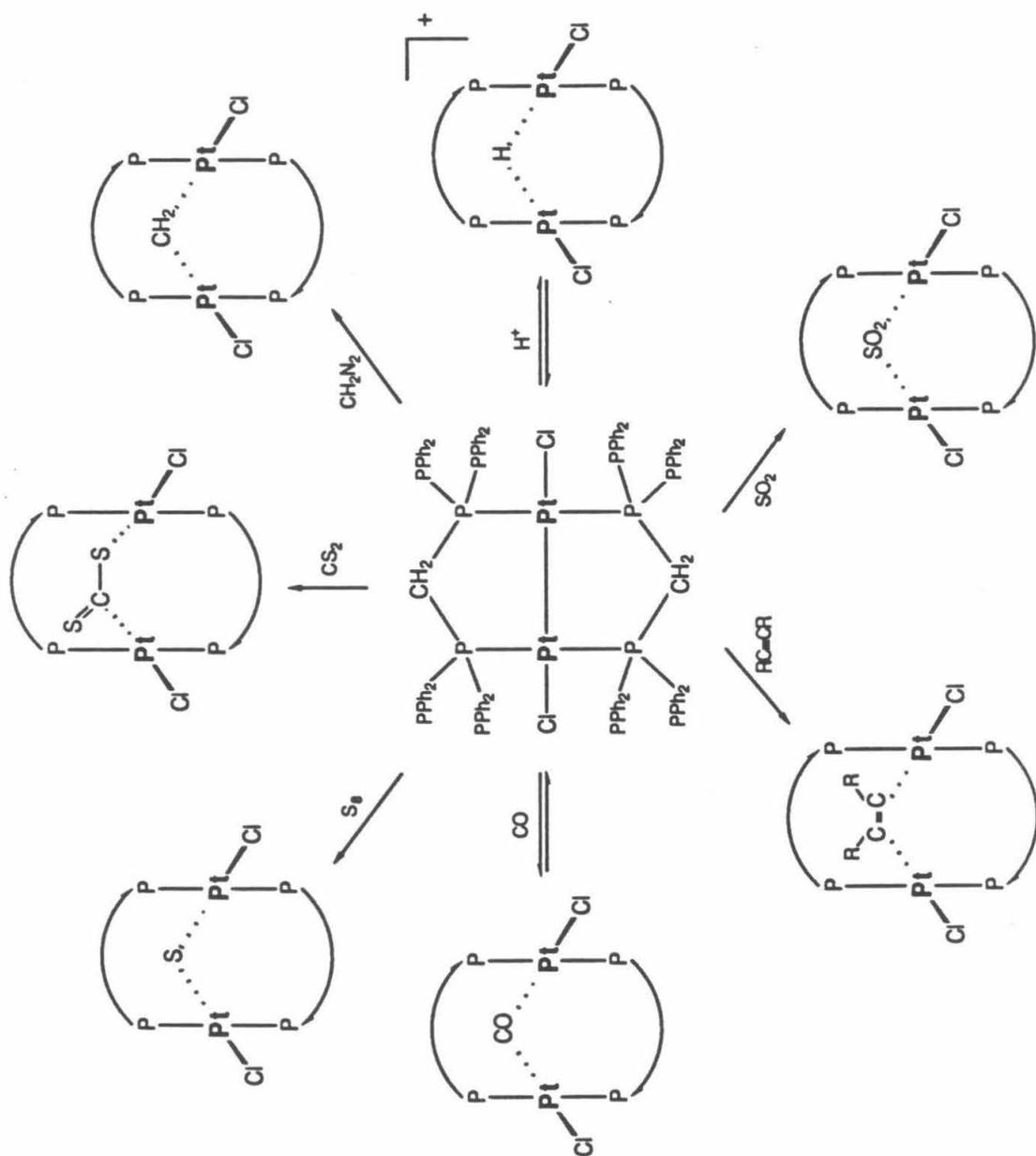
Table 1.1

Selected Bond Lengths and Bond Angles of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2^\dagger$

<u>Bond Lengths (Å)</u>		<u>Bond Angles (°)</u>	
Pt(1) - Pt(2)	2.651(1)	Cl(1) - Pt(1) - Pt(2)	178.5(2)
Pt(1) - Cl(1)	2.408(5)	P(1) - Pt(1) - P(2)	175.6(2)
Pt(2) - Cl(2)	2.401(5)	Cl(1) - Pt(1) - P(1)	90.2(2)
Pt(1) - P(1)	2.294(7)	Cl(1) - Pt(1) - P(2)	89.6(2)
Pt(1) - P(2)	2.264(7)	Pt(2) - Pt(1) - P(1)	89.3(2)
Pt(2) - P(3)	2.259(7)	Pt(2) - Pt(1) - P(2)	91.0(2)
Pt(2) - P(4)	2.250(7)		
		Cl(2) - Pt(2) - Pt(1)	175.7(2)
		P(3) - Pt(2) - P(4)	172.9(2)
		Cl(2) - Pt(2) - P(3)	94.3(2)
		Cl(2) - Pt(2) - P(4)	92.2(2)
		Pt(1) - Pt(2) - P(3)	88.7(2)
		Pt(1) - Pt(2) - P(4)	85.0(2)

[†]Data taken from reference 26.

Figure 1.5 Known thermal chemistry of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$.



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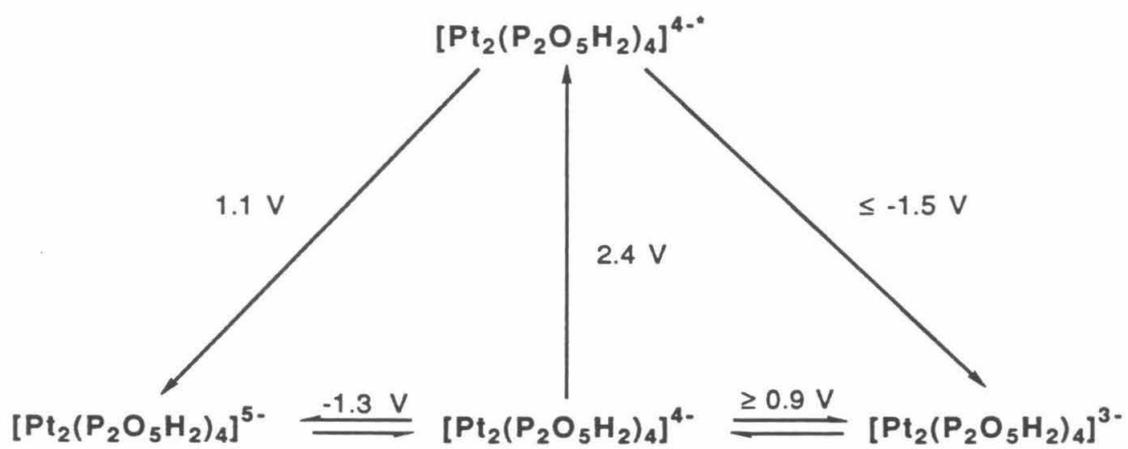
CHAPTER 2**Electron Transfer Studies of $(\text{TBA})_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]$**

INTRODUCTION

Several experiments were performed which attempted to establish the ground and excited state redox properties of $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$, Pt_2 , in water.¹⁻⁴ Cyclic voltammograms were measured; however, the redox couples of Pt_2 were irreversible, suggesting that the oxidized and reduced species of Pt_2 are very short lived. Hence only estimates of the oxidation and reduction potentials were obtained. Fairly good approximations for the excited state redox values were obtained through oxidative and reductive quenching studies.^{1,3-5} From these investigations it was determined that Pt_2 was both a strong reducing agent^{3,4} and a strong oxidant in water.¹ The modified Latimer diagram for Pt_2 in water is illustrated in Figure 2.1.

Since the potassium salt of Pt_2 is only soluble in water, attempts to find routes to new photochemical reactions were hampered by the limited solvent selection. Metathesis of the potassium salt to tetrabutylammonium, TBA^+ , yielded a salt soluble in a variety of solvents. Along with the increase in solvent selection, the reactivity of the dimer appeared to increase in non-aqueous media, which led to new questions concerning the reducing power of the dimer. This chapter describes electron transfer quenching experiments designed to probe the reducing power of Pt_2 in acetonitrile.

Figure 2.1 Modified Latimer diagram for $\text{K}_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]$ (E in V vs. SCE) in water.



EXPERIMENTAL

Materials

All solvents, unless otherwise noted, were reagent grade and used as received without further purification.

$K_4[Pt_2(P_2O_5H_2)_4]$ and $(TBA)_4[Pt_2(P_2O_5H_2)_4]$: The potassium salt of Pt_2 was prepared by the published procedure.⁶ Metathesis of the potassium salt to the tetrabutylammonium salt was carried out in the dark following, for the most part, the published procedure.⁷ For improved results: $K_4[Pt_2(P_2O_5H_2)_4]$ (0.25 g) was dissolved in 10 mL of water. A 4:1 molar excess of tetrabutylammonium chloride (TBACl) in 30 mL of CH_2Cl_2 was added to the Pt_2 solution with vigorous shaking. The platinum complex was extracted into the organic layer (yellow) and separated. The $(TBA)_4[Pt_2(P_2O_5H_2)_4]$ was precipitated from CH_2Cl_2 by adding an excess of dry diethyl ether (~ 100 mL).

Pyridinium hexafluorophosphates. The pyridinium hexafluorophosphate quenchers were prepared and recrystallized by Dr. Janet Marshall and used as received.⁸

Benzophenone and its derivatives. Benzophenone (Aldrich, 99%), 4-methylbenzophenone (Aldrich, 97%), 4-fluorobenzophenone (Aldrich, 97%), 4-methoxybenzophenone (Aldrich, 97%), 4-dimethoxyaminobenzophenone (Aldrich, 98%) were recrystallized from hot 2:1:1 (v/v/v) water:acetone:ether and used without further purification.

Methylviologen hexafluorophosphate. Methylviologen dichloride hydrate (Aldrich) was dissolved in a minimum amount of water to which a saturated solution of potassium hexafluorophosphate, KPF_6 , (Aldrich, 98%) was added. Methylviologen hexafluorophosphate precipitated immediately, was filtered, washed with water, and recrystallized from hot

2:1:1 (v/v/v) water:acetone:ether yielding long white needles.

Tetrabutylammonium hexafluorophosphate, TBAPF₆. Equal molar concentrations of saturated tetrabutylammonium bromide, TBABr, (Aldrich, 98%) in water and saturated KPF₆ in water were added together. TBAPF₆ precipitated, was filtered, and recrystallized from hot absolute ethanol yielding white needles.

Physical Measurements

Electrochemical Measurements. All cyclic voltammograms were measured using EG&G PAR electronics, Models 175 (Universal Programmer), 173 (Potentiostat/Galvanostat) and 179 (Digital Coulometer) with a Houston Instruments 2000 X-Y Recorder. The reference electrode was a standard calomel electrode, SCE; both the working and auxiliary electrodes were platinum wires. The cyclic voltammograms were measured in acetonitrile (Burdick and Jackson, UV grade) solutions, purged with N₂, with 0.1 M TBAPF₆ as the supporting electrolyte. Measurements were made at both 200 and 500 mV sec⁻¹ scan rates.

Stern-Volmer Quenching Procedures. Rate constants obtained from luminescence lifetime quenching data were measured using a Nd:YAG pulsed laser system described previously.⁹ Rate constants obtained from emission intensity quenching data were measured using an emission spectrophotometer constructed at Caltech and described elsewhere.¹⁰ All emission spectra were recorded and corrected for monochromator and photomultiplier tube response using the OLIS fluorimetry programs.¹¹

Samples for all electron transfer measurements were prepared by placing Pt₂ (~1 mg) was placed in a two-compartment spectrophotometric cell (referred to as a quenching cell) equipped with two teflon vacuum

valves and consisting of a square cuvette (1 cm pathlength) and a round bottom flask (pyrex, ~15 mls.). Exactly 4 mL of acetonitrile (Burdick and Jackson, UV grade) was added to the round bottom flask of the quenching cell and the solution was freeze/pump/thawed a minimum of 6 times on a high vacuum line (limiting pressure $< 10^{-3}$ Torr). Quencher solutions were added to the square compartment in aliquots of known concentration; the quencher solvent was distilled off under vacuum to leave behind known quantities of the solid quencher. The quenching cell was re-evacuated while the already evacuated solution was sealed in the round bottom compartment to maintain the oxygen-free conditions of the experiment.

RESULTS and DISCUSSION

As in the case of the cyclic voltammograms measured for Pt₂ in water, those measured in acetonitrile showed irreversible redox couples.

Initial electron transfer quenching experiments were performed using a series of pyridinium quenchers with reduction potentials of -0.67 V to -1.85 V vs. SCE (see Table 2.1).⁸ (Energy transfer quenching can be ruled out since the pyridiniums' triplet energy is too high to make it a competitive process with electron transfer.) Data were collected using luminescent lifetime techniques (excitation wavelength 355 nm and emission wavelength 515 nm) and the electron transfer rates were calculated by fitting the lifetime and concentration values to the Stern-Volmer equation:

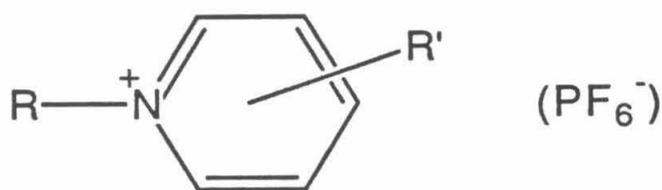
$$\tau_0/\tau = 1 + k_q[Q]\tau_0 .$$

(τ_0 is the phosphorescent excited state lifetime in the absence of quencher, τ is the phosphorescent excited state lifetime at various quencher concentrations [Q], and k_q is the quenching rate constant.) An example of Stern-Volmer quenching is illustrated in Figure 2.2.

The value of k_q for all pyridinium quenchers used was on the order of $10^9 - 10^{10} \text{ M}^{-1} \text{ sec}^{-1}$, which is the upper limit for diffusion controlled reactions in acetonitrile. Even 2,6-dimethyl-4-methoxy-N-methyl-pyridinium, with a reduction potential of -1.85 V, had a quenching rate constant of $1.5 \times 10^9 \text{ M}^{-1} \text{ sec}^{-1}$, which is remarkably high (Table 2.1). The high quenching rate constants and the relatively insignificant drop off with increasing negative reduction potentials which were observed can be explained by electrostatic interactions. In

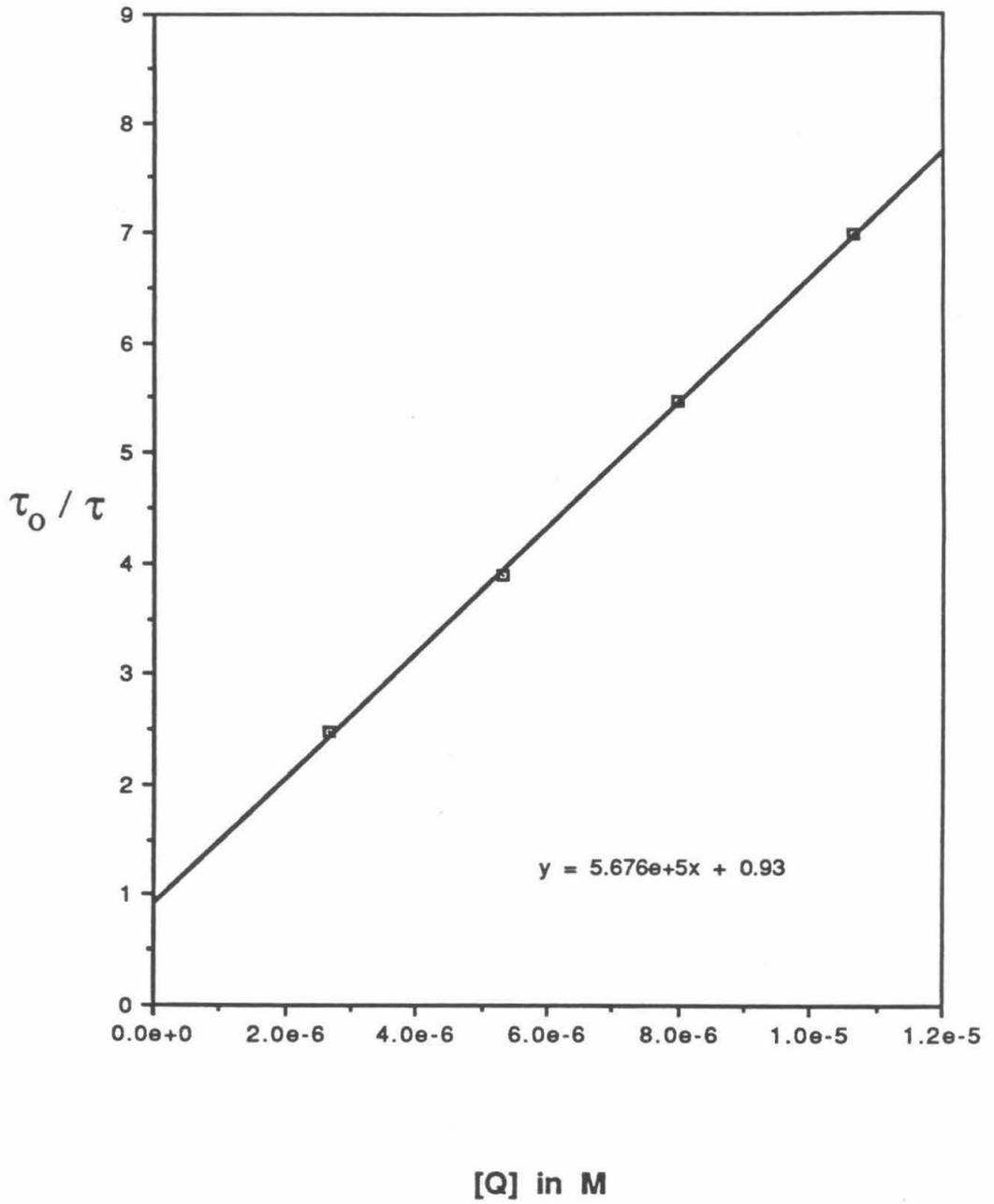
Table 2.1

Electron Transfer Data Using Pyridiniums



Quencher	R	R'	E (A/A ⁻) V vs. SCE	k _q
4-Cyano- N-methylpyridinium	CH ₃	4-CN	-0.67	2.0 × 10 ¹⁰
4-Carbomethoxy- N-methylpyridinium	CH ₃	4-CO ₂ CH ₃	-0.78	3.2 × 10 ¹⁰
4-Amido- N-methylpyridinium	CH ₃	3-CONH ₂	-1.14	1.4 × 10 ¹⁰
4-Ethylpyridinium	CH ₃		-1.36	4.1 × 10 ¹⁰
2-Methoxy- N-methylpyridinium	CH ₃	2-OCH ₃	-1.48	6.7 × 10 ¹⁰
4-Methoxy- N-methylpyridinium	CH ₃	4-OCH ₃	-1.49	5.9 × 10 ¹⁰
2,6-Dimethyl-4-methoxy- N-methylpyridinium	CH ₃	2,6-(CH ₃) ₂ 4-OCH ₃	-1.85	3.2 × 10 ¹⁰

Figure 2.2 Stern-Volmer plot of the electron transfer quenching of $(\text{TBA})_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]$ with 2-methoxy-N-methylpyridinium ($E_{1/2} \sim -1.48$ V vs. SCE).



solution, the positively charged pyridinium quenchers are attracted to the high negative charge (4-) of Pt_2 ; therefore, they congregate about the dimer, setting up a non-random, non-uniformly distributed, quencher-dimer localized system. When the laser pulse excites the platinum dimer, one electron transfer to a closely associated quencher molecule is easily facilitated, resulting in the observed high quenching rate constants. Along the same lines, attempts to gather quenching rate data using methylviologen, MV^{2+} , as the electron acceptor caused precipitation of an orange solid presumed to be Pt_2 tight-ion paired with the MV^{2+} cation. These results indicated that in order to get a more accurate description of the bimolecular electron transfer rate, neutral or negatively charged electron acceptors are required.

The second series of quenchers used were derivatized benzophenones with reduction potentials ranging from -1.86 V to -2.03 V vs. SCE (see Table 2.2). The cyclic voltammograms of the derivatized benzophenones showed quasi-reversible reduction couples. The peak to peak separation in all samples was significantly greater than 59 mV and the peak heights for the forward and back waves were not equivalent; however, reproducible $E_{1/2}$ values at different scan rates were obtained.

Electron transfer rates between Pt_2 and the various derivatized benzophenones were determined using emission intensity quenching techniques since lifetime data obtained from laser excitation of the phosphorescent excited state exhibited biphasic behavior. The observed biphasic behavior of the lifetime data was ascribed to energy transfer from the excited benzophenone quencher to the Pt_2 dimer. The electronic absorption spectrum of the derivatized benzophenones clearly

showed significant absorption at 355 nm. Thus, the derivatized benzophenones were excited by the laser pulse, then underwent energy transfer to the ground state Pt_2 to produce Pt_2^* . The biphasic lifetime decay reflected the fact that Pt_2 excitation to Pt_2^* was no longer occurring only from the laser pulse, but occurred as a result of the energy transfer mechanism as well.

In using the emission intensity quenching technique, Pt_2 was excited by continuous irradiation at 436 nm (a wavelength at which none of the derivatized benzophenones had any absorption). The decrease in phosphorescent intensity was monitored with respect to increasing quencher concentrations (see Figure 2.3). Once again Stern-Volmer kinetics:

$$\Phi_0/\Phi = 1 + k_q[\text{Q}]\tau_0$$

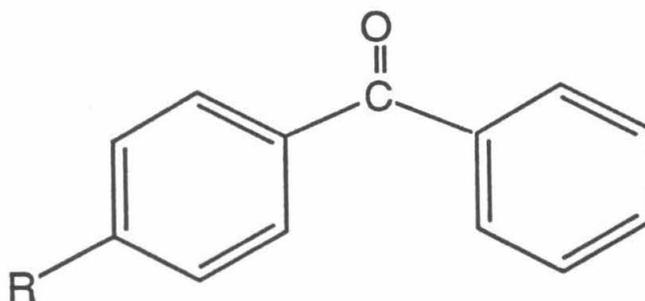
were used to determine the quenching rate constants; see Table 2.2. (Φ_0 is the intensity of the emission band at 516.7 nm in the absence of quencher, and Φ is the intensity of the emission band at 516.7 nm at various quencher concentrations.)

The quenching rate constant obtained for benzophenone, whose reduction potential is quite similar to that of 2,6-dimethyl-4-methoxy-N-methyl-pyridinium, was an order of magnitude less than that obtained for the pyridinium quencher ($1.0 \times 10^8 \text{ M}^{-1} \text{ sec}^{-1}$ for benzophenone vs. $1.5 \times 10^9 \text{ M}^{-1} \text{ sec}^{-1}$ for the pyridinium quencher). This result supports the notion that the strong electrostatic attraction between the positively charged pyridinium quenchers and Pt_2 led to quenching rate constants which were anomalously high.

The quenching rate constant obtained for the quencher with the most

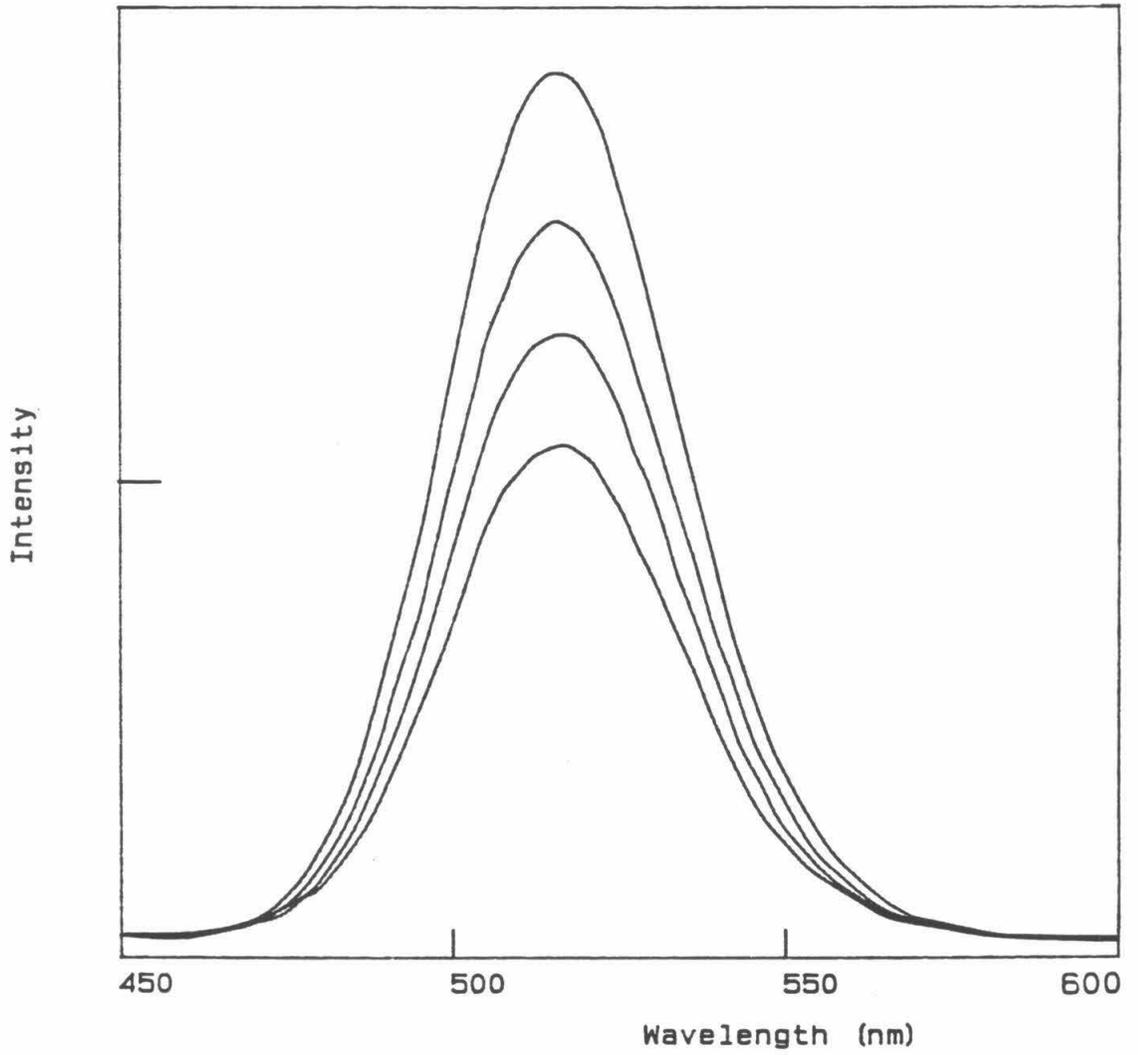
Table 2.2

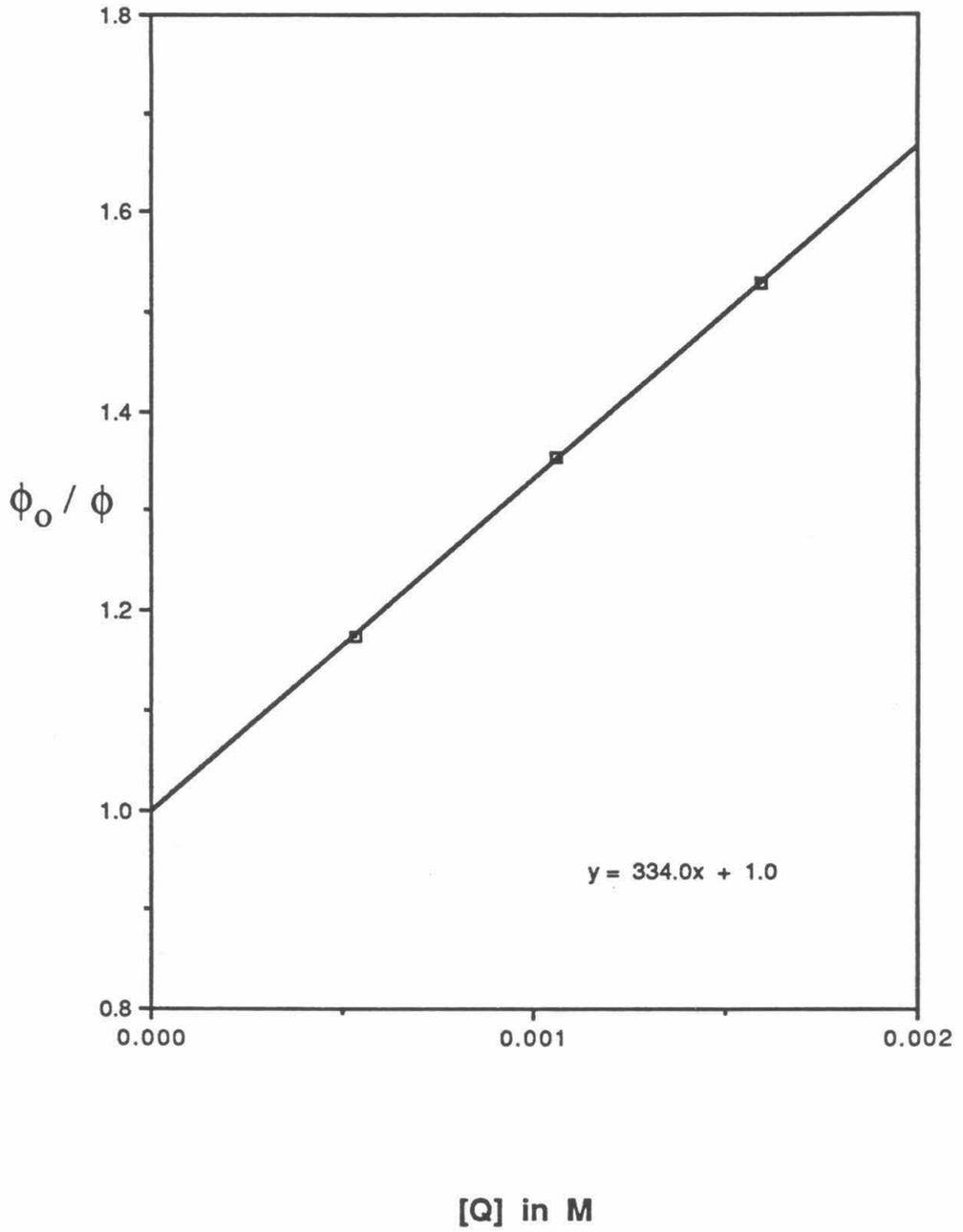
Electron Transfer Data Using Benzophenones



Quencher	R	E (A/A ⁻) V vs. SCE	k _q
Benzophenone	H	-1.86	1.0 × 10 ⁸
4-Methylbenzophenone	CH ₃	-1.89	8.5 × 10 ⁷
4-Fluorobenzophenone	F	-1.89	8.5 × 10 ⁷
4-Methoxybenzophenone	OCH ₃	-1.98	4.8 × 10 ⁷
4-Dimethylaminobenzophenone	N(CH ₃) ₂	-2.03	4.6 × 10 ⁷

- Figure 2.3**
- (a) Emission intensity quenching of $(\text{TBA})_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]$ with 4-dimethylaminobenzophenone (5.3×10^{-4} M aliquot additions of quencher).
 - (b) Stern-Volmer plot of the electron transfer quenching of $(\text{TBA})_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]$ with 4-dimethylaminobenzophenone.



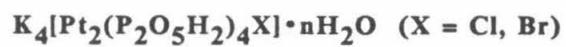


negative reduction potential, 4-dimethylaminobenzophenone, was $4.6 \times 10^7 \text{ M}^{-1} \text{ sec}^{-1}$, suggesting the Pt_2^* has an oxidation potential of about -2.0 V.

A plot of $RT/\ln k_q$ vs. the quencher reduction potential, should yield a slope whose value is 0.5 by Marcus theory for outer-sphere electron transfer;^{12,13} however for the benzophenone data the fit to Marcus theory was poor. One reason for the poor fit may be that the electron transfer is not occurring by a purely outer-sphere mechanism, but may be partially facilitated by an inner-sphere process as well. However, since the reduction potential range of the quenchers is less than 0.2 V, and the change in the quenching rate constant is less than one order of magnitude, the data range is insufficiently broad to allow any definite conclusions to be drawn.

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CHAPTER 3**Translational Symmetries in the Linear-Chain Semiconductors**

INTRODUCTION

Materials having anisotropic electrical conductivity, whether built from aromatic organic molecules or transition metal complexes, have been the subject of much contemporary research.¹⁻⁷ The electronic properties of a linear chain are determined by the extent of interaction between neighboring atoms along the chain and by the translational symmetry of the linear chain. While the importance of the overlap between atomic orbitals of neighboring atoms in the linear chain has been discussed,⁸⁻¹⁰ little attention has been directed to the desired characteristics of the translational symmetry¹¹ required for optimization of a property such as electrical conductivity. Indeed, the mechanisms of electrical transport in these materials are unknown. To date, the variety of translational symmetry types in linear chains has been limited. $K_2Pt(CN)_4Br_{0.3} \cdot 3.2H_2O$ is typical of the near perfect $(A)_n$ symmetry found in most linear chain systems made from mononuclear transition metal complexes.^{5,12} Tetrathiofulvalenium tetracyanoquinodimethanide, TTF-TCNQ, with $(A)_n$ along each of the two segregated stacks of TTF cations and TCNQ anions, is representative of 1-dimensional organic conductors.^{1,4,13} Salts of tetramethyltetraselenafulvalene, such as $[(TMTSF)_2(ReO_4)]$, also have this symmetry, and are members of the first class of linear chain compounds that has been found to undergo a transition to a superconductive phase.¹⁴⁻¹⁶

Our group has been concerned with the bonding in d^8-d^8 transition metal dimers, such as $Rh_2b_4^{2+}$ ($b =$ diisocyanopropane).¹⁸⁻²² The rhodium dimer has interesting photophysical properties resulting from metal-metal interactions. For example, the lowest excited electronic state (${}^3A_{2u}; (d\sigma)^2(d\sigma^*)^1(p\sigma)^1$) possesses a relatively strong Rh-Rh

bond.^{20,21} Interestingly, some of the rhodium isocyanides are 1-dimensional materials.²² In our studies of d^8-d^8 platinum complexes, we discovered a material that contained a linear chain of $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$ units bridged by a halide ion²³⁻²⁶ and having a metallic luster. The 300 K X-ray structure of Pt_2X ($\text{Pt}_2\text{X} = \text{K}_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{X}]$; X = Cl, Br, I) has been reported;²⁷ it revealed an apparent $(\text{AAB})_n$ linear chain system that prompted us to investigate the properties of this unusual system. Some of the possible translational geometries for the three-atom repeat unit of Pt_2X are shown in Figure 3.1. In this chapter, focus is placed on the structures of the 1-dimensional materials, their Raman spectra, and which of these structural schemes are valid for the Pt_2X system.

During the course of this work, the preparation and room temperature X-ray structures of an analogous 1-dimensional system, $\text{Ni}_2(\text{CH}_3\text{CS}_2)_4\text{I}$ and $\text{Pt}_2(\text{CH}_3\text{CS}_2)_4\text{I}$, were reported.²⁸ Also, the 300 K structure of $\text{K}_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}] \cdot 3\text{H}_2\text{O}$ has been independently determined and reported, together with Raman and infrared data for $\text{K}_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{X}] \cdot n\text{H}_2\text{O}$ (X = Cl, Br, I).²⁹ Band theory calculations have been done for the general Pt_2X system.¹¹

Figure 3.1 Translational symmetries for $\text{K}_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{X}]$ (X = Cl, Br).



- Scheme 1: (AAB)_n -P_t^{2.5}—X—P_t^{2.5}—X—P_t^{2.5}—P_t^{2.5}—X—P_t^{2.5}—P_t^{2.5}—X—
- Scheme 2: (AABCCB)_n -P_t²—P_t²—X—P_t³—P_t³—X—P_t²—P_t²—X—
- Scheme 3: (ABC)_n -P_t²—P_t³—X—P_t²—P_t³—X—P_t²—P_t³—X—
- Scheme 4: -P_t^{2.5+δ}—P_t^{2.5+δ}—X—P_t^{2.5-δ}—P_t^{2.5-δ}—X—P_t^{2.5+δ}—P_t^{2.5+δ}—X—
- Scheme 5: -P_t^{2.5+δ}—P_t^{2.5-δ}—X—P_t^{2.5+δ}—P_t^{2.5-δ}—X—P_t^{2.5+δ}—P_t^{2.5-δ}—X—

EXPERIMENTAL

Materials

Preparation of $K_4[Pt_2(P_2O_5H_2)_4Cl] \cdot 3H_2O$, Pt_2Cl : The preparations of both Pt_2Cl and Pt_2Br are similar in that Pt_2 and Pt_2X_2 ($Pt_2 = K_4[Pt_2(P_2O_5H_2)_4]$ and $Pt_2X_2 = K_4[Pt_2(P_2O_5H_2)_4X_2]$; X = Cl, Br, I) are the reactants. However, an improved method for preparing the reduced parent metal dimer, $K_4[Pt_2(P_2O_5H_2)_4] \cdot 2H_2O$, was used.²⁵ $K_4[Pt_2(P_2O_5H_2)_4]$ (0.08 g) and $K_4[Pt_2(P_2O_5H_2)_4Cl_2]$ (0.08 g) were dissolved in a minimum amount of water together with 0.4 g KNO_3 . Absolute methanol was added dropwise until a dark precipitate began to form. A few drops of water were added to clarify the solution. The solution was then cooled slowly to 5°C and allowed to stand for 2 days. The bronze crystals were filtered and dried. [Analysis calculated for $K_4[Pt_2(P_2O_5H_2)_4Cl] \cdot 3H_2O$ (%): Pt, 32.19; P, 20.45; Cl, 2.93. Found: Pt, 32.4; P, 20.6; Cl, 2.82.]

Alternatively, 2 mL of chlorine water were added to a solution of 0.4 g $K_4[Pt_2(P_2O_5H_2)_4]$ in 5 mL of H_2O . A saturated solution of KCl in 5 mL H_2O was added next, followed by several drops of methanol to see if the dark precipitate was observed. Crystals were obtained in the same manner as that stated above.

Preparation of $K_4[Pt_2(P_2O_5H_2)_4Br] \cdot 3H_2O$, Pt_2Br : Pt_2Br was prepared from $K_4[Pt_2(P_2O_5H_2)_4]$ by the published procedure.²⁷

The appearance of Pt_2X under transmitted light shows some sample dependence resulting in subtle differences in the Raman data.³⁰

Physical Measurements

Crystal Structure Determinations. The general features of both

crystal structure determinations were the same. The parameters of Pt_2Cl at 22 K, and Pt_2Cl at 300 K were refined in the tetragonal space group, $P4/mbm$ (#127). A summary of the crystal and intensity collection data for both structures is given in Table 3.1. The crystals were mounted with epoxy on glass fibers. The 300 K diffraction data were taken on a Syntex $P2_1$; low temperature data sets were taken on a locally modified Syntex $P\bar{1}$ diffractometer.³¹ Both diffractometers were equipped with graphite-monochromated $\text{Mo K}\alpha$ X-radiation ($\lambda = 0.71073 \text{ \AA}$). Cell dimensions, given in Table 3.1, were determined from the setting angles of 15 reflections with $25^\circ < 2\Theta < 30^\circ$. The temperatures reported for the Syntex $P\bar{1}$ data sets have error limits of $\pm 1 \text{ K}$.

Data were collected by Θ - 2Θ scans at $2^\circ/\text{minute}$ from $(2\Theta(K\alpha_1)-1)^\circ$ to $2\Theta(K\alpha_2)+1^\circ$. The three check reflections collected after every 97 reflections showed a slight decrease in intensity in all crystals ($\sim 3\%$). Variances, $\sigma^2(I)$, were assigned on the basis of counting statistics plus an additional term, $(0.024 \cdot I)^2$, to account for fluctuations proportional to the diffracted intensity. Intensities were corrected for Lorentz and polarization effects and crystal decay. Using the crystal dimensions and the linear absorption coefficient, the data were corrected for absorption by Gaussian integrations over an $8 \times 8 \times 8$ grid. All structures were refined using full-matrix least-squares techniques, minimizing $\sum w(F_o^2 - F_c^2)^2$, where $w = 1/\sigma^2(F_o^2)$. For both structures, final difference maps were somewhat noisy, particularly in the plane containing the potassium ions and the water molecules. Alternative descriptions of the water would be consistent with the data. Residual peaks and holes (up to $\pm 2.9 \text{ e}\text{\AA}^{-3}$) near the potassium ion positions suggest that some changes could be made in their description, but we

Table 3.1

Crystal and Intensity Data Collection Summary.

(Space group is $P4/mbm$ (#127), with $z=2$.)

	Pt ₂ Cl at 300 K	Pt ₂ Cl at 22 K
Molecular Formula	K ₄ [Pt ₂ (P ₂ O ₅ H ₂) ₄ Cl] ₂ ·3H ₂ O	K ₄ [Pt ₂ (P ₂ O ₅ H ₂) ₄ Cl] ₂ ·3H ₂ O
a, Å	13.283(3)	13.170(4)
c, Å	8.125(3)	8.077(3)
crystal dimensions, mm	0.10×0.14×0.26	0.22×0.15×0.08
Absorption Coefficient μ (Mo K α) cm ⁻¹	114.8	116.9
Transmission Coefficient (min, max)	0.221, 0.343	0.193, 0.407
2 θ Limits, deg	3.0-75.0	5.0-80.0
Number of Reflections Collected	4112	3335
Unique Reflections	2095	1428
Number of Reflections with I>0	1820	1333
Refined Parameters	63	61
R ^a	0.056	0.036
GOF ^b	2.2	2.0

^a $R = (\sum |F_o - |F_c||) / \sum F_o$, based upon data with I>0^b $GOF = [\sum w(F_o^2 - F_c^2)^2 / (n-p)]^{1/3}$

did not develop a good alternative model. The maps suggest that the halide and potassium occupancies may be non-stoichiometric, but we have not been able to pursue this. Calculations were done with programs of the CRYM X-ray computing system plus ORTEP;³² scattering factors were taken from the standard compilations.³³ Hydrogen atoms were ignored in the calculations.

Pt₂Cl at 22 K. A different Pt₂Cl crystal was used for the low temperature study than that which was used for the room temperature structure. This second crystal was cooled to 22 K in incremental steps at which time the unit cell dimensions were determined using the same 15 reflections as used to determine the cell dimensions at room temperature. No extra reflections were found at low temperature, thus confirming the choice of cell dimensions.

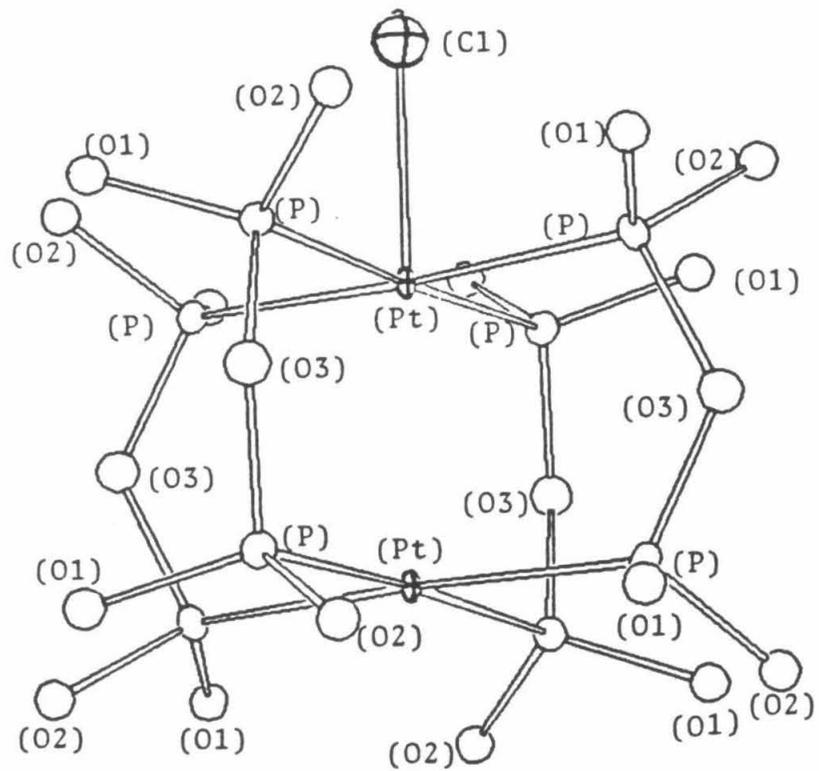
During data collection, the temperature was maintained at 22±1 K. A full orthorhombic data set was taken and used for refinement in both orthorhombic and tetragonal unit cells; 3335 reflections were scanned in the range $5^\circ < 2\theta < 80^\circ$. Systematic absences in the diffractometer data of $h0l$, where $h = 2n+1$, and $0kl$, where $k = 2n+1$, were consistent with space group P4/mbm (#127). The data were merged to give 1428 independent reflections, all of which were used in the structure solution and refinement. Although lattice cell constants and the goodness of fit did not indicate that an orthorhombic distortion had occurred, checks were made for both a doubled unit cell and for an orthorhombic distortion. Scans were taken in reciprocal space along [100], [010], [001], and [110] axes with the output of the counter circuit connected to a chart recorder. No extra reflections were observed along any axis that would indicate a radical change in cell dimension.³⁴⁻³⁶

Initially, refinement proceeded with ordered platinum atoms but two-site disordered bridging chloride ions. Because this yielded highly anisotropic platinum thermal parameters ($U_{11} = U_{22} = 0.0029(1)$, $U_{33} = 0.0082(1)$), an alternative model was refined, with the platinum atom also disordered between two sites. Allowing the two half-populated platinum atoms to move along the stacking axis gave a Pt-Pt site separation of $0.142(2)$ Å,²⁹ or Pt-Pt bond distances of $2.685(2)$ Å, $2.827(2)$ Å, or $2.969(2)$ Å, depending, as will be discussed later, on the overall translational symmetry. The final R value and goodness of fit were essentially the same as in the ordered model. An ORTEP of the Pt_2Cl unit is presented in Figure 3.2a, and a partial drawing of the structure illustrating the 1-dimensional stacking unit is shown in Figure 3.2b; atomic coordinates are given in Table 3.2.

Pt₂Cl at Room Temperature. An earlier collected and refined data set was re-evaluated based upon the findings of Pt_2Cl at 22 K. The bridging halide was represented by a two site model with isotropic thermal parameters. One fully populated water site, W(1), was found, with additional sites of lower partial population in the region near K(1). The structure was refined with both a single platinum atom site and with a two-site model; both gave the same values for the R-index and the goodness of fit. Because spectral data indicate that both Pt_2 and Pt_2Cl_2 are present in the crystal, the parameters for the model with two platinum sites are reported herein. A partial drawing of the structure illustrating the 1-dimensional stacking unit is shown in Figure 3.2b; atomic coordinates are listed in Table 3.3.

Raman Spectroscopy. The Raman data were collected by Dr. Steven D.

- Figure 3.2**
- (a) ORTEP drawing of a $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}]^{4-}$ molecule.
 - (b) ORTEP drawings of all $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{X}]^{4-}$ structures.
Ellipsoids are drawn at the 50% probability level.



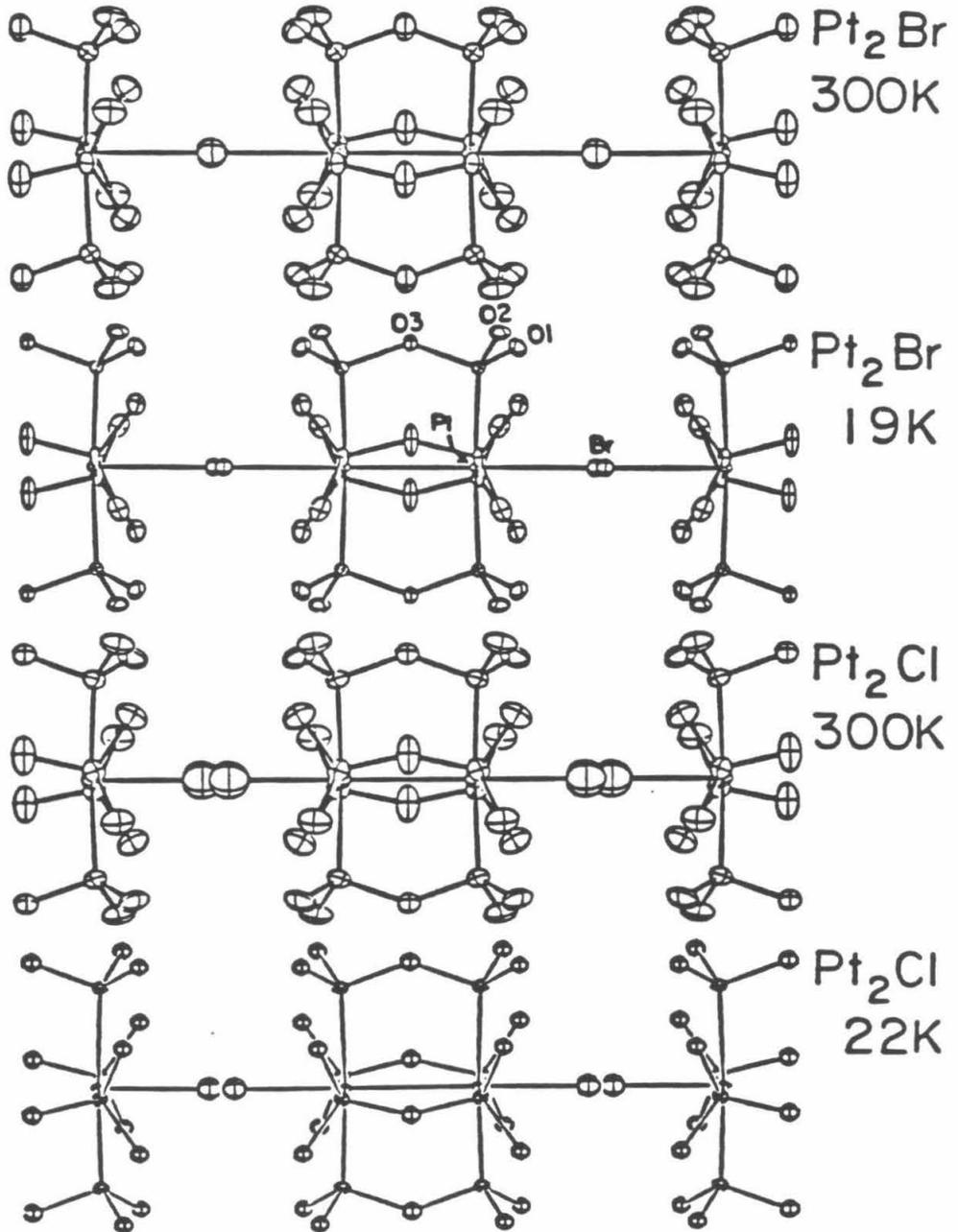


Table 3.2

Parameters for Pt₂Cl at 22 K. x, y, z and $U_{eq} \times 10^4$

Atom	x	y	z	U_{eq}^a	POP†
PtA	0	0	3338(2)	0.22(3)*	2.0
PtB	0	0	3162(1)	0.22(3)*	2.0
Cl	0	0	316(9)	287(10)	2.0
K1	2772(2)	7772	0	150(3)	4.0
K2	3892(2)	8892	5000	134(3)	4.0
P	1173(1)	1325(1)	3174(1)	90(2)	16.0
O1	2161(2)	1073(2)	2140(4)	122(6)	16.0
O2	807(2)	2353(2)	2572(4)	129(7)	16.0
O3	1659(4)	1510(3)	5000	156(10)	8.0
W1	1030(10)	6030	5000	1.5(2)*	3.0
W2	0	5000	2272(60)	1.3(6)*	0.6
W3	0	5000	909(60)	0.8(7)*	0.5
W4	649(40)	5649	0	1.4(6)*	0.7
W5	612(30)	5612	2835(41)	1.1(5)*	1.2

$$^a U_{eq} = 1/3(U_{11} + U_{22} + U_{33})$$

* B (iso)

† Atoms per unit cell

Table 3.3

Parameters for Pt₂Cl at 300 K. x, y, z and $U_{eq} \times 10^4$

Atom	x	y	z	U_{eq}^a	POP†
PtA	0	0	3321(2)	1.0(.3)*	2.0
PtB	0	0	3180(2)	1.1(.3)*	2.0
Cl	0	0	373(8)	668(18)	2.0
K1	2742(3)	7742	0	402(5)	4.0
K2	3874(3)	8874	5000	389(4)	4.0
P	1174(1)	1304(1)	3181(2)	212(2)	16.0
O1	2158(3)	1035(3)	2189(6)	340(8)	16.0
O2	833(3)	2325(3)	2570(6)	363(9)	16.0
O3	1640(5)	1492(5)	5000	437(14)	8.0
W1	1014(14)	6014	5342(21)	3.1(3)*	3.2
W2	410(54)	5410	2186(76)	7.9(14)*	1.6
W3	0	5000	586(**)	5.2(24)*	0.4
W4	632(**)	5632	0	4.1(20)*	0.4
W5	4373(**)	9373	1134(**)	3.0(26)*	0.4

^a $U_{eq} = 1/3(U_{11} + U_{22} + U_{33})$ ** esd ≥ 100

* B (iso)

† Atoms per unit cell

Conradson at the Los Alamos National Laboratory, Los Alamos, New Mexico. Raman spectra of single crystals of Pt_2Cl and Pt_2Br were obtained using 676.4 nm excitation at several temperatures in the range 25-300 K. The crystals were mounted in a cell containing *ca.* 150 psi (room temperature) of He in order to minimize local heating. The cell was coupled to the cold end of an Air Products Displex closed-cycle cryostat and the temperature monitored and controlled using thermocouples. The spectra were obtained using a SPEX 1401 double monochromator with excitation from a Spectra Physics 171 Kr^+ laser.

RESULTS

Given the 300 and 19 K structure of Pt_2Br reported earlier,³⁷ these efforts were directed towards establishing the structural variations that occur in Pt_2Cl and Pt_2Br and how they may be related to the conduction properties observed in these systems.

Structural Data on Pt_2Cl : Pt_2Cl retains a tetragonal unit cell on cooling to 22 K. The Pt_2Cl unit cell dimensions contracted with temperature in a fashion similar to that observed in Pt_2Br .³⁷ It is interesting to note that the unit cell did not double in size along any unit cell axis, although this has been seen in other 1-dimensional systems. Also, no evidence for an orthorhombic distortion was found nor was any tilting of the Pt-Cl-Pt unit with respect to the stacking axis observed.

Tables 3.2 and 3.3 list the final parameters of the models used to describe the two structures. Tables 3.4 and 3.5 compare selected bond distances and angles of Pt_2 , Pt_2Br (300 and 19 K) and Pt_2Cl (300 and 22 K). The Pt_2Cl structures were refined with a two-site model for the bridging chloride and both one- and two-site models for the platinum atoms. At both room temperature and at 22 K the chloride ion sites were well-separated and refined smoothly. At room temperature the Pt could be refined in a single site; changing to a two-site model gave nearly identical results. The resulting $(\text{AABCCB})_n$ translational symmetry is in better accord with spectral data, which will be discussed below, and thus we have accepted it. Upon cooling Pt_2Cl to 22 K, the chloride and platinum positions remain disordered, and the two-site model is retained. This model is not statistically better than a

Table 3.4

Selected Bond Distances

	Pt ₂ ^a	Pt ₂ Br ^b	Pt ₂ Br ^c	Pt ₂ Cl ^c	Pt ₂ Cl ^c
T(K)	300	300	19	300	22
Pt-Pt	2.925(1)	2.793(1)	2.781(1)	2.729(2) 2.958(2)	2.685(2) 2.969(2)
Pt-X (short)	----	2.699(1)	2.579(4)	2.281(7) 2.395(7)	2.441(7) 2.299(7)
Pt-X (long)	----	----	2.778(4)	2.886(7) 3.001(7)	2.951(7) 2.809(7)
ΔX	----	0	0.199(5)	0.606(9)	0.510(10)
Pt-P	2.320(5)	2.334(1)	2.331(1)	2.334(2) 2.331(2)	2.335(2) 2.331(2)
P-O(H)	1.579(9)	1.562(4)	1.572(4)	1.577(4)	1.581(3)
P=O	1.519(9)	1.505(4)	1.509(4)	1.514(4)	1.518(3)
P-O (bridging)	1.623(6)	1.618(3)	1.621(6)	1.622(6)	1.626(5)
P····P	2.980(6)	2.935(2)	2.924(2)	2.956(2)	2.949(1)
O(H)····O	2.505(19)	2.487(5)	2.482(5)	2.476(6)	2.479(4)
K(1)····O(H)	2.848(14)	2.915(5)	2.861(8)	2.884(10)	2.828(7)
K(1)····O	2.847(13)	2.912(6)	2.887(8)	2.873(10)	2.850(7)

^aReferences 23 and 40.^bReference 27.^cThis work.

Table 3.5

Selected Bond Angles

	Pt ₂ ^a	Pt ₂ Br ^b	Pt ₂ Br ^c	Pt ₂ Cl ^c	Pt ₂ Cl ^c
T(K)	300	300	19	300	22
Pt-Pt-P	90.67(10)	91.74(2)	91.75(3)	92.78(6)	93.25(5)
				89.97(6)	89.76(5)
Pt-P-O(H)	114.0(5)	113.4(1)	113.1(1)	114.3(2)	114.7(1)
				112.7(2)	112.7(1)
Pt-P=O	118.0(5)	117.3(2)	117.5(2)	118.7(2)	118.4(1)
				117.7(2)	117.1(1)
Pt-P-O (bridge)	110.3(4)	110.4(2)	110.4(2)	109.0(2)	108.7(2)
				111.7(2)	112.1(2)
P-O-P	133.3(9)	130.2(3)	128.9(4)	131.4(4)	130.2(3)

^aReferences 23 and 40.^bReference 27.^cThis work.

single-site platinum model (the R factor and goodness of fit are each improved by only 0.01) but, as in the room-temperature case, it has been adopted because of spectroscopic results. Listings of observed and calculated structure factors and anisotropic thermal parameters are tabulated in Appendix I.

Raman Spectroscopy: The Raman spectra of Pt_2Cl and Pt_2Br at *ca.* 25 K are shown in Figure 3.3; these spectra are similar to those reported by Kurmoo and Clark^{29b}. The wavenumbers of the bands observed for Pt_2Cl are significantly different from those reported earlier,^{29b} owing to the fact that different excitation wavelengths were used to obtain the Raman spectra. For Pt_2Cl , we employed 514.5 nm excitation, which is close to the IVCT absorption maximum, while Kurmoo and Clark employed 647.1 nm excitation. We note that the Raman spectrum of Pt_2Br also exhibits unusual relative intensity changes with excitation wavelength. The origin of these differences has been investigated by Swanson and coworkers.³⁰ Kurmoo and Clark concluded that, on the basis of X-ray work at room temperature, both Pt_2Cl and Pt_2Br consist of polar dimers (Scheme 3 with $(\text{ABC})_n$ translational symmetry) and assigned the band observed at $\sim 152\text{ cm}^{-1}$ in the Pt_2Cl spectrum to the Pt-Pt stretch. The Pt(III)-Pt(III) stretch in the fully oxidized form, Pt_2Cl_2 , is observed at 158 cm^{-1} and the Pt(II)-Pt(II) stretch in the reduced form, Pt_2 , is at 115 cm^{-1} . If a polar dimer represents the true structure for Pt_2Cl , we would expect a Pt-Pt stretching mode intermediate in wavenumber to those found for Pt_2 and Pt_2Cl_2 . Accordingly, we believe that the $\sim 152\text{ cm}^{-1}$ band (155 cm^{-1} in our spectrum) of Pt_2Cl should be assigned to the Pt(III)-Pt(III) stretch of the oxidized

unit in a chain with $(AABCCB)_n$ translational symmetry. The band at 119 cm^{-1} in Pt_2Cl is attributed to the Pt(II)-Pt(II) stretch of the reduced unit in the chain. The fact that the two bands assigned to Pt-Pt stretching in Pt_2Cl are quite close to those observed for the fully oxidized and fully reduced forms suggests strong valence localization for Pt_2Cl . The band at $\sim 301\text{ cm}^{-1}$ in Pt_2Cl is attributed to the symmetric Pt(III)-Cl stretch.

The Raman spectrum of Pt_2Br shows bands at ~ 117 and $\sim 122\text{ cm}^{-1}$. While the band at 117 cm^{-1} is quite close to the Pt(II)-Pt(II) stretching frequency of Pt_2 , the 122 cm^{-1} feature is significantly lower than the 132 cm^{-1} Pt(III)-Pt(III) stretching frequency of Pt_2Br_2 .²⁹ Furthermore, the Pt-Br stretch of Pt_2Br is 211 cm^{-1} , by contrast to 223 cm^{-1} for Pt_2Br_2 . The energies for the Pt-Br and "Pt(III)-Pt(III)" stretching modes are inconsistent with complete valence localization and $(AABCCB)_n$ translational symmetry. The two Pt-Pt stretches could result from coupling of two Pt-Pt bonds of equal or nearly equal strengths, as would be expected for Schemes 3, 4, and 5 (Figure 3.1). However, the appearance of only one band attributable to a Pt-Br stretching mode suggests that Scheme 4, the symmetrically modified $(AAB)_n$ structure, provides the best description. The bands at ~ 132 and $\sim 223\text{ cm}^{-1}$ are sample dependent in their intensities and wavenumbers (note that Clark and Kurmoo reported a band at 137 cm^{-1}). We attribute these two bands to Pt-Pt and Pt-Br stretches, respectively, of a local state. The presence of local states in these systems has been investigated by Swanson and coworkers.³⁰ The observed Raman bands and their assignments are given in Tables 3.6 and 3.7.

The relative intensities of the bands attributed to the Pt-Pt and

Figure 3.3 Raman spectra of $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{X}]^{4-}$ (X = Cl, Br) at resonance with the intervalence band.

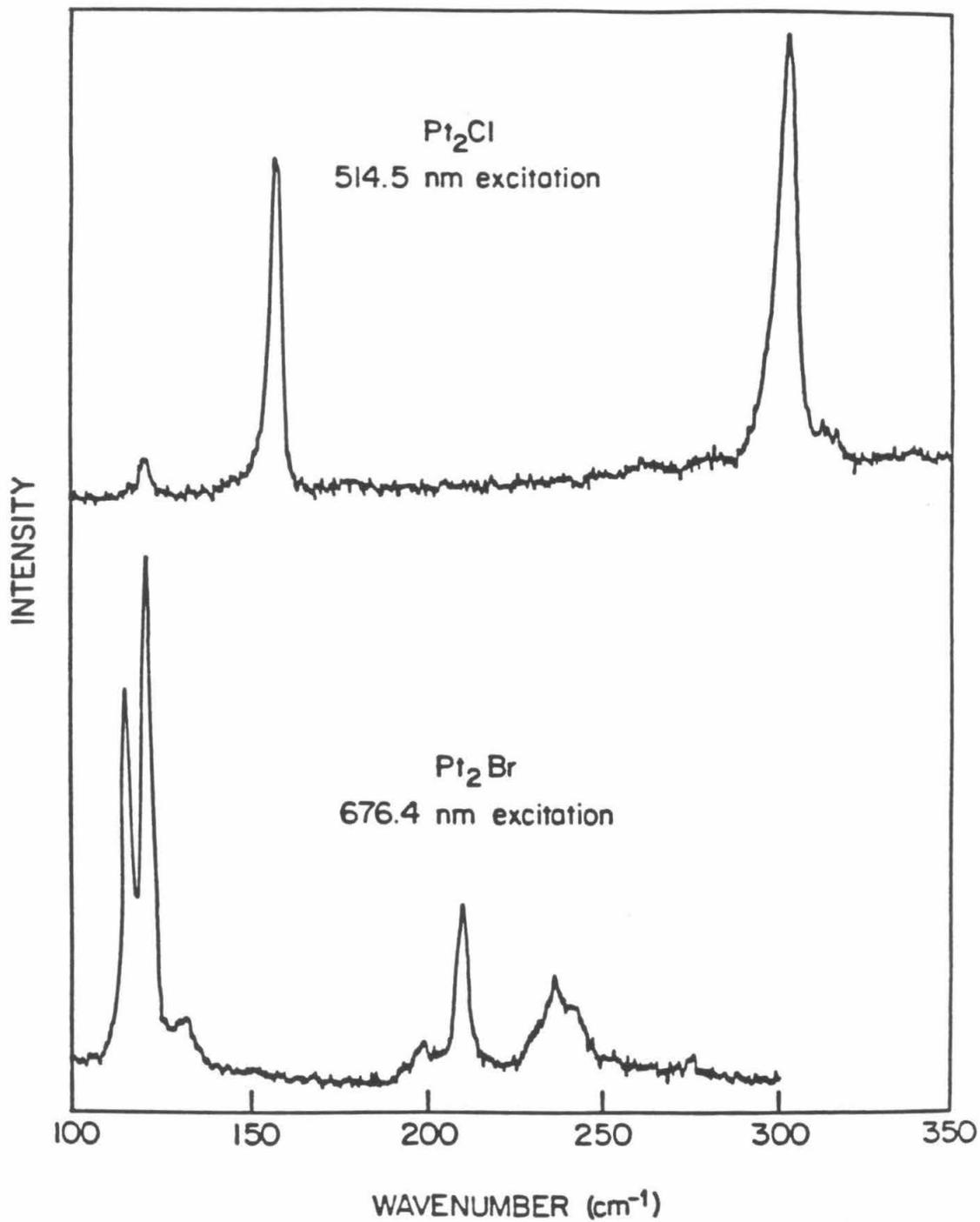


Table 3.6

Raman Bands for Pt₂Cl

Assignment	This work ^a (cm ⁻¹)	Clark and Kurmoo ^b (cm ⁻¹)
ν_1 (Pt ^{II} -Pt ^{II})	119 (mw)	125.8 (mw)
ν_2 (Pt ^{III} -Pt ^{III})	155 (vs)	152.3 (vs)
ν_3'	260 (mw)	263 (mw)
ν_3 (Pt ^{III} -Cl)	301 (vs)	291.3 (vs)
$\nu_1 + \nu_3$	419 (mw)	418 (w)
$\nu_2 + \nu_3$	453 (m)	443 (ms)
$\nu_3' + \nu_3$	557 (mw)	552 (w)
$2\nu_3$	598 (m)	583 (ms)

Abbreviations: m=medium, v=very, w=weak, s=sharp.

^a $\lambda_0=514.5$ nm, 25 K, single crystal.

^b Reference 29, $\lambda_0=647.1$ nm, 80 K, K[ClO₄] disk matrix.

Table 3.7

Raman Bands for Pt₂Br

Assignment	This work ^a (cm ⁻¹)	Clark and Kurmoo ^b (cm ⁻¹)
δ (Pt-Pt-Br)	93 (w)	93 (w)
ν_1 (Pt-Pt)	116.5 (vs)	117 (vs)
ν_2 (Pt-Pt)	122 (vs)	122 (vs)
ν_2' (Pt-Pt) ^c	132 (mw)	137 (mw)
?	193 (mw)	195 (mw)
ν_3 (Pt-Br)	211 (m)	210 (m)
ν_3' (Pt-Br) ^c	223 (w)	223 (w)
$2\nu_1$	231 (sh)	
$\nu_1 + \nu_2$	236 (m)	
$2\nu_2$	241 (m)	239 (m)
$\nu_2 + \nu_3$	332 (mw)	330 (mw)
$3\nu_2$	360 (mw)	355 (mw)
$2\nu_2 + \nu_3$	454 (w)	450 (w, br)
$4\nu_2$	474 (w,br)	475 (w, br)

Abbreviations: m=medium, v=very, w=weak, s=sharp, sh=shoulder, br=broad.

^a $\lambda_0=676.4$ nm, 25 K, single crystal.

^b Reference 29, $\lambda_0=647.1$ nm, 80 K, K[ClO₄] disk matrix.

^c Raman bands attributed to local states; see text.

Pt-X stretching modes also provide information concerning the structures of these materials. The $\nu(\text{Pt-Cl})$ is intense relative to the $\nu(\text{Pt-Pt})$ bands of the Pt_2Cl complex, while $\nu(\text{Pt-Br})$ is weak relative to the $\nu(\text{Pt-Pt})$ bands of Pt_2Br . These observations are consistent with a significant shift of the Cl atom from the central position between the dimer units but with the Br atom being quite close to this special position.

The Raman spectra of Pt_2Cl and Pt_2Br do not change significantly in going from 25 K to room temperature, indicating that there is no structural phase transformation over this temperature range. The energies and widths of the Raman bands do vary slightly with temperature, as is expected.

DISCUSSION

A tetragonal unit cell has been found for the parent dimer, Pt_2 , and the 1-dimensional materials, Pt_2Cl and Pt_2Br . The c axis length is very similar for Pt_2Cl and Pt_2Br systems, and only slightly reduced for Pt_2 . Although the c axis lengths are similar, there is a significant difference in the halide separations between the chloride (0.606(9) Å) and the bromide (0.199(5) Å)³⁷. This difference is larger than expected from the difference in radius between Cl and Br (0.15 Å), although similar to that found for Wolframs's red salts.³⁸ We conclude that the $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4$ cores and the K^+ ions form a stable arrangement. We note the disorder found in the water positions and the fact that no disorder was evident in the K^+ ion positions. Replacement of K^+ by Na^+ causes a major structural change in Pt_2 (from tetragonal to an orthorhombic unit cell);²⁵ thus, we conclude that the alkali metal ion

metal ion has an important structural influence.

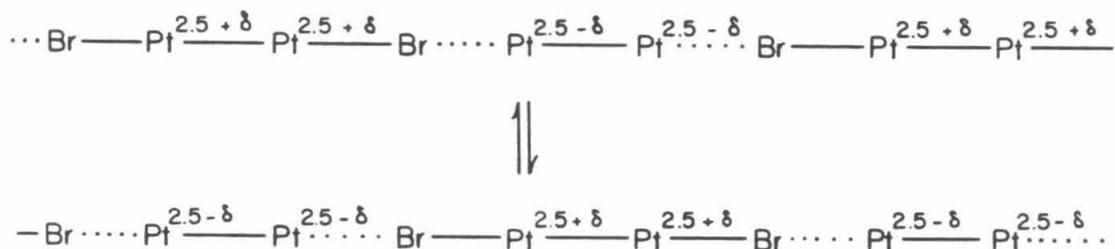
This regular array of potassium ions and Pt-Pt units fixes the coordination environment for bridging chloride or bromide (with potassium as the counter ion) as a tetragonal cavity approximately 5.34 Å long. This distance corresponds to the Pt-Pt distance between neighboring dimers along the stacking axis. From the structures of fully oxidized Pt_2X_2 complexes, one finds the following Pt-X bond distances: Pt-Cl = 2.407(2), Pt-Br = 2.572(1), and Pt-I = 2.742(1) Å.^{27,29,39} Thus a chloride ion is not constrained along the *c* axis, a bromide ion just fits, and an iodide ion is too large for the cavity (which may explain why the iodide complex does not crystallize in the same space group).

The translational symmetries $((\text{AAB})_n$, $(\text{ABC})_n$ and $(\text{AABCCB})_n$) shown in Schemes 1-5 (Figure 3.1) illustrate different structural arrangements for Pt_2X , and imply different patterns of platinum oxidation states; however, attempts by Butler *et al.*³⁷ to assign the oxidation states of the platinum atoms by ESCA did not yield much useful information. In fact, the ESCA measurements provided no conclusive evidence to suggest multiple platinum oxidation states or differing chemical environments for the platinum, bridging bromide or iodide, or ligated phosphorus atoms.

From earlier work on Pt_2Br at 19 K,³⁷ the Pt_2Br structure was fitted by a single site model for the platinum atoms and a two-site model for the bridging bromide ion. Of the translational symmetries schematically described in Figure 3.1, the only ones that permit relatively well-defined platinum atom positions and two-site disorder of

the bridging halide are Schemes 3, 4, and 5 with $(ABC)_n$, symmetrically modified $(AAB)_n$, and asymmetrically modified $(AAB)_n$ translational symmetries, respectively. The separation between the linear chains of atoms in these structures is large, 9.3 Å, and thus the interchain structural correlation is expected to be weak. A randomly ordered set of chains of Schemes 3, 4, and 5 will have equal numbers of chains pointing to the left and to the right; an average yields the X-ray diffraction results observed for Pt_2Br at 19 K.

The polar dimer of $(ABC)_n$ translational symmetry and the asymmetrically modified $(AAB)_n$ translational symmetry should yield spectra in which there are two bands attributable to Pt-Br stretches. Since only one such Raman band is observed for Pt_2Br , schemes 3 and 5 can be ruled out. Thus Scheme 4, the "intermediate" between the extremes of $(AAB)_n$ and $(AABCCB)_n$, best describes the Pt_2Br structure. Scheme 4, with the symmetrically modified $(AAB)_n$ translational symmetry, requires both a short and a long Pt-Br bond; at 19 K, bond distances of 2.579(4) and 2.778(4) Å were found. At least in the limit of a short chain, we would not expect a large energy barrier for halide "atom transfer",⁴⁰ as shown below:



Halide motion should be thermally activated and, at 300 K, would be rapid. The 300 K X-ray structural results are consistent with either

fast halide motion in an upper vibrational state, or a centrally positioned bromide. We recall, however, that the room temperature Raman spectrum shows evidence for a band attributable to a Pt-Br stretch, demonstrating that the bromide ion remains slightly shifted, even at room temperature, from the special position equidistant from neighboring Pt atoms. Scheme 4 also allows for a very slight variation in Pt-Pt bond length, such that it may or may not (depending upon the sensitivity of the technique and the value of δ) be possible to resolve two distinct Pt-Pt bonds. Raman data do, in fact, give evidence for two bands attributable to Pt-Pt stretches, these being very close in wavenumber (*ca.* 116.5 and 122 cm^{-1}), whereas the X-ray data do not require two distinct Pt-Pt bond lengths. (Comparison of the isotropic thermal ellipsoids for the Pt atoms in Pt_2Br at 19 K³⁷ and in Pt_2Cl at 22 K shows that the Pt atom thermal motion in Pt_2Br (0.074 Å) is larger than that in Pt_2Cl (0.053 Å).) Thus there may, in fact, be two distinct Pt-Pt bonds in the Pt_2Br structure. The Pt-Pt distance of 2.781(1) Å at 19 K, which reflects a formal Pt-Pt bond order of 0.5, is intermediate between that observed for the reduced dimer, Pt_2 , $d_{(\text{Pt-Pt})} = 2.925(1)$ Å^{23,41} and the fully oxidized Pt_2Br_2 , $d_{(\text{Pt-Pt})} = 2.716(1)$ Å.^{29c,39}

The translational symmetry of the Pt_2Cl structure is different from that of Pt_2Br . The Pt_2Cl structural data were refined using a two-site model for the halide and both a one-site and two-site model for the platinum centers giving nearly identical results. The Pt_2Cl Raman spectrum resembles a superposition of the Raman spectra of Pt_2 and Pt_2Cl_2 . Quite simply, Pt_2Cl appears, at both room temperature and 22

K, to be a combination of equal amounts of Pt_2 and Pt_2Cl_2 alternating along the chain. Therefore, the translational symmetry of Pt_2Cl is assigned as $(\text{AABCCB})_n$ (Scheme 2, Figure 3.1).

Whangbo and Canadell have performed band calculations using extended Hückel methods and halide p- and metal d-orbitals on a model for the Pt_2X systems, $\text{Pt}_2(\text{HCS}_2)_4\text{I}$.¹¹ They have discussed the results both in terms of a possible Peierls instability of a metallic system, and in terms of localized electronic states of a non-metallic system. In neither limit did the $(\text{ABC})_n$ translational symmetry represent an energy minimum; instead $(\text{AABCCB})_n$ was predicted to be more stable than $(\text{AAB})_n$ by 4 kJ mol^{-1} . The results show that the subtle balance of forces (i.e., electron-phonon coupling, etc.) in the Pt_2X systems gives rise to different ground state structures for Pt_2Cl and Pt_2Br .

Pt_2Br is a moderate semiconductor, with conductivity 10^6 times greater than that of Reihlen's green, $[\text{Pt}^{\text{II}}(\text{etn})_4][\text{Pt}^{\text{IV}}(\text{etn})_4\text{Br}_2] \cdot 4\text{H}_2\text{O}$ (etn = ethylamine)³⁸, but 10^6 times weaker than that of $\text{K}_2[\text{Pt}(\text{CN})_4]\text{Br}_{0.3} \cdot 3.2\text{H}_2\text{O}$. We note that an analogous system that uses iodide ion as the bridging unit between metal dimers, $\text{Pt}_2(\text{CH}_3\text{CS}_2)_4\text{I}$, forms a structure at 300 K with translational symmetry that is virtually $(\text{AAB})_n$.^{28,29} The conductivity of this material is comparable to that of Pt_2Br .

In conclusion, the translational symmetry, $(\text{AAB})_n$, which we anticipated for Pt_2Br , was not found in any structure, although the observed structure is only slightly different from this translational symmetry. Consideration of the cavity size between Pt_2 units indicates

that bromide is more appropriate than chloride for creating a highly symmetric 1-dimensional material with $(AAB)_n$ symmetry. Chloride, which is not as well suited for the K^+/Pt_2 framework as bromide, leads to the $(AABCCB)_n$ structure. The Pt_2Cl and Pt_2Br complexes provide an interesting example of how competing interactions in a quasi 1-dimensional system lead to different ground state structures. An understanding of the relationship between the structures of these materials, the nature of the local states, and the macroscopic properties such as conductivity is essential in developing materials with useful properties.

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CHAPTER 4

Pressure Dependence of the Electronic Spectra of

$K_4[Pt_2(P_2O_5H_2)_4X_n]$ (X = Cl, Br; n = 0, 1, 2)

INTRODUCTION

An important characteristic of quasi-one-dimensional, halide-bridged, mixed-valence transition metal complexes, MX, is that their physical properties may be controlled by varying the transition metal complex ions, the halogen, and external pressure.¹⁻⁷ These materials exhibit an intense intervalence charge-transfer excitation, the IVCT band, that is polarized along the chain axis. Resonance Raman spectra obtained by exciting into the IVCT band are also highly polarized along the chain axis and the frequencies of the Raman-enhanced modes have been used along with IVCT band energy as a measure of the extent of valence delocalization.^{1c} These highly anisotropic semiconductors are generally in the trapped-valence limit, with the metal atoms in alternating valence states. These and similar systems have recently been recognized as examples of commensurate charge-density-wave (CDW) systems.^{6,8,9} They exhibit a large distortion of the halogen from the central position between the metal dimers due to a Peierls instability.

Linear MMX chain complexes, $K_4[Pt_2(P_2O_5H_2)_4X] \cdot 3H_2O$, $X=Cl, Br$ (Pt_2X), are of particular interest¹⁰⁻¹⁴ because they are more valence delocalized and show smaller distortions of the halide sublattice relative to the MX systems. Early structural work on Pt_2Br indicated that the halogen was centrally located between the Pt dimers.¹¹ However, theoretical calculations suggested that there should be a distorted halide sublattice and some mixed-valence character.¹³ Recent experimental studies have shown that there is a small distortion of the halogen from the central position in Pt_2Br .^{10,14} In contrast, Pt_2Cl is in the strongly trapped-valence limit with a large distortion of the chloride from the central position (see Chapter 3).

Original interest in investigating these types of systems under high pressure was inspired by the research of Interrante and coworkers.¹⁵ They investigated similar linear chain MX complexes and found that the conductivity could be increased by as much as nine orders of magnitude at 14.0 GPa, presumably because of a decrease in the Peierls gap with increasing pressure. This tremendous pressure dependence made these types of materials good candidates for spectroscopic studies. Tanino, *et al.*⁵ measured the effects of high pressure on the lattice parameters, optical gaps, luminescence peaks, and the X-ray absorption near edge structure of MX complexes. It was concluded that the Peierls gap decreases with increasing pressure.

The effects of pressure on the electronic spectra of Pt_2Cl , Pt_2Br , $(\text{TBA})_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]$ (Pt_2), and $\text{K}_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}_2]$ (Pt_2Cl_2) have been investigated in order to gain insight into the structural changes that can occur in these materials. The energy maximum of the IVCT band provides a direct diagnostic of the extent of valence delocalization in the monohalide complexes.¹⁶ In addition, study of the other bands in the electronic spectra should yield valuable information on the effects of pressure on the electronic structures of the complexes. The magnitude and direction of the pressure-induced changes of the energy of the transitions provide information on changes in ligand field splitting and changes in the bond strength of the Pt-Pt complexes.

EXPERIMENTAL

Materials

All solvents were reagent grade and used without further purification.

The platinum dimers $K_4[Pt_2(P_2O_5H_2)_4]$, Pt_2 , and $K_4[Pt_2(P_2O_5H_2)_4]$, Pt_2Cl_2 , were prepared by published procedures.^{14,17,18} Linear chain complexes, $K_4[Pt_2(P_2O_5H_2)_4Cl]$, Pt_2Cl , and $K_4[Pt_2(P_2O_5H_2)_4Br]$, Pt_2Br , were prepared as described in the experimental section of Chapter 3.

Physical Measurements

The high pressure electronic absorption measurements were performed by Dr. Mary Ann Stroud with technical assistance from Dr. Stephen F. Agnew, Dr. Philip D. Stroud and Mr. Douglas G. Eckhart at the Los Alamos National Laboratory, Los Alamos, New Mexico.

Pressure was generated using a gasketed Merrill-Basset diamond anvil cell with an inconel gasket. The high energy working range of the cell was limited to 41700 cm^{-1} because of the strong absorption by nitrogen impurities in the diamonds. A ruby fluorescence pressure calibration method was used to determine pressures. Absorption measurements were obtained at room temperature using a Perkin Elmer 330 spectrophotometer. The instrument was modified for DAC work by using a 100 W tungsten-halogen lamp source and by adding a Harrick 5 to 1 beam condenser to the sample chamber.¹⁹

Fluorescence measurements were obtained on a SPEX Model 1403 3/4-m double monochromator equipped with a Princeton Applied Research photon counting system. Typically, the 514.5 nm line of a Spectra-Physics 171 argon laser was used as the exciting line. The cell was allowed to

equilibrate after each pressure increase and ruby fluorescence was routinely checked before and after the data collection to insure that the pressure had not varied significantly. When possible, the fluorescence from several rubies located in different parts of the cell were measured and the average value was used to estimate the pressure. Pressure gradients of as large as 0.4 GPa were observed at high pressures. Because of problems with sample degradation, the laser power at the sample was kept to less than 20 mW. For Pt_2Cl_2 , Pt_2 , and Pt_2Br , all pressure effects were reversible. Degradation of the Pt_2Cl sample was a problem. Peak positions were always reversible; however, intensity changes and weak additional features in the release spectra were observed when sample degradation had occurred.

For Pt_2Cl_2 , a 3-5% mixture of Pt_2Cl_2 in KClO_4 was prepared and mixed for several minutes in a dental amalgamator. The sample was loaded into the cell immediately after mixing. Pressure was applied to the sample until the salt fused and the sample became transparent. The cell was released to zero pressure before the run was started. The Pt_2 crystals were ground in mineral oil under flowing N_2 . When the crystals were ground in air, the electronic absorption spectra suggested that the sample was reacting. Samples prepared by minimal grinding of Pt_2Br single crystals with KClO_4 in a mortar and pestle resulted in the best obtainable electronic spectra. Single crystals of Pt_2Cl were ground in air and under flowing N_2 . Mineral oil or KClO_4 was used as the pressurizing medium.

Sample-dependent high-energy scattering was observed in the electronic spectra. After correction for this scattering tail had been made, peak locations were determined using a cursor. Polynomial fits

of all the energy vs. pressure electronic data were determined. For many of the electronic transitions, the ambient-pressure peak locations were somewhat dependent on the sample preparation, though the pressure-induced shifts were not. Most likely, the major cause of these differences was varying particle size. In order to compare the relative pressure-induced shift of the energy of a transition in the different runs, an energy correction factor was subtracted from each run. The correction factor, ν_0 , was determined by minimizing the standard deviation of a polynomial fit to the data. Wherever a significant energy difference resulted between runs, the ν_0 values are noted.

RESULTS and DISCUSSION

Pt_2 , Pt_2Cl_2 , Pt_2Cl , and Pt_2Br contain Pt_2P_8 units of D_{4h} symmetry built by four bi-nucleating pyrophosphite ligands, $[(\text{HO}_2\text{P})_2\text{O}^{2-}]$ (see Figure 1.1). The potassium salts of the fully reduced complex, Pt_2 , and the monohalides, Pt_2Cl and Pt_2Br , are isostructural with the metal-metal bond aligned along the z axis of the tetragonal unit cell, $P4/mbm$, as illustrated in Figure 3.2b. In the monohalides, the Pt_2 units form an infinite linear chain in the z direction. The larger the distortion of the halogen from the central position, the larger the trapped-valence character of the chain. In Pt_2Cl , which is in the trapped-valence limit, the halogen is displaced 0.25 Å from the central position at 22 K¹⁴ and the charge on the alternating Pt dimers is close to Pt(II)-Pt(II) and Pt(III)-Pt(III). In the Pt_2Br complex, the halogen is displaced only 0.1 Å from the central position at 19 K,¹⁴ and the charge on the Pt atoms is more delocalized. The ambient-pressure energies of all the electronic transitions studied as a function of pressure are listed in Table 4.1.

Pt_2Cl_2 . Two electronic transitions were studied as a function of pressure in Pt_2Cl_2 , $d_{xz}, d_{yz} \rightarrow d\sigma^*$ and $\sigma(X) \rightarrow d\sigma^*$ (Figure 4.1).²⁰ The $d\sigma \rightarrow d\sigma^*$ transition has been assigned to an absorption (at about 46,500 cm^{-1}) that is too high in energy to be observed in a diamond anvil cell. The smoothed absorption spectra of Pt_2Cl_2 at 0.1 and 9 GPa are shown in Figure 4.2.

The $d_{xz}, d_{yz} \rightarrow d\sigma^*$ transition shifts blue with increasing pressure at a rate of 300 $\text{cm}^{-1}/\text{GPa}$ from its position at 28,100 cm^{-1} . The shift is linear over the pressure range studied (Figure 4.3a). Increasing pressure enhances the overlap between the d_{z^2} orbitals on adjacent Pt

Table 4.1

Ambient-Pressure Energies of Electronic Transitions (in cm^{-1})

Transition	Pt ₂ Cl ₂	Pt ₂	Pt ₂ Br	Pt ₂ Cl
$d_{xz}, d_{yz} \rightarrow d\sigma^*$	28,100			
$\sigma(X) \rightarrow d\sigma^*$	34,500		32,800	35,400
$d\sigma^* \rightarrow p\sigma$		26,700	27,600	26,400
$d_{xz}, d_{yz} \rightarrow p\sigma$		40,200		
$d\sigma^* \rightarrow d_{x^2-y^2}$		35,300		
unassigned				25,600
IVCT			16,100	18,600

Figure 4.1 Energy level diagram for $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}_2]^{4-}$ at ambient pressure.

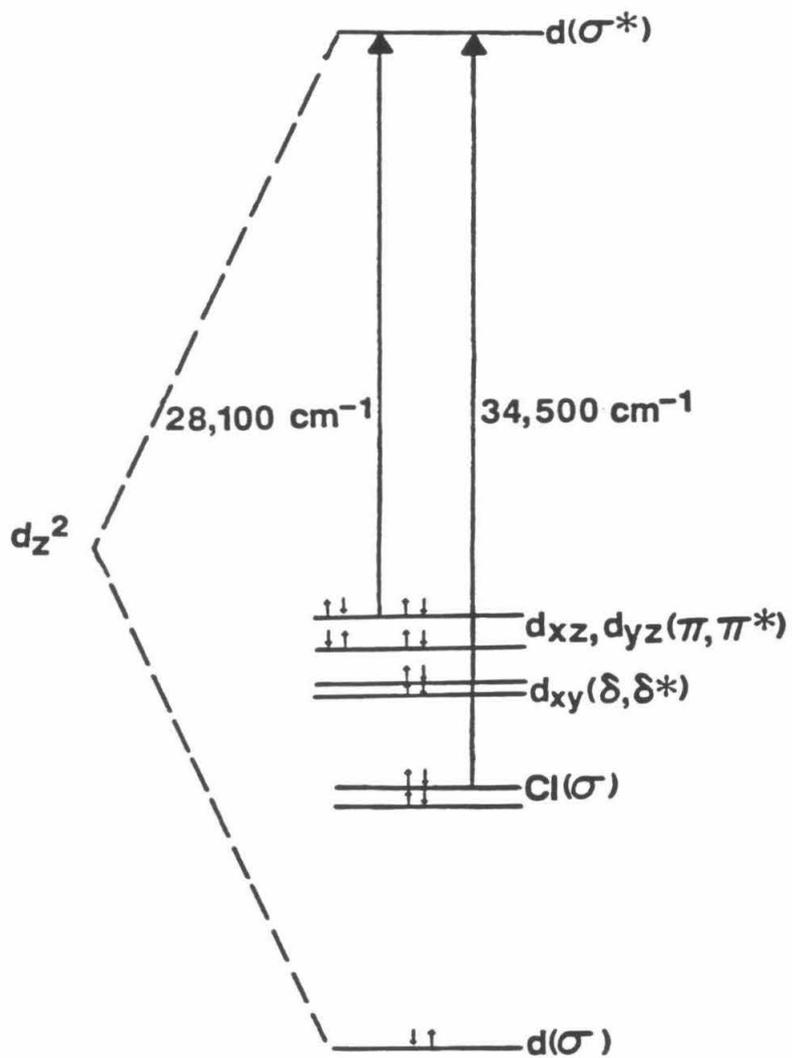
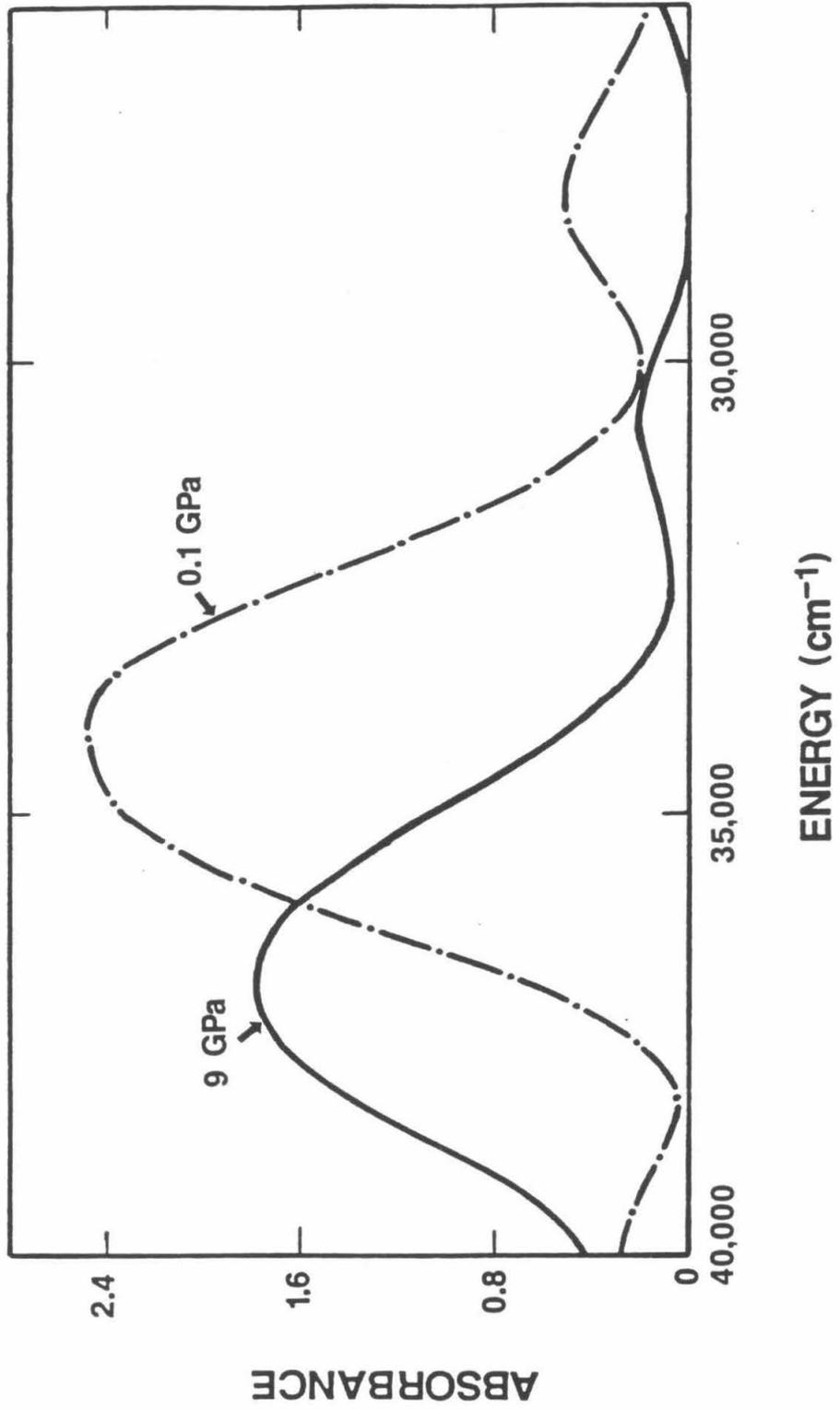
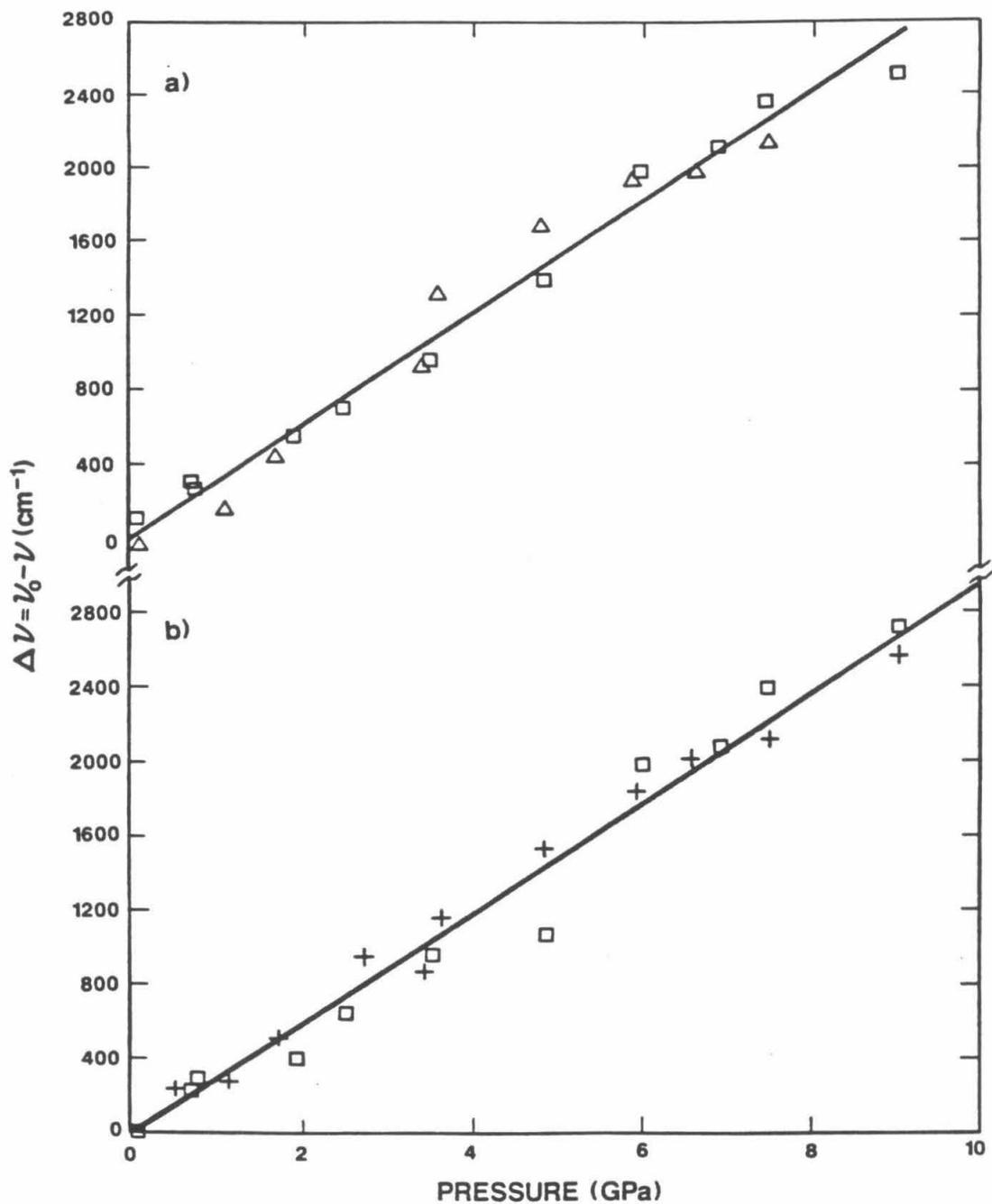


Figure 4.2 Smoothed electronic absorption spectra of $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}_2]^{4-}$ at 0.1 and 9.0 GPa.



- Figure 4.3**
- (a) Pressure shift of the $d_{xz}, d_{yz} \rightarrow d\sigma^*$ transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}_2]^{4-}$.
 - (b) Pressure shift of the $\sigma(\text{Cl}) \rightarrow d\sigma^*$ transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}_2]^{4-}$.



atoms, thereby destabilizing $d\sigma^*$. The two sets of d_{xz}, d_{yz} orbitals interact only weakly, so the shift of the $d_{xz}, d_{yz} \rightarrow d\sigma^*$ transition is attributed to the pressure-induced destabilization of $d\sigma^*$. Any increase in the splitting of the d_{xz}, d_{yz} orbitals would reduce the magnitude of the observed blue shift.

The energy of the $\sigma(Cl) \rightarrow d\sigma^*$ transition also blue shifts linearly, at a rate of $300 \text{ cm}^{-1}/\text{GPa}$ from $34,500 \text{ cm}^{-1}$ (Figure 4.3b). Accordingly, this blue shift is attributed mainly to the destabilization of $d\sigma^*$ with increasing pressure.

Pt₂. The MO diagram of the d^8-d^8 Pt₂ complex is given in Chapter 1 (Figure 1.2); however, for convenience the energy levels are presented again in Figure 4.4. The effects of pressure on three electronic transitions, $d_{xz}, d_{yz} \rightarrow p\sigma$, $d\sigma^* \rightarrow p\sigma$, and $d\sigma \rightarrow d_{x^2-y^2}$,²¹ have been investigated (Figure 4.5).

The $d_{xz}, d_{yz} \rightarrow p\sigma$ transition in Pt₂ ($40,200 \text{ cm}^{-1}$ at ambient pressure) exhibits a linear blue shift of $190 \text{ cm}^{-1}/\text{GPa}$ with increasing pressure (Figure 4.6a). An explanation of the observed shift is that the p_z orbitals are coupled strongly to d_{z^2} orbitals on adjacent Pt atoms, while the d_{xz}, d_{yz} orbitals experience only very weak π interactions. This pattern of interactions with increasing pressure results in a destabilization of p_z relative to d_{xz}, d_{yz} .

The $d\sigma^* \rightarrow d_{x^2-y^2}$ energy (ambient-pressure energy of $35,400 \text{ cm}^{-1}$) exhibits unusual nonlinear behavior with increasing pressure. The $d\sigma^* \rightarrow d_{x^2-y^2}$ transition shifts red up to 5.5 GPa at an initial rate of about $230 \text{ cm}^{-1}/\text{GPa}$, then levels off and begins to shift blue until at least 10.0 GPa (Figure 4.6b). This transition is expected to be affected both by increased splitting of $d\sigma$ and $d\sigma^*$, resulting from

Figure 4.4 Energy level diagram for $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$ at ambient pressure.

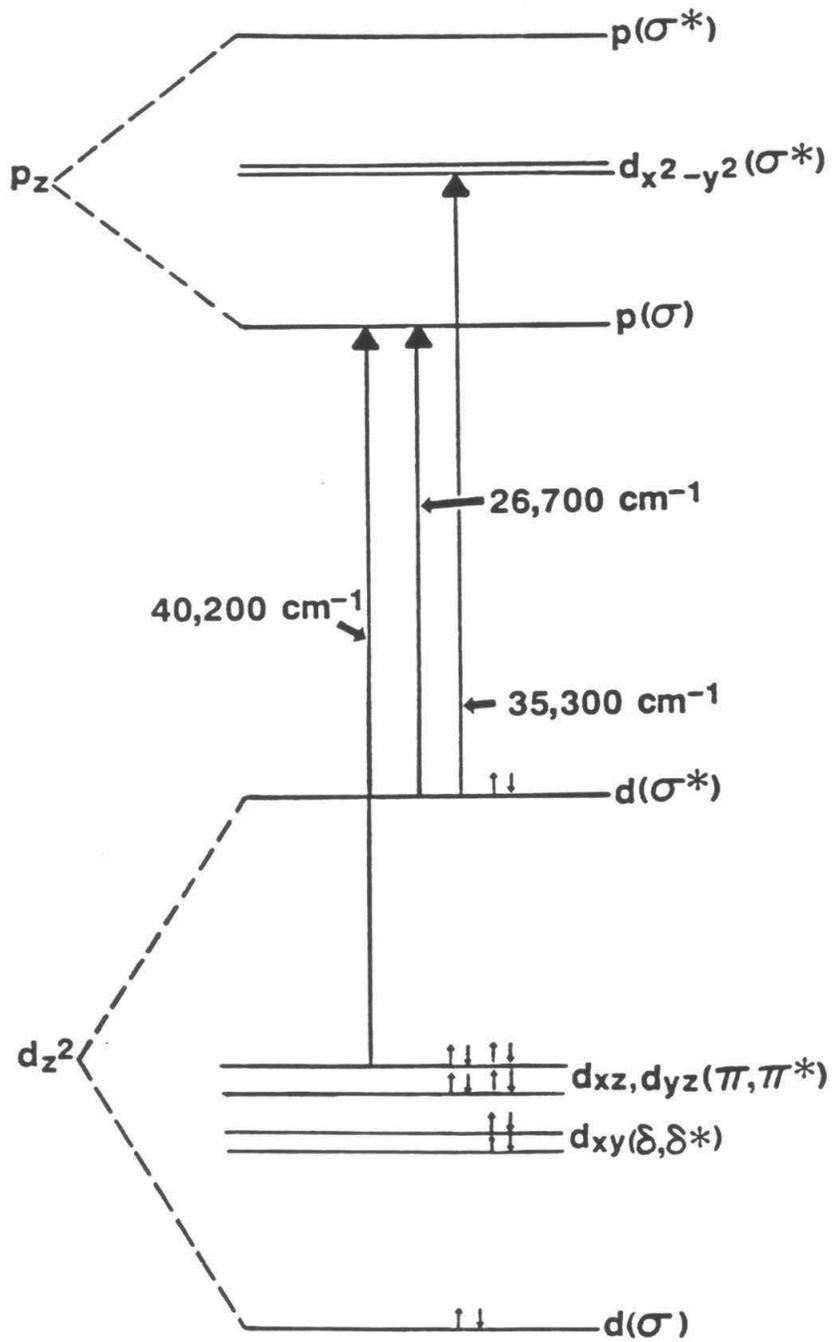
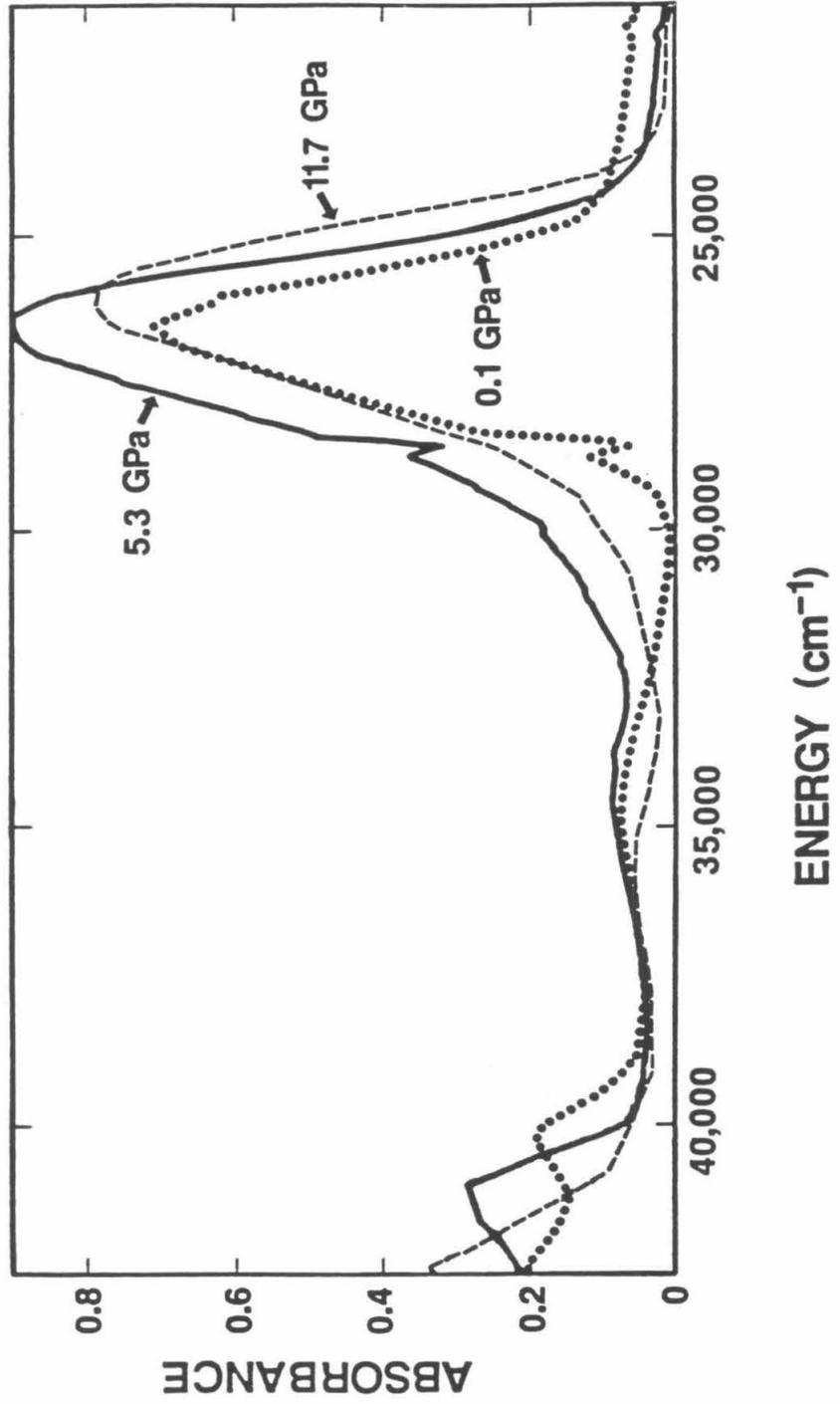
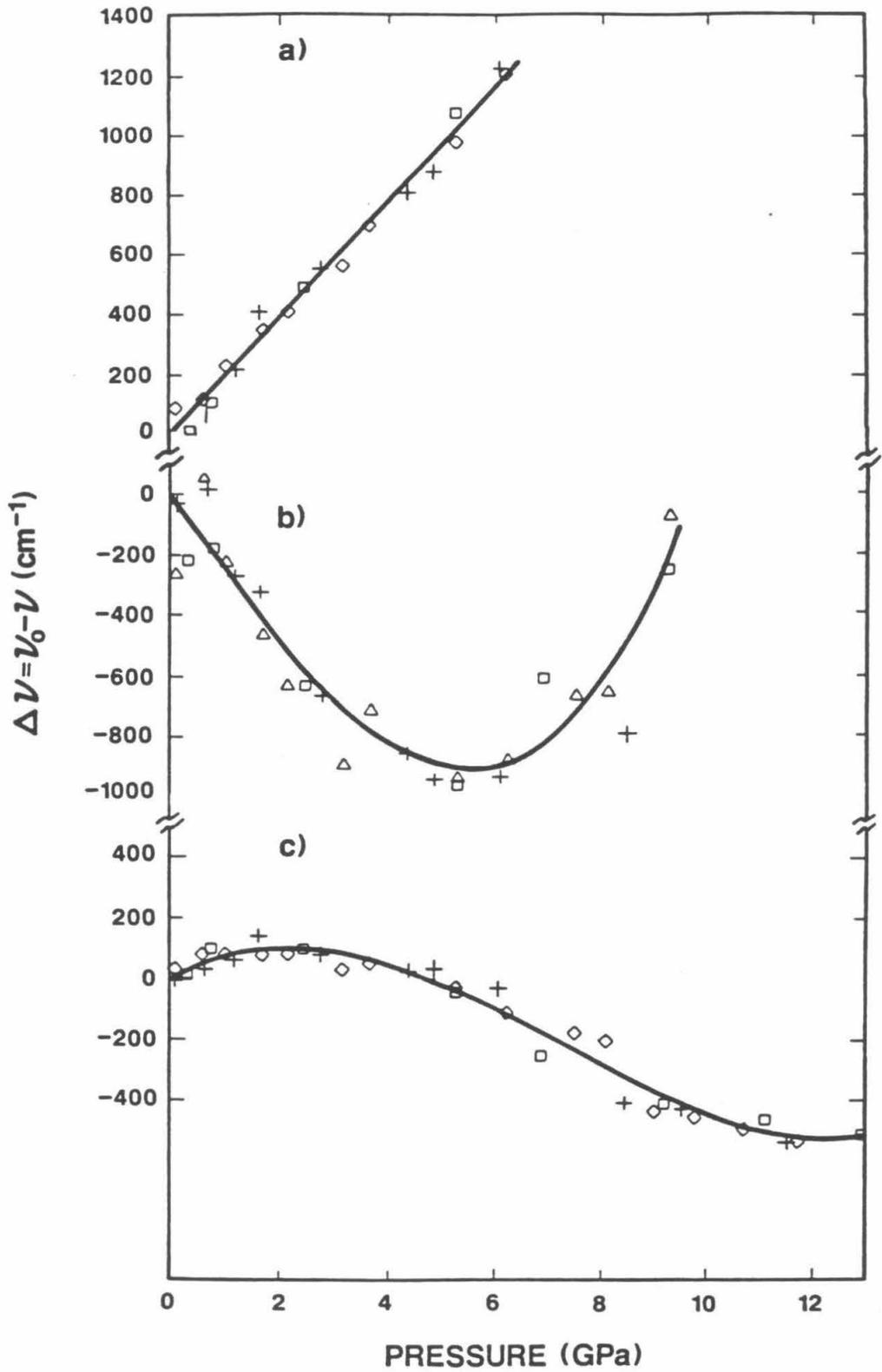


Figure 4.5 Smoothed electronic absorption spectra of $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$ at 0.1, 5.2, and 11.7 GPa.



- Figure 4.6**
- (a) Pressure shift of the $d_{xz}, d_{yz} \rightarrow p\sigma$ transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$.
 - (b) Pressure shift of the $d\sigma^* \rightarrow d_{x^2-y^2}$ transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$.
 - (c) Pressure shift of the $d\sigma^* \rightarrow p\sigma$ transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$.



greater overlap of the d_{z^2} orbitals on adjacent Pt atoms, and an increased ligand field that destabilizes $d_{x^2-y^2}$ relative to d_{z^2} . In the absence of large changes in the ligand field splitting, an increase in the overlap of adjacent d_{z^2} orbitals would lead to a red shift, and it is reasonable that this is the dominant effect at low pressure. At higher pressures, increased ligand field splitting between $d_{x^2-y^2}$ and d_{z^2} begins to dominate and the $d\sigma^* \rightarrow d_{x^2-y^2}$ transition shifts to the blue. The strong nonlinearity of the ligand field effect is attributed to nonlinear compression of the P-Pt bonds as the pressure is increased. At low pressure, compression along the *a* and *b* crystallographic axes results in compression of the interchain separation and little or no change in the P-Pt bonds. At higher pressure, as the interchain interactions stiffen, the P-Pt bonds begin to compress, thereby destabilizing $d_{x^2-y^2}$.

There is little shift in the 26,700 cm^{-1} energy of $d\sigma^* \rightarrow p\sigma$ up to 5.0 GPa, though the transition does appear to initially shift slightly to the blue. At higher pressures, a significant red shift is observed; this shift appears to level off at ca. 11.0 GPa. Increased splitting of $d\sigma$ and $d\sigma^*$ would result in a red shift of the $d\sigma^* \rightarrow p\sigma$ transition. The $d\sigma^*$ and $p\sigma$ levels are coupled, and any increase in the coupling as the pressure is raised would tend to repel $d\sigma^*$ and $p\sigma$, resulting in a blue shift. For the $d\sigma^* \rightarrow p\sigma$ transition, there is probably a delicate balance between these two effects. Because a weak shift is observed, other effects also may significantly influence the behavior of this transition. At higher pressures, when the transition begins to red shift, increased coupling between the phosphorus σ orbital and $d\sigma^*$ occurs. Increased $\sigma(\text{P})-d_{z^2}$ coupling is expected at higher pressures in

light of the observed pressure shift of the $d\sigma^* \rightarrow d_{x^2-y^2}$ transition. This effect suggests that the energy of the $d\sigma^*$ orbital may be tuned in binuclear complexes by varying the bridging ligand. For monohalide complexes, the ligands may be used to control the extent of delocalization of the chain.

The weak $22,000 \text{ cm}^{-1}$ absorption attributable to the triplet $d\sigma^* \rightarrow p\sigma$ transition²¹ initially shifts weakly blue with increasing pressure. The weak $31,800 \text{ cm}^{-1}$ shoulder is probably the $d\sigma^* \rightarrow d_{x^2-y^2}$ triplet,²¹ which shifts red with increasing pressure.

Pt₂Br. A diagram of the electronic transitions in Pt₂Br is shown in Figure 4.7. The absorption spectrum of Pt₂Br is a superposition of electronic transitions in Pt₂ and Pt₂Br₂, modified slightly depending on the degree of delocalization, and an additional broad low energy band, which is the IVCT band characteristic of the mixed-valence complex itself (Figure 4.8). The highest energy transition observed in this study, $32,800 \text{ cm}^{-1}$ at ambient pressure, shifts to higher energy with increasing pressure at a rate of $270 \text{ cm}^{-1}/\text{GPa}$ (Figure 4.9a). The rate of the shift is within 10% of the observed shift of the $\sigma(\text{Cl}) \rightarrow d\sigma^*$ transition in Pt₂Cl₂ and it is linear over the pressure range studied. The band position at ambient pressure is very close to that of $\sigma(\text{Br}) \rightarrow d\sigma^*$ in Pt₂Br₂ ($32,800 \text{ cm}^{-1}$),^{17,22} and the absorption is intense. Based on this evidence, the band is assigned to the $\sigma(\text{Br}) \rightarrow d\sigma^*$ transition.

The pressure-induced energy change of the $27,600 \text{ cm}^{-1}$ transition is shown in Figure 4.9b. The energy of the transition shifts weakly to the red with increasing pressure and appears to level off above 5.0 GPa. Based on our analysis of the Pt₂Cl₂ and Pt₂ results, a red shift

Figure 4.7 Energy level diagram for $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Br}]^{4-}$ at ambient pressure.

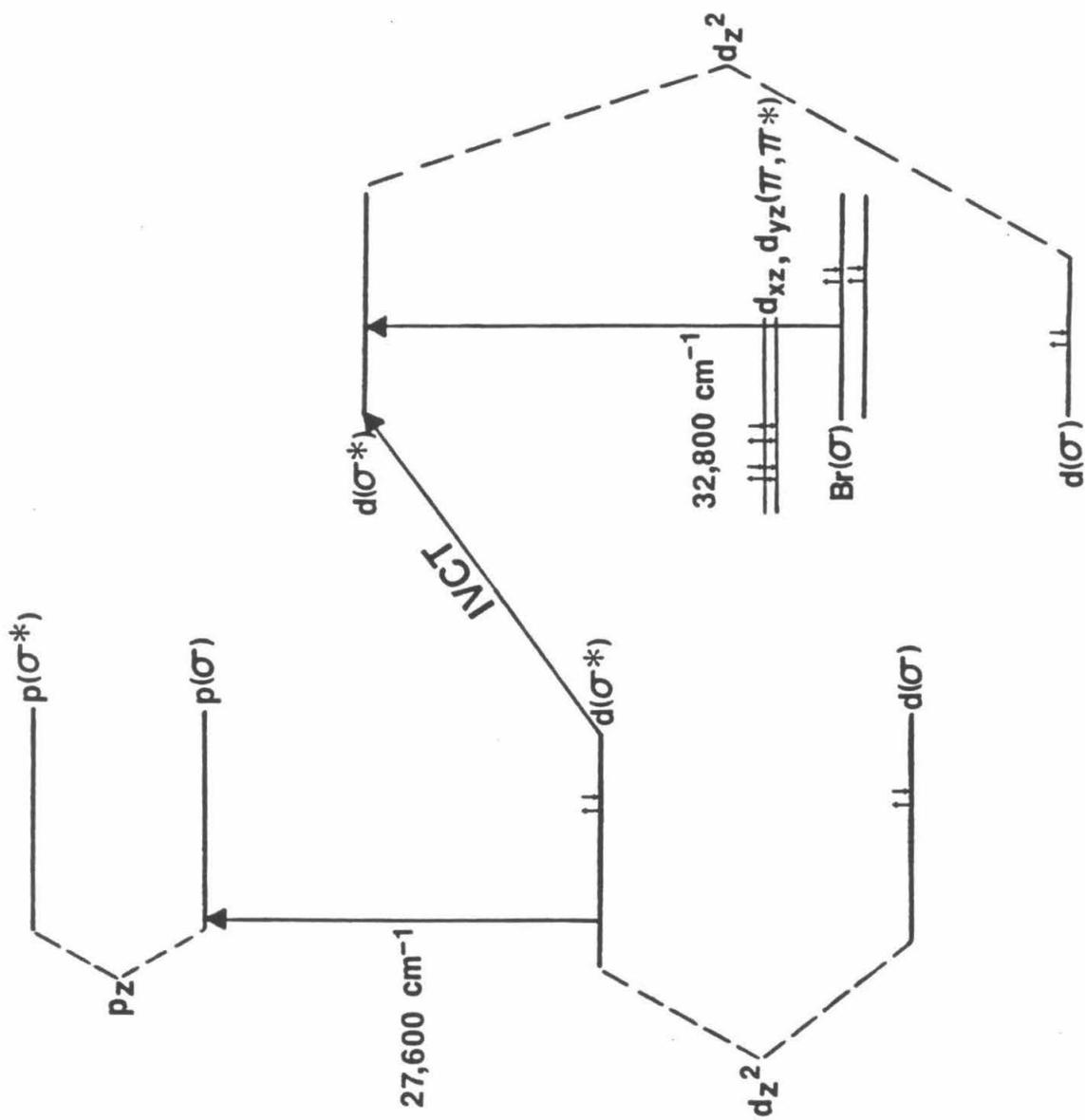
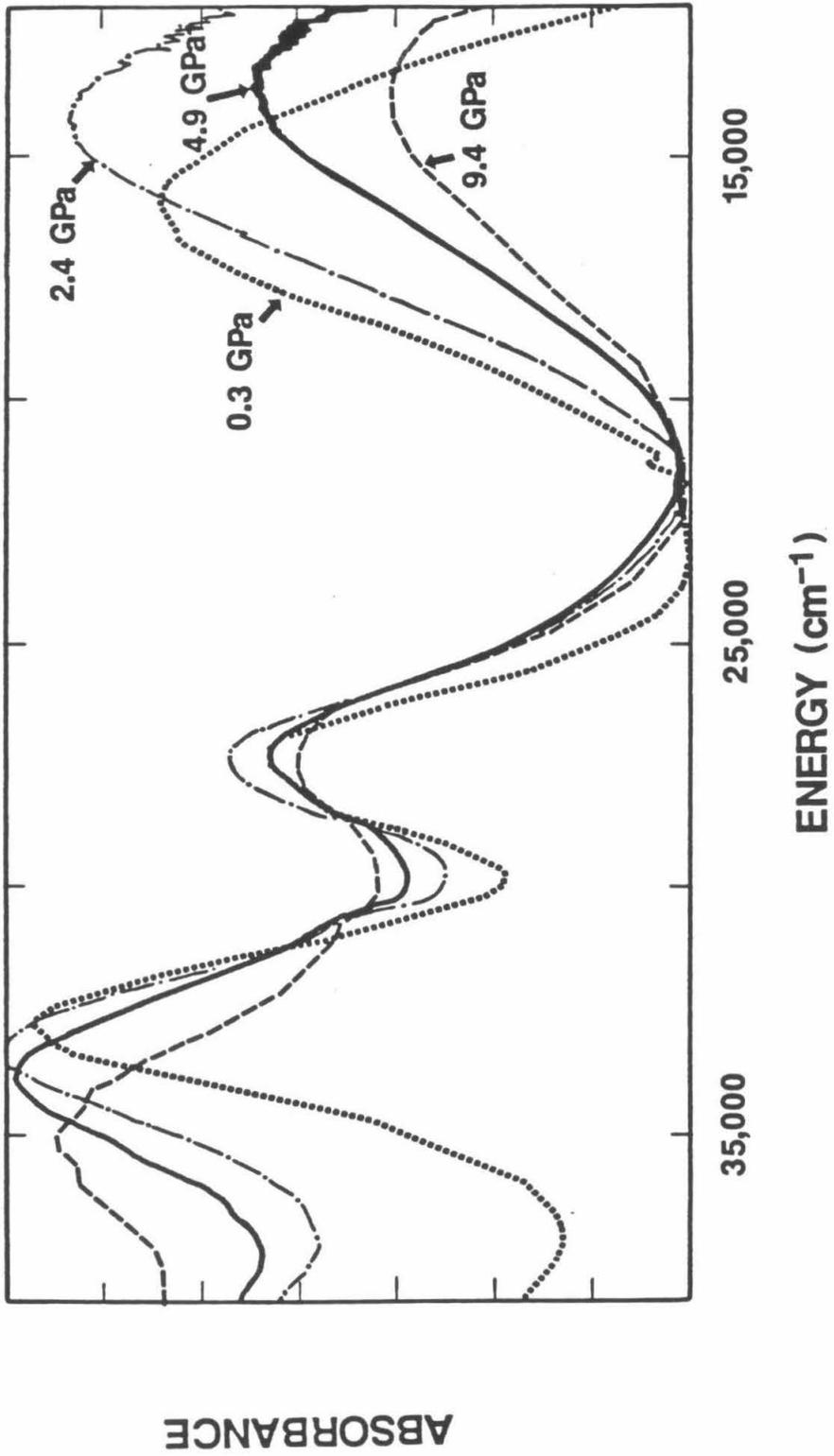
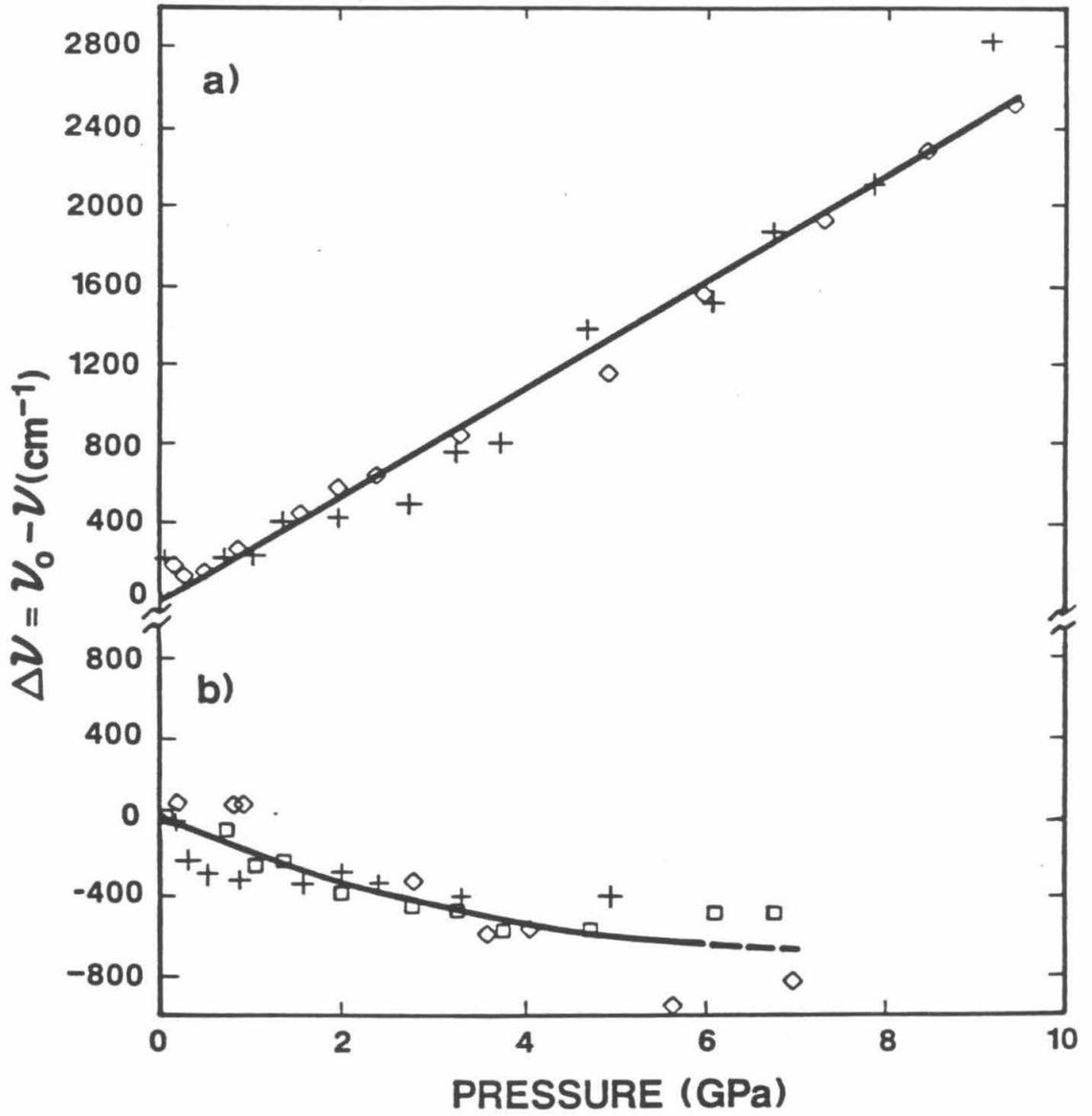


Figure 4.8 Smoothed electronic absorption spectra of $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Br}]^{4-}$ at 0.3, 2.4, 4.9, and 9.4 GPa.



- Figure 4.9**
- (a) Pressure shift of the $\sigma(\text{Br}) \rightarrow d\sigma^*$ transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Br}]^{4-}$.
 - (b) Pressure shift of the $d\sigma^* \rightarrow p\sigma$ transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Br}]^{4-}$ ($\nu_0 = 27,200, 27,600, \text{ and } 27,900 \text{ cm}^{-1}$).



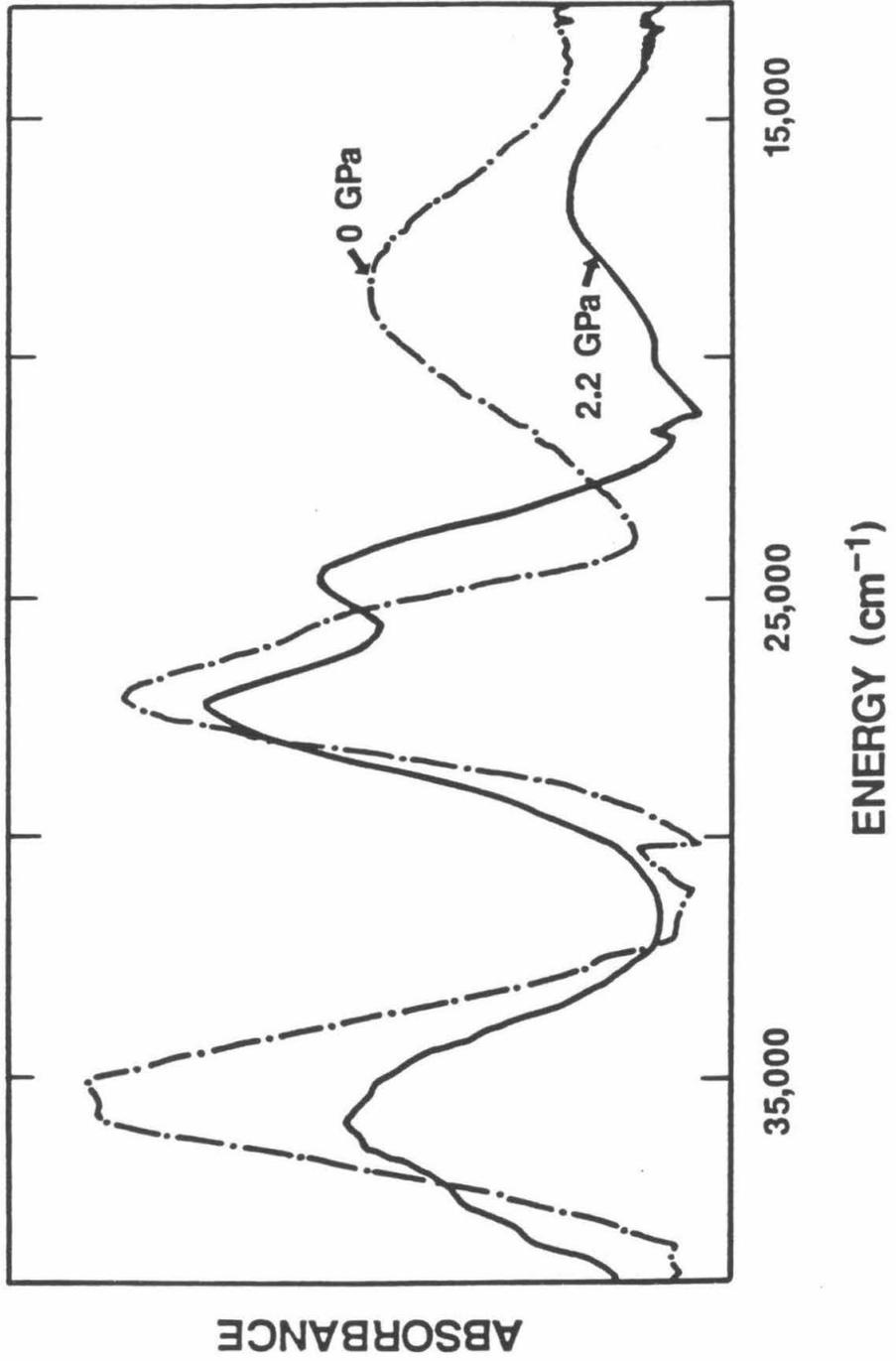
of $d\sigma^* \rightarrow p\sigma$ is plausible, while a blue shift is expected for $d_{xz}, d_{yz} \rightarrow d\sigma^*$. In Pt_2Br , a red shift of $d\sigma^* \rightarrow p\sigma$ also is expected at high pressure, because of enhanced intermolecular interactions along the metal-metal axis. Accordingly, the intense $27,600 \text{ cm}^{-1}$ band is assigned to $d\sigma^* \rightarrow p\sigma$.

Pt_2Cl . The energy level diagram for Pt_2Br (Figure 4.7) can also be utilized in discussing the spectrum of Pt_2Cl , where four electronic transitions have been studied as a function of pressure (Figure 4.10). The highest energy transition observed for Pt_2Cl ($35,400 \text{ cm}^{-1}$ at ambient pressure) shifts to higher energy with increasing pressure at a rate of $290 \text{ cm}^{-1}/\text{GPa}$ (Figure 4.11a). The shift is linear with pressure and its magnitude is within 10% of the observed shift of the $\sigma(Cl) \rightarrow d\sigma^*$ transition in Pt_2Cl_2 . The energy of the transition at ambient pressure is very close to that of $\sigma(Cl) \rightarrow d\sigma^*$ in Pt_2Cl_2 (Table 4.1). Since the transition is also very intense, it is assigned to $\sigma(Cl) \rightarrow d\sigma^*$.

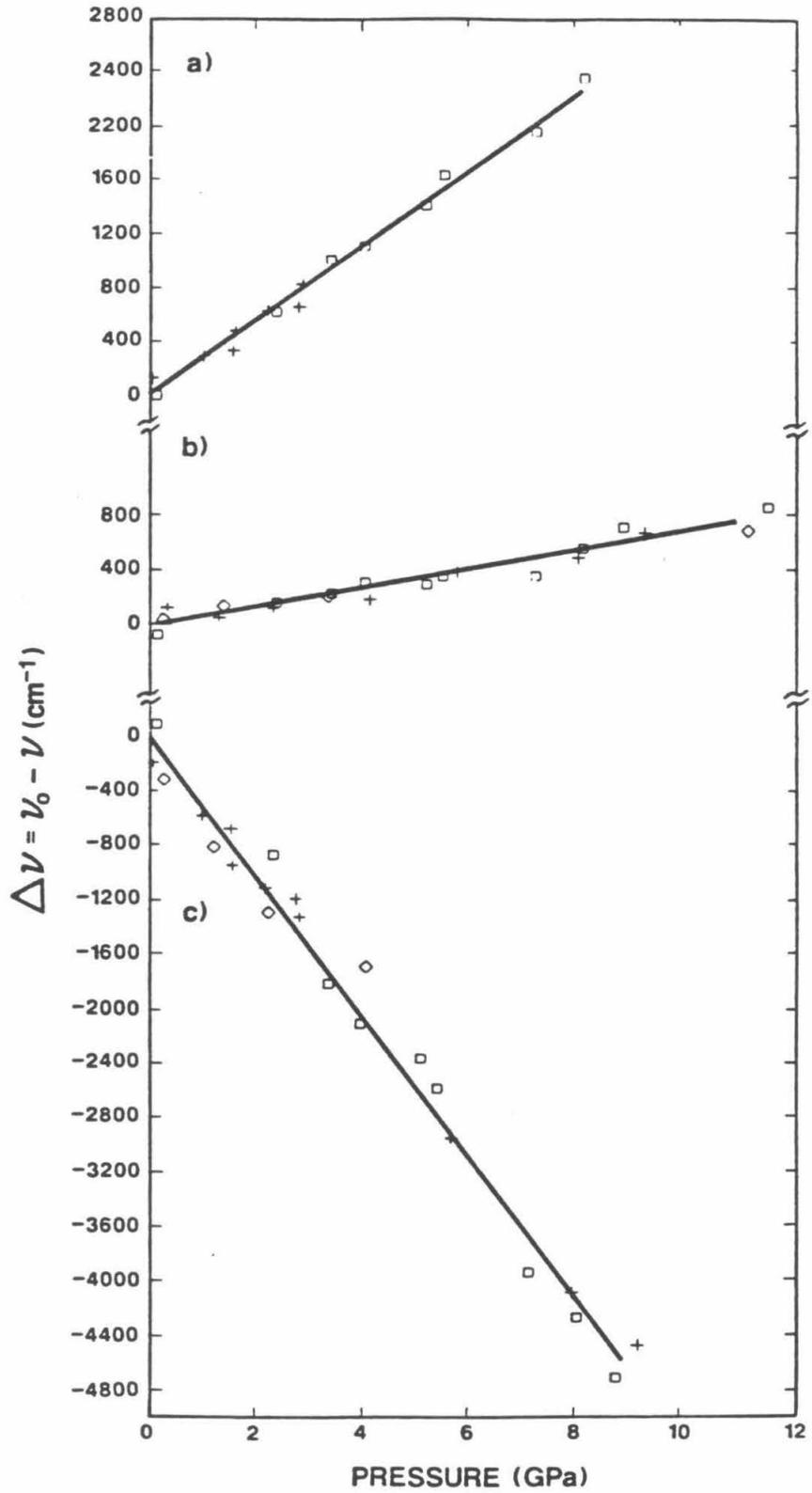
The weak blue shift, $75 \text{ cm}^{-1}/\text{GPa}$, of the transition at $26,700 \text{ cm}^{-1}$ is illustrated in Figure 4.11b. A blue shift of the $d_{xz}, d_{yz} \rightarrow d\sigma^*$ band has been observed in Pt_2Cl_2 and can be inferred in Pt_2 by taking the difference between the pressure-induced shift of $d_{xz}, d_{yz} \rightarrow p\sigma$ and that of $d\sigma^* \rightarrow p\sigma$; however, the shifts in Pt_2Cl_2 and Pt_2 are over twice that of the blue shift observed for the $26,700 \text{ cm}^{-1}$ transition in Pt_2Cl . A blue shift is also plausible for $d\sigma^* \rightarrow p\sigma$ if an increase in the coupling between $d\sigma^*$ and $p\sigma$ is the dominant pressure effect. Accordingly, the intense $26,700 \text{ cm}^{-1}$ band is assigned to the $d\sigma^* \rightarrow p\sigma$ transition.

There is a very strong pressure-induced linear red shift ($-510 \text{ cm}^{-1}/\text{GPa}$) of the Pt_2Cl electronic transition with an ambient-pressure

Figure 4.10 Smoothed electronic absorption spectrum of $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}]^{4-}$ at 0 and 2.2 GPa.



- Figure 4.11**
- (a) Pressure shift of the $\sigma(\text{Cl}) \rightarrow d\sigma^*$ transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}]^{4-}$.
 - (b) Pressure shift of the $d\sigma^* \rightarrow p\sigma$ transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}]^{4-}$ ($\nu_0 = 26,200, 26,800,$ and $27,100 \text{ cm}^{-1}$).
 - (c) Pressure shift of the unassigned transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}]^{4-}$.



energy of $25,600 \text{ cm}^{-1}$. A weak red shift of the $d\sigma^* \rightarrow p\sigma$ transition has been observed for Pt_2Br and at pressures greater than 5.0 GPa for Pt_2 . The magnitude of the red shift of the $25,600 \text{ cm}^{-1}$ band is several times larger and for this reason assignment to $d\sigma^* \rightarrow p\sigma$ is unlikely. (In addition, the $26,700 \text{ cm}^{-1}$ has already been assigned to the $d\sigma^* \rightarrow p\sigma$ transition.) Another unusual characteristic of this transition is that the intensity of the peak varies considerably with the method of sample preparation. The peak has been observed previously in the electronic spectra of samples in KClO_4 pellets,¹⁰ and there is no peak in the Pt_2Br electronic spectrum that exhibits similar behavior. Because of these observations, the peak has not been assigned to a metal-based transition of the ground-state structure. The transition may result from a defect state present in the crystal. Resonance Raman spectra of Pt_2Cl show evidence for fine structure that grows in when the excitation wavelength is tuned to the red of the IVCT band. This fine structure was attributed to a polaronic local state resulting from a deficiency of K^+ ions.¹² On the basis of the Peierls-Hubbard $3/4$ filled band model developed by Baeriswyl and Bishop,⁶ the energies of the defect electronic transitions are predicted to be to the blue as well as to the red of the IVCT band.

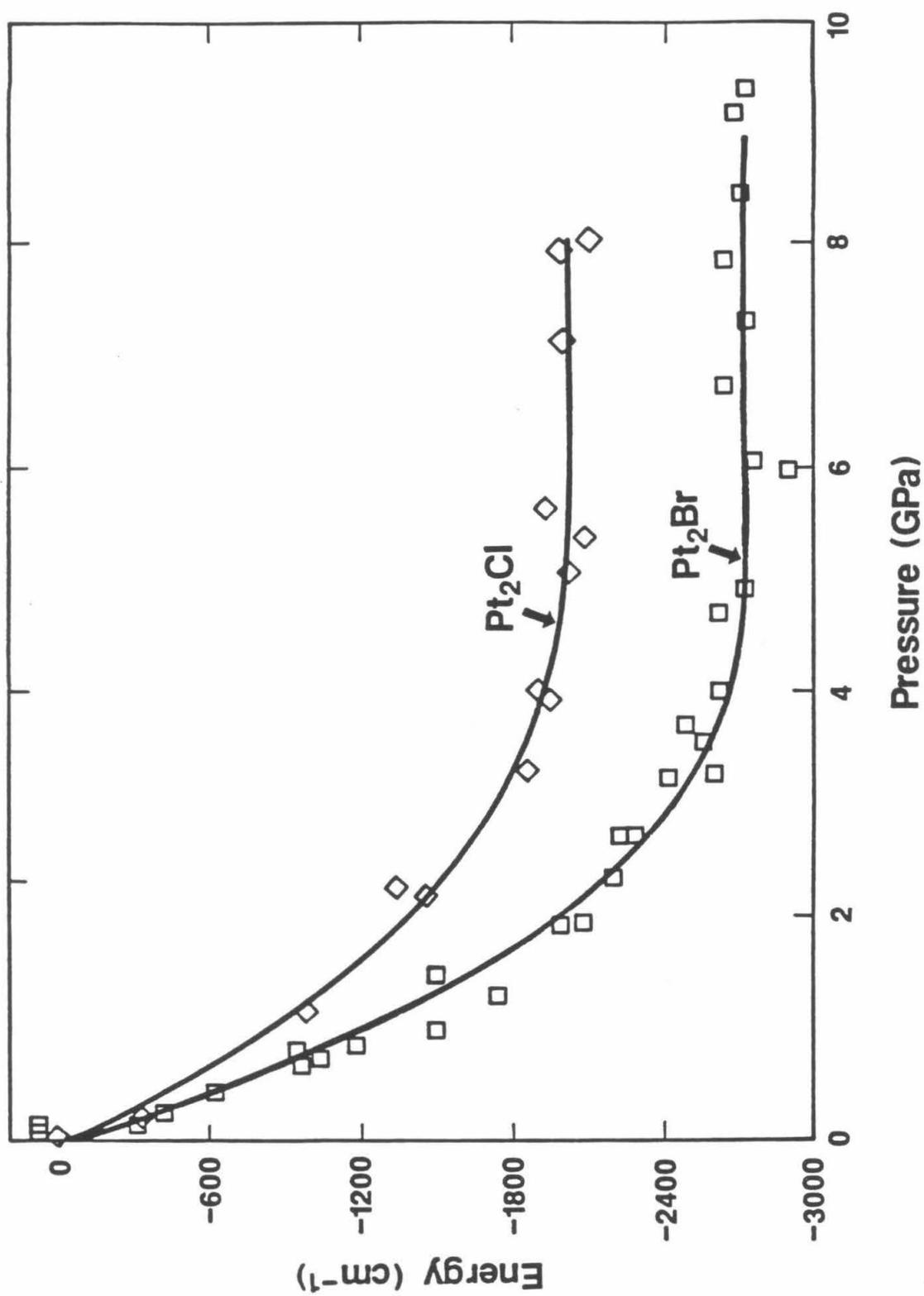
Intervalence Charge-Transfer Excitation. In the limit of a strongly trapped-valence material with a large distortion of the halogen from the central position, the IVCT band corresponds to $d\sigma^* \rightarrow d\sigma^*$ electron transfer along the chain axis from Pt_2 to an adjacent Pt_2X_2 (Figure 4.9). The energy of the IVCT band at ambient pressure in the trapped-valence Pt_2Cl complex is $18,600 \text{ cm}^{-1}$; it is significantly lower ($16,100 \text{ cm}^{-1}$) in the more delocalized Pt_2Br complex. (The lower

the energy of the IVCT band, the more delocalized the complex.¹⁶⁾ In Pt_2Cl and Pt_2Br , as the pressure increases, the IVCT band initially shifts strongly to lower energy and then appears to level off above 4.0 GPa (Figure 4.12). The initial red shift is approximately $-750 \text{ cm}^{-1}/\text{GPa}$ in Pt_2Cl and $-1000 \text{ cm}^{-1}/\text{GPa}$ in Pt_2Br . By way of comparison, the absorption edge of Wolfram's red salt shifts to lower energy at a rate of $-1600 \text{ cm}^{-1}/\text{GPa}$ up to 3.4 GPa.^{5c}

The strong red shift of the IVCT band with increasing pressure is consistent with a continuous change toward a symmetric structure with the halogen located equidistant between adjacent Pt dimers. As the halogen becomes more centrally located, there will be a decrease in the energy difference between $d\sigma^*$ orbitals on adjacent Pt dimers. The red shift of the IVCT band will be further enhanced by a decrease in the halide distortion, because the extent of delocalization along the z axis will increase the band character of the $d\sigma^*$ orbitals. Band calculations by Whangbo and Canadell on a related MMX system support these conclusions.¹³

In interpreting pressure-induced changes of the IVCT band, it is important to separate the effects of structural changes from those of changes in orbital overlap. In this case we can establish an upper limit to the red shift caused by increasing overlap by comparing the shifts of the IVCT bands in Pt_2Cl and Pt_2Br . The red shift observed for the IVCT band in Pt_2Br is 4/3 times as large as that observed for Pt_2Cl . Accordingly, at least the increased red shift of the IVCT band in Pt_2Br is attributed to the movement of the Br atom towards the center position relative to the surrounding Pt_2 fragments.

Figure 4.12 Relative energy shifts of the IVCT bands in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Br}]^{4-}$ ($\nu_0 = 15,700, 16,200, \text{ and } 16,400 \text{ cm}^{-1}$) and $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}]^{4-}$ ($\nu_0 = 18,400, 18,500, \text{ and } 18,800 \text{ cm}^{-1}$).

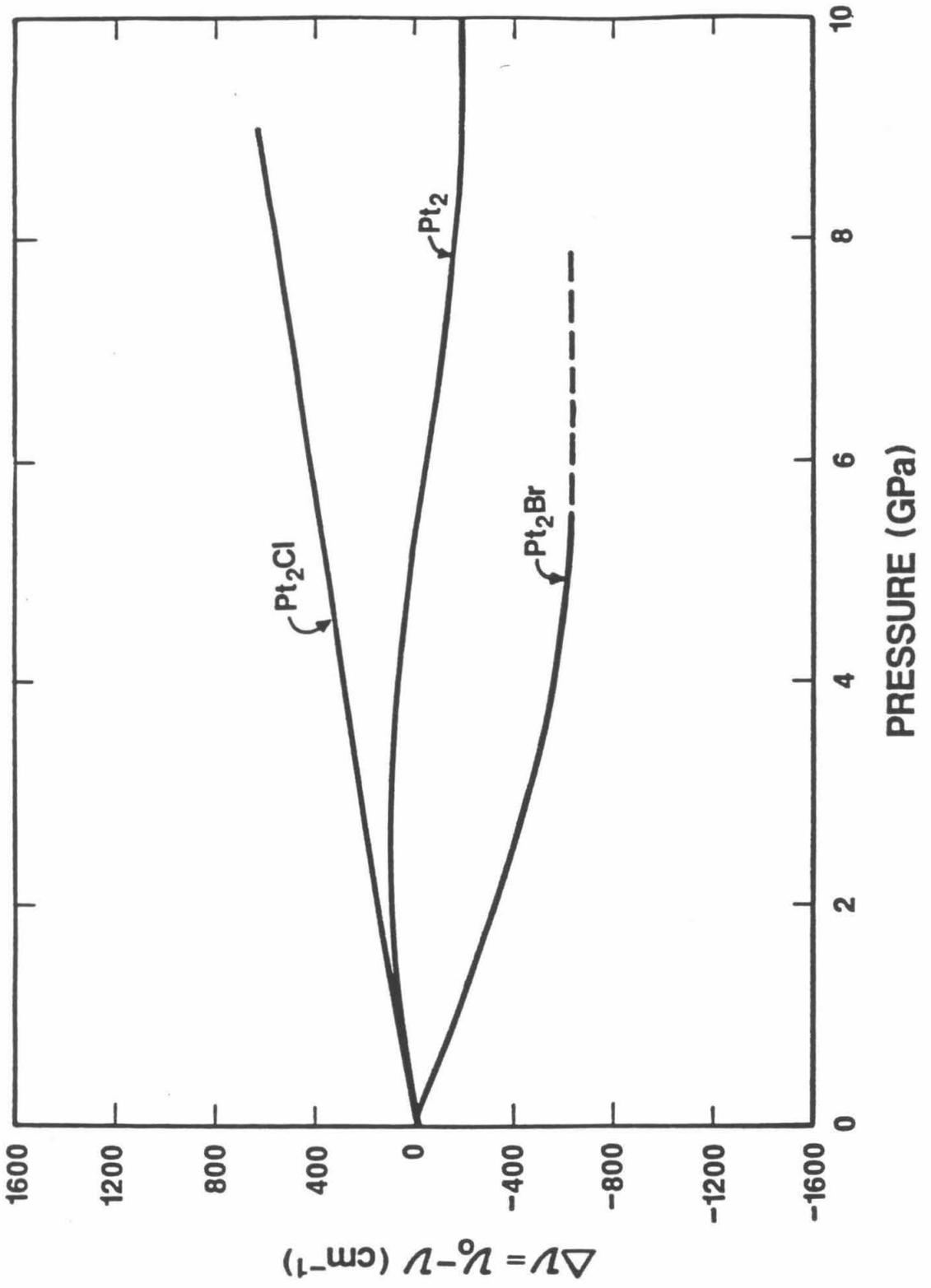


Comparisons of Pt_2Cl_2 , Pt_2 , and Pt_2X Transitions.

A pressure-induced blue shift of the $d_{xz}, d_{yz} \rightarrow d\sigma^*$ transition in Pt_2Cl_2 of approximately $300 \text{ cm}^{-1}/\text{GPa}$ was observed and a shift of $190 \text{ cm}^{-1}/\text{GPa}$ can be inferred for Pt_2 . It is reasonable for the $d\sigma^*$ orbital in Pt_2Cl_2 to be destabilized more rapidly than in Pt_2 because it is coupled to the halogen orbitals. By way of comparison, the $\text{Mn}_2(\text{CO})_{10}$ $d_{xz}, d_{yz} \rightarrow d\sigma^*$ transition blue shifts at an initial rate of $500 \text{ cm}^{-1}/\text{GPa}$ in the solid state.²³ However, the pressure-induced shift of the transition in $\text{Mn}_2(\text{CO})_{10}$ is influenced by stabilization of the $d\pi$ levels because of increased π -backbonding with the carbonyl ligands.

The pressure-induced shifts of the $d\sigma^* \rightarrow p\sigma$ transition in Pt_2 , Pt_2Br , and Pt_2Cl are compared in Figure 4.13. A weak shift was observed in all three complexes; however, the direction of the shift varies, thereby indicating that there is a delicate balance among several competing effects. Increased coupling between the d_{z^2} and p_z orbitals, which may be somewhat larger in the monohalides because of the presence of the bridging halide, blue shifts $d\sigma^* \rightarrow p\sigma$. This appears to be the dominant effect in Pt_2Cl . The observed red shift in Pt_2Br can be explained in terms of band broadening of the $d\sigma^*$ and $p\sigma$ orbitals due to enhanced intermolecular interactions along the z axis. The band broadening should be larger in Pt_2Br than in Pt_2Cl , since Pt_2Br becomes delocalized more rapidly with increasing pressure.²⁴ A red shift of the $d\sigma^* \rightarrow p\sigma$ transition also results from an increase in coupling of $d\sigma^*$ with the phosphorus σ orbitals, an effect that becomes important in Pt_2 at high pressure.

Figure 4.13 Relative energy shifts of the $d\sigma^* \rightarrow p\sigma$ transition in $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}]^{4-}$, $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4]^{4-}$, and $[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Br}]^{4-}$.



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CHAPTER 5

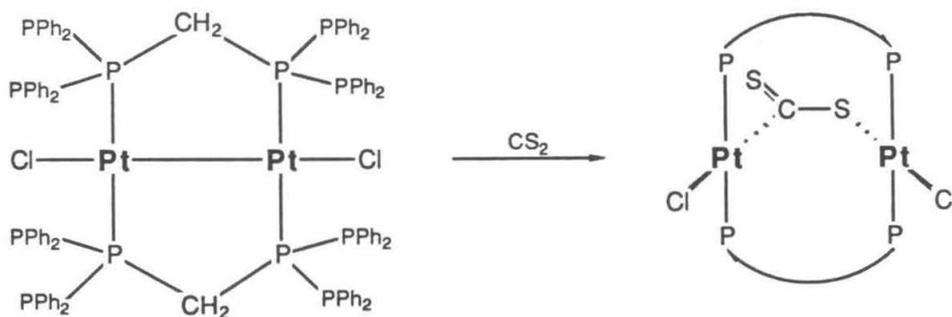
Photophysics of $[\text{Pt}_2(\mu\text{-dppm})_2\text{Y}_2]^{n+}$

(Y = Cl, Br; n = 0 or Y = PPh₃, PPh₂Me, PPhMe₂; n = 2)

INTRODUCTION

Adjacent coordination sites are necessary for concerted, small molecule addition to a photo-excited metal center to occur. $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ (dppm = bis(diphenylphosphino)methane) and its analogs are good candidates for the study of this type of reactivity. The ligand arrangement about each platinum metal center is square planar with each platinum atom acting as the other platinum's fourth ligand.^{1,2} The two bridging dppm ligands are opposite of one another; thus, unlike in the unbridged $d^9\text{-}d^9$ complex $\text{Pt}_2(\text{CNCH}_3)_6^{2+}$ which adopts a staggered structure,³ this molecule adopts a more planar (or eclipsed) structure with an $\sim 40^\circ$ degree twist (see Figure 1.4). Overall, these bridged complexes have four sites available for substrate addition, two adjacent sites above and two below the molecular plane.

These molecules are known to possess a rich thermal chemistry dominated by small molecule addition across the metal-metal bond⁴⁻¹⁰ (see Figure 1.5). For example, reacting $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ with CS_2 results in the formation of an A-frame $d^8\text{-}d^8$ dimer in which the CS_2 group bridges the metal centers through a carbon and a sulfur atom.⁷



Although the thermal reactivity of these complexes have been explored, no investigations into their photochemistry have been reported.

Having the necessary structural requirements for concerted, small molecule addition, it is worthwhile to evaluate the photochemical potential of this system by gaining some understanding of the nature of the excited state and the factors which govern its reactivity.

EXPERIMENTAL

Materials

All solvents used in the synthesis reactions were used as received without further purification. Solvents used in spectroscopic experiments were purified, if necessary, dried, freeze/pump/thaw degassed (4 cycles), then bulb-to-bulb distilled under vacuum into storage flasks (glass solvent pots equipped with teflon vacuum valves) and again freeze/pump/thaw degassed (3 cycles). Dichloromethane (Burdick and Jackson, UV grade) was dried over calcium hydride and stored over freshly activated Linde 3A molecular sieves. The sieves were activated by heating under dynamic vacuum ($<10^{-3}$ Torr) for 24 - 48 hours. Significant purification of 2-methyltetrahydrofuran, 2-MeTHF, (Aldrich, 99%) was necessary. The 2-MeTHF was first distilled from a 0.5% suspension of CuCl_2 followed by distillation from potassium hydroxide pellets. Finally the 2-MeTHF was dried over calcium hydride and stored over sodium metal and benzophenone. Absolute ethanol (USI Chemicals Co., punctillious) was dried over sodium metal and stored over freshly activated 3A Linde molecular sieves. Methanol (Burdick and Jackson, UV grade) was stored over freshly activated Linde 3A molecular sieves.

$\text{Pt}(\text{dppm})\text{Cl}_2$: K_2PtCl_4 (Aldrich, 98%) and bis(diphenylphosphino)methane, dppm, (Aldrich, 97%) were used to prepare $\text{Pt}(\text{dppm})\text{Cl}_2$ by a method similar to that used by Westland to prepare $\text{Pt}(\text{dppe})\text{Cl}_2$ (dppe = 1,2-bis(diphenylphosphino)ethane).¹¹ K_2PtCl_4 (0.50 g) and dppm (0.45 g) were refluxed in a mixture of 17 mL ethanol and 17 mL concentrated HCl for 4 1/2 hours. During the course of the reaction the desired product precipitated from the hot mixture as a fine powder. The reaction was completed when the initial red color of K_2PtCl_4 disappeared

and only the white powder product remained. The product was clean and further purification was not necessary. (If desired, the complex may be recrystallized from hot dimethylformamide, DMF, by adding a mixture of petroleum ether and diethyl ether.) The yield is essentially 100%. [Analysis calculated for $\text{Pt}_2(\text{dppm})\text{Cl}_2$ (%): C, 46.17; H, 3.41. Found: C, 45.99; H, 3.43.]

$\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$: The procedure of Brown *et al.*¹² was followed with slight modification. $\text{Pt}(\text{dppm})\text{Cl}_2$ (1.2 g) was suspended in 20 mL of nitrogen purged ethanol in a round bottom flask equipped with a stir bar. Sodium borohydride, NaBH_4 , (0.5 g) was suspended in 15 mL of nitrogen purged ethanol in an addition funnel. Over a 40 - 60 minute period, the NaBH_4 mixture was added dropwise to the $\text{Pt}(\text{dppm})\text{Cl}_2$ mixture with rapid stirring; the white mixture began to take on a yellow-brown tinge during the addition. After all the NaBH_4 was added, the mixture was stirred for another 5 minutes, then filtered. The solid was rinsed with methanol and water to get rid of any unreacted NaBH_4 . The impure $[\text{Pt}_2(\mu\text{-dppm})_2(\mu\text{-H})(\text{H})_2]\text{Cl}$ solid was prepared for the next reaction step by further rinsing with ethanol, diethyl ether, and lastly benzene. (On occasion, the impure $[\text{Pt}_2(\mu\text{-dppm})_2(\mu\text{-H})(\text{H})_2]\text{Cl}$ was dried *in vacuo* for 1 hour as suggested in the original procedure.)¹² Next, impure $[\text{Pt}_2(\mu\text{-dppm})_2(\mu\text{-H})(\text{H})_2]\text{Cl}$ (1.0 g) was suspended in 45 mL of benzene and 1.5 mL concentrated HCl and refluxed for 15 minutes. The hot yellow solution was decanted off from the red-brown residue, allowed to cool, and precipitated by adding ~15 mL iso-octane. The final product recrystallized as the lemon-yellow benzene solvate $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{C}_6\text{H}_6$ in ~70% yield. In most instances, the benzene solvate obtained directly from the reaction was pure; however if further cleaning was necessary,

there were several methods of purification depending upon which solvate was desired. Compounds studied gave acceptable elemental analysis and clean NMR spectra which matched published results.

$\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ dissolves in only a few solvents, mainly halocarbons. Recrystallizing $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ from the solvent of choice and precipitating with iso-octane will yield crystals of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ with the corresponding solvate. $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{C}_6\text{H}_6$ is not soluble in benzene. If the sample obtained from the reaction mixture needs to be cleaned and the benzene solvate is desired, any impurities present must be washed from the solid. The easiest way to remove the impurities is to stir a suspension of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{C}_6\text{H}_6$ in boiling benzene for ~15 minutes, then filter. Repeat this procedure as often as necessary to obtain pure product.

Occasionally during the course of recrystallization, white crystalline specks were observed in addition to the desired yellow $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{solvate}$ crystals. The white impurity is believed to be $[\text{Pt}_2(\mu\text{-dppm})_2(\mu\text{-H})\text{Cl}_2]\text{Cl}$. Addition of a few drops of triethylamine, NEt_3 , clears up this contamination problem.

$\text{Pt}(\mu\text{-dppm})_2\text{Br}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$, $[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_3)_2](\text{PF}_6)_2$, $[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_2\text{Me})_2](\text{PF}_6)_2$, $[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPhMe}_2)_2](\text{PF}_6)_2$: These platinum complexes were prepared by the published procedures.^{12,13}

$\text{Pd}_2(\mu\text{-dppm})_2\text{Cl}_2$: This palladium complex was prepared by method 2 of the published procedure.⁸

Physical Measurements

Electronic Absorption Spectroscopy: Electronic absorption spectra were measured using a Cary Model 17 or Shimadzu UV-260 spectrophoto-

meter. Low temperature, 77 K, electronic absorption spectra were measured using the Cary Model 17 spectrophotometer equipped with a quartz optical dewar (copper block base). Samples for the absorption measurements were prepared by placing $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ (~1 mg) into an extra long spectrophotometric cell consisting of a square cuvette (1 cm pathlength) and a teflon vacuum valve.

Emission and Excitation Spectroscopy: Emission spectra were recorded using an emission spectrophotometer constructed at Caltech and described elsewhere.¹⁴ Samples for these measurements (solids and solutions) were placed in NMR tubes equipped with a 14/20 ground glass joint and capped with a septum. Samples were run in aerated, N_2 purged, and degassed solutions; no significant differences were observed. Emission and excitation spectra were measured using a Perkin-Elmer spectrofluorimeter MPF-66 interfaced to a Perkin-Elmer computer. Emission and excitation spectra measured at 77 K were obtained by placing the sample in an NMR tube and immersing the tube in a liquid-nitrogen-filled quartz finger dewar. Samples for these measurements were placed in 3 mm (o.d.) quartz EPR tubes capped with a septum and parafilm. Lower temperature emission and excitation spectra (down to 16 K) were measured by placing the solid sample (adhered on a quartz disk with silicon grease) into a CTI Model 21 cryocooler. The temperature inside the cryocooler was varied with a small resistive heater; a chromel/Au-Fe thermocouple was used to measure the temperature.

Emission Lifetime Measurements: Luminescent lifetime data were collected using a pulsed Nd:YAG laser system described previously.¹⁵ The third harmonic (355 nm) was used as the excitation line. Low

temperature, 77 K experiments were carried out using the before mentioned liquid-nitrogen-filled quartz finger dewar. Samples for these measurements (solids and solutions) were placed in NMR tubes equipped with a 14/20 ground glass joint and capped with a septum. Samples were run in aerated, N_2 purged, and degassed solutions; no significant differences were observed. Variable temperature experiments were performed using a continuous-flow nitrogen gas dewar in which the nitrogen gas was cooled by travelling through copper tubing immersed in a liquid-nitrogen-filled dewar. The temperature was regulated by the speed of gas flow and measured by a copper vs. copper nickel (copper-constantan) thermocouple. Variable temperature samples were placed in extra long 5 mm (o.d.) quartz EPR tubes and capped with a septum and parafilm.

Other. Elemental analyses were performed by Mr. Larry Henling and Mr. Fenton Harvey at the Caltech Analytical Laboratory, Pasadena, California.

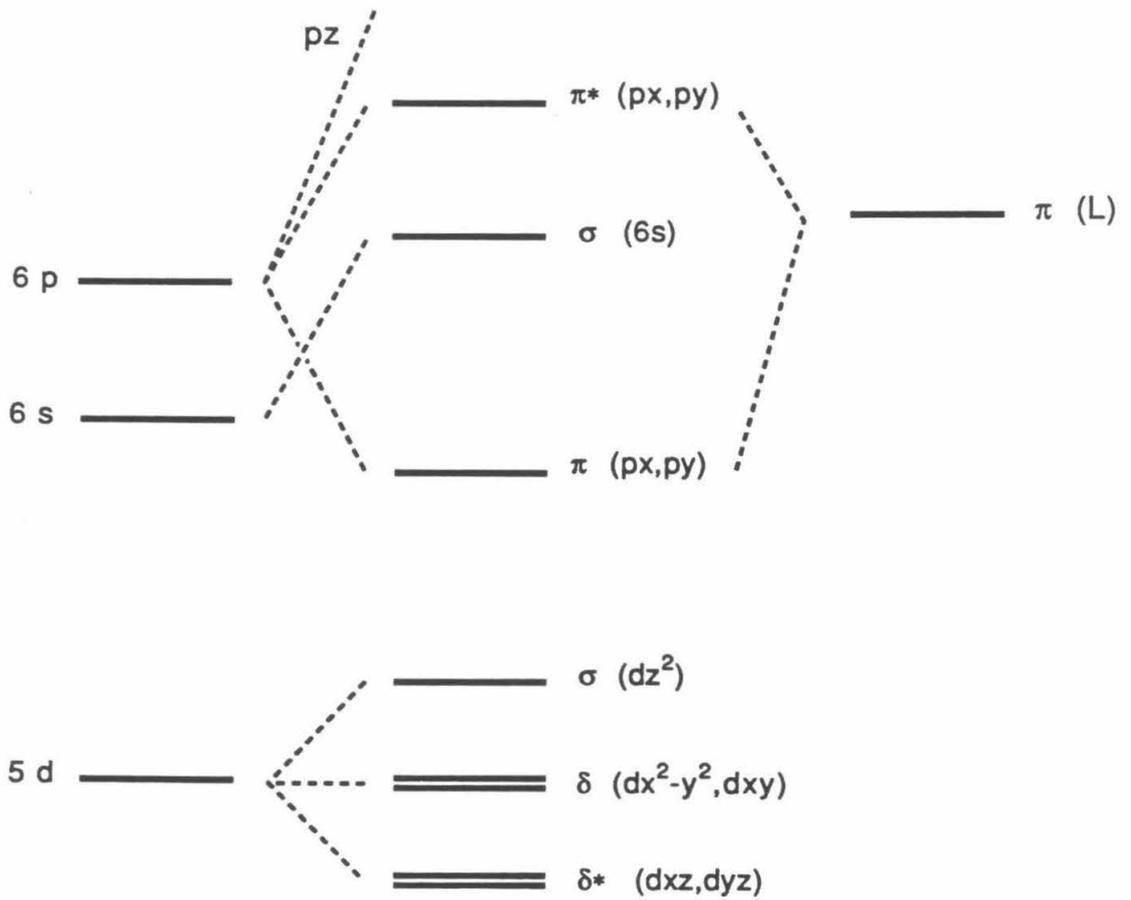
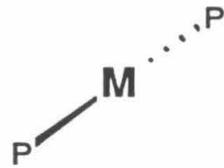
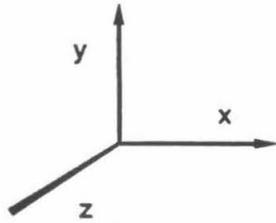
RESULTS and DISCUSSION

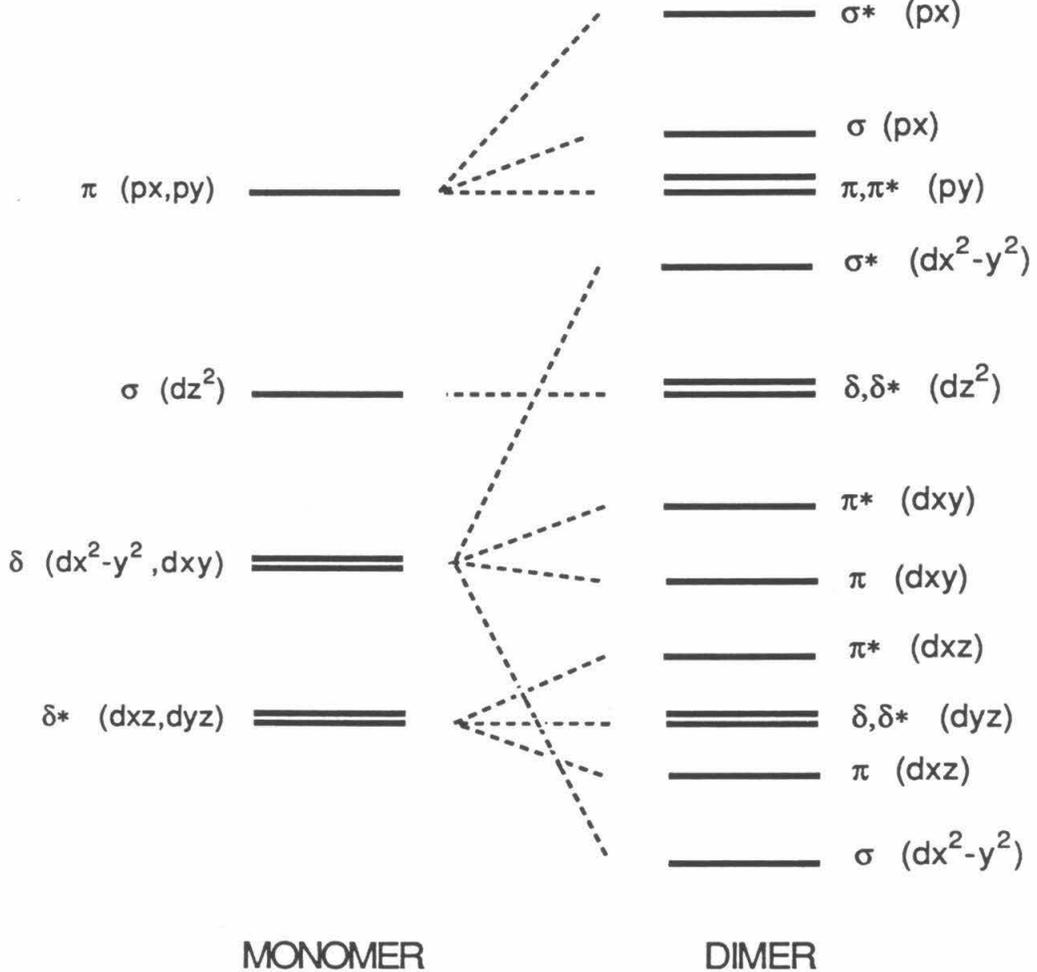
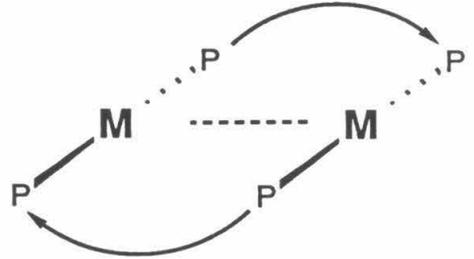
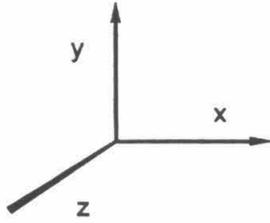
In order to discuss the origin of the chemical reactivity of these d^9-d^9 dimers, a description of their electronic distribution would prove useful; however no detailed molecular orbital (MO) diagram for these systems has been proposed. There may be several reasons for this, one being that the electronic absorption spectra of these complexes are quite complicated and the bands are not well resolved, making the measurement of exact band maxima, peak intensities, and band widths difficult.

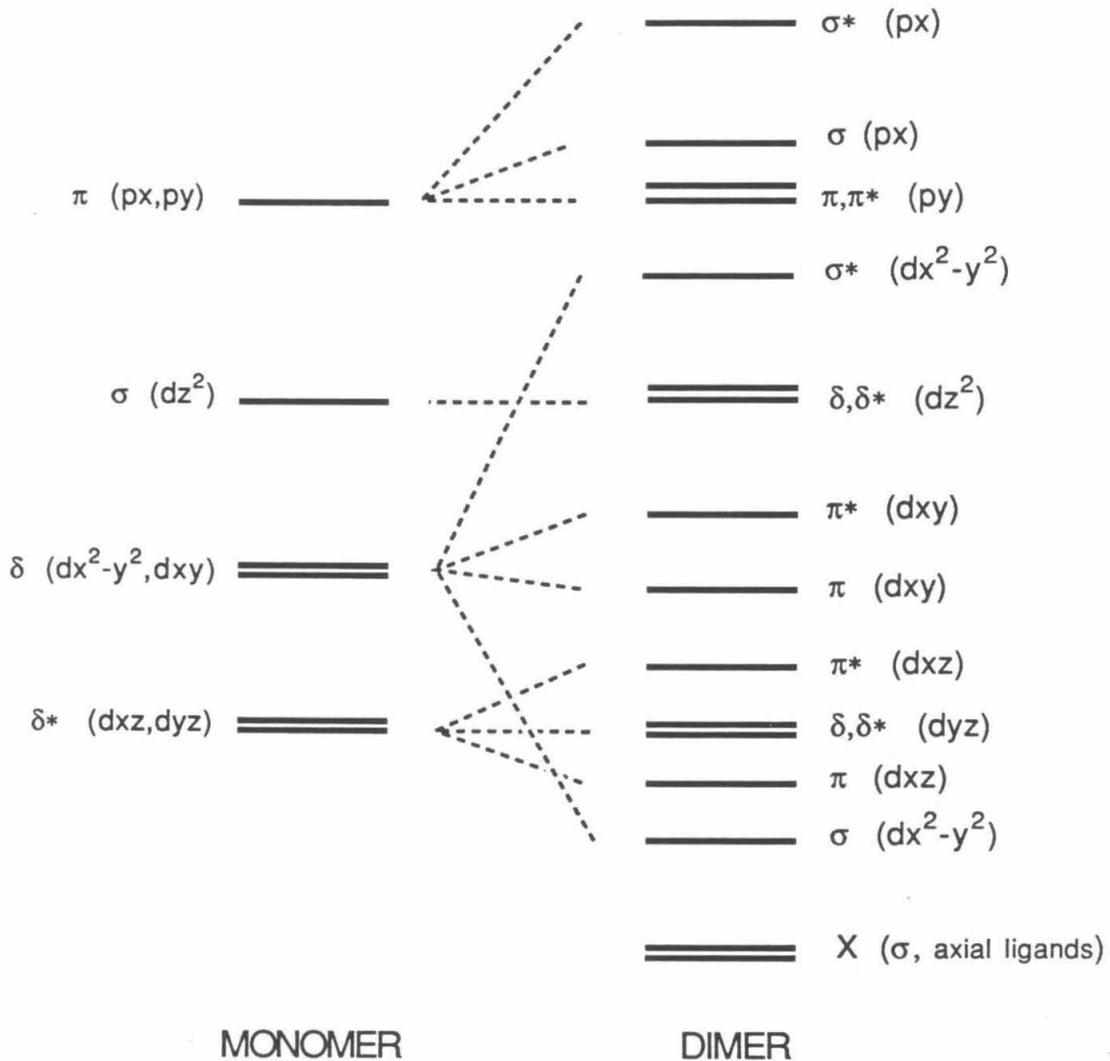
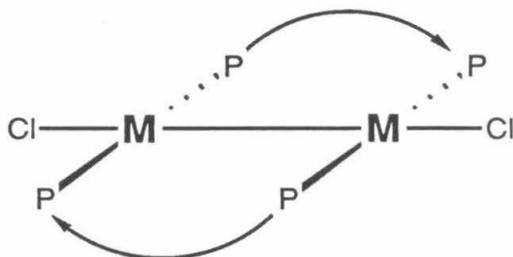
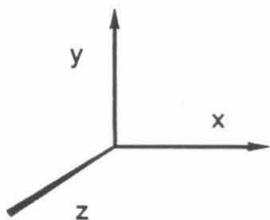
A plausible MO diagram may be constructed from MO diagrams of monomeric d^{10} complexes which possess similar geometries and ligand environments. The MO diagram for $[\text{Au}(\text{PEt}_3)_2]^+$ is shown in Figure 5.1a.¹⁶ Bringing two such diagrams together along the x-axis (designated as the metal-metal bonding axis) results in the diagram illustrated in Figure 5.1b. When axial ligands are added to $\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4^{4-}$, the resulting d^7-d^7 MO diagram has the same basic orbital ordering as the d^8-d^8 MO diagram except with two additional $\sigma(\text{X})$ orbitals and two fewer metal localized electrons (see Figures 4.1 and 4.4). Therefore, addition of axial halides to the $d^{10}-d^{10}$ MO diagram constructed in Figure 5.1b, along the x axis, will most likely change the energies of several of the levels, without perturbing the orbital ordering very much. The MO diagram presented in Figure 5.1c takes these perturbations into account.

From Figure 5.1c, the highest occupied molecular orbital, HOMO, is expected to be $\delta^*(d_{z^2})$. However, depending upon how close the monomer d_{z^2} and d_{xy} , $d_{x^2-y^2}$ levels are, the possibility that $\pi^*(d_{xy})$ is the HOMO cannot be completely ruled out.

- Figure 5.1**
- (a) Molecular orbital diagram for MP_2 (taken from reference 16).
 - (b) Proposed molecular orbital diagram for $(\text{MP}_2)_2$.
 - (c) Proposed molecular orbital diagram for $(\text{MP}_2)_2\text{X}_2$.







The lowest unoccupied molecular orbital, LUMO, is predicted to be $\sigma^*(d_{x^2-y^2})$; however, the $\pi(py)$ orbital cannot be immediately excluded from consideration. In $d^{10}-d^{10}$ systems the HOMO level is $d\sigma^*(d_{x^2-y^2})$ and the LUMO is p_x, p_y ; however addition of the axial halide along the x axis in this d^9-d^9 systems may raise the $d\sigma^*(d_{x^2-y^2})$ level up sufficiently, making the $\pi(p_y)$ level lower in energy. On the other hand, photochemically induced oxidative-addition reactions of unbridged d^9-d^9 systems¹⁷⁻¹⁹ show metal-metal bond cleavage and production of Pt(II) $15e^-$ radical species as the predominant reaction intermediate. Similarly, in bridged systems, where the dissociative pathway is unlikely, substrate addition across the metal-metal bond occurs with net metal-metal bond cleavage.⁴⁻⁶ Therefore, the assignment of the LUMO as $d\sigma^*(d_{x^2-y^2})$ is given additional support considering the observed thermal and photochemical reactivity.

Group theory analysis of the proposed MO diagram predicts several orbitally allowed electronic transitions which are listed in Table 5-1. The first six transitions are most likely to be the lowest energy transition assuming that the $\sigma^*(d_{x^2-y^2})$ and $\pi\pi^*(p_y)$ orbitals are close in energy.

Crystals of the platinum(I) complexes are yellow in color; whereas, the palladium(I) analogs are more intensely orange colored. The room temperature electronic absorption spectrum of $Pt_2(\mu-dppm)_2Cl_2$ in dichloromethane first reported by Alves *et al.*²⁰ is shown in Figure 5.2. Three bands are observed; however, they are not well resolved and do not allow for any theoretical analysis. Electronic absorption band maxima of $Pd_2(\mu-dppm)_2Cl_2$ have been reported by Benner *et al.*;⁸ the

Figure 5.2 Electronic absorption spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ in dichloromethane (300 K).

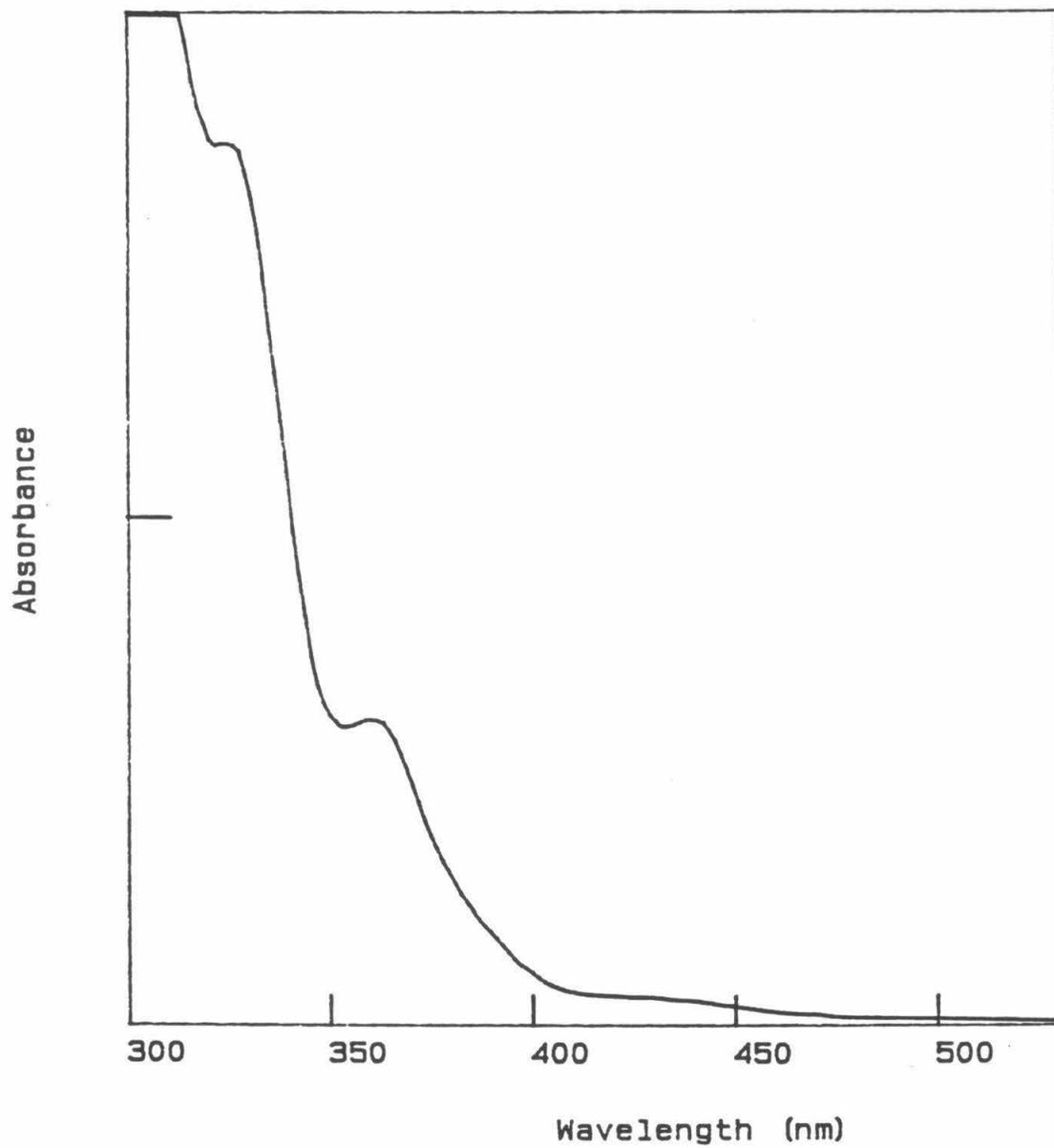
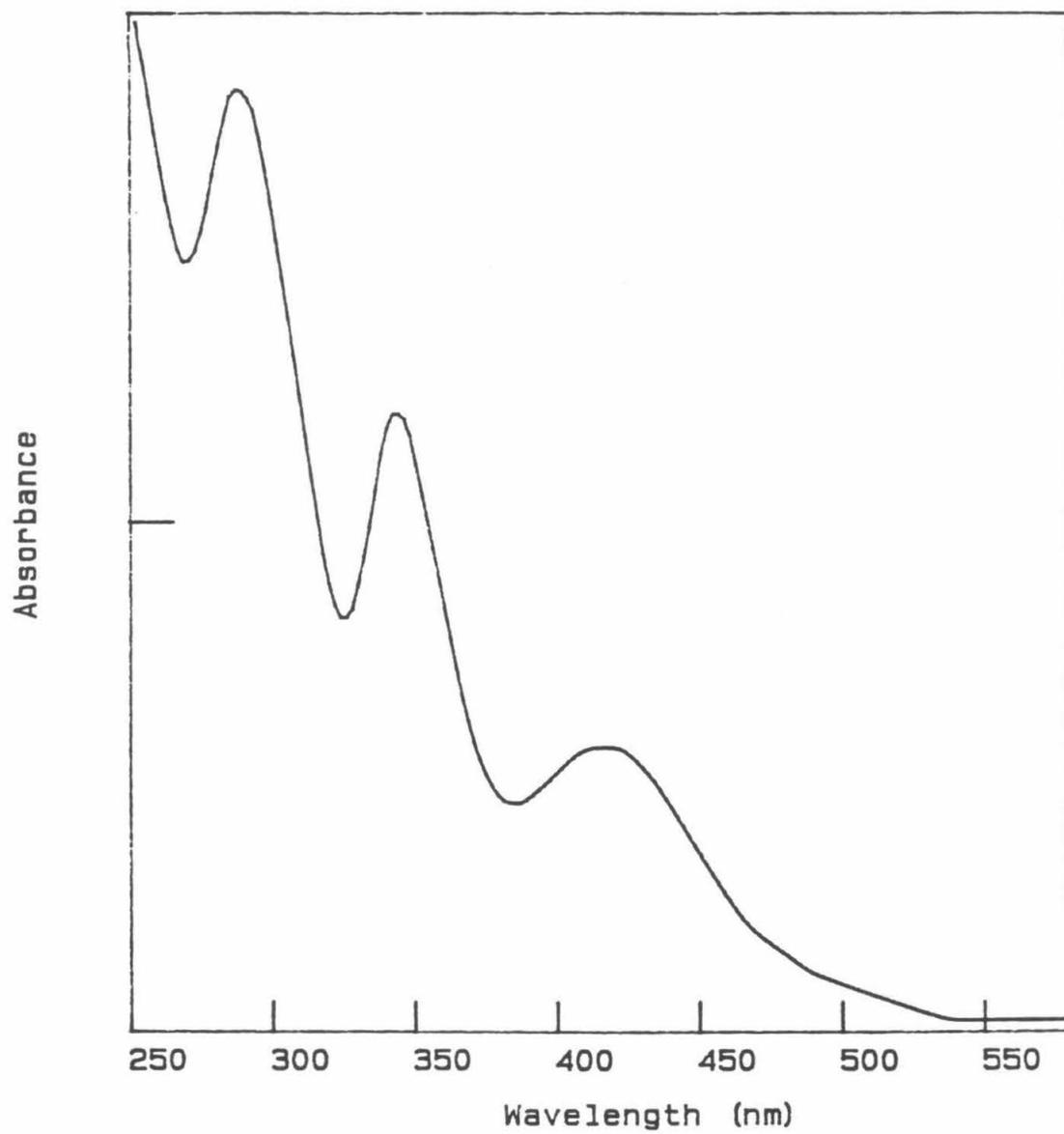


Figure 5.3 Electronic absorption spectrum of $\text{Pd}_2(\mu\text{-dppm})_2\text{Cl}_2$ in dichloromethane (300 K).

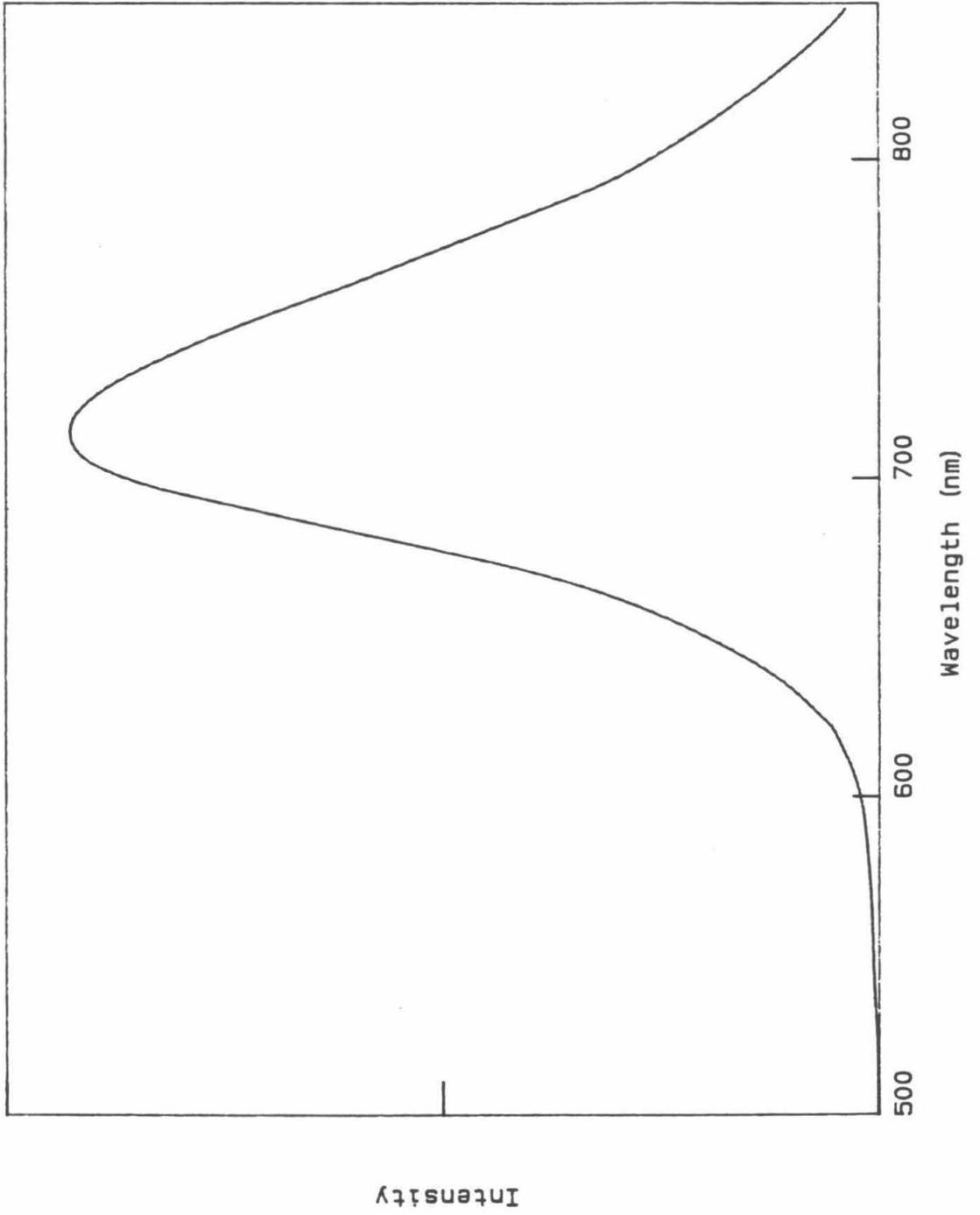


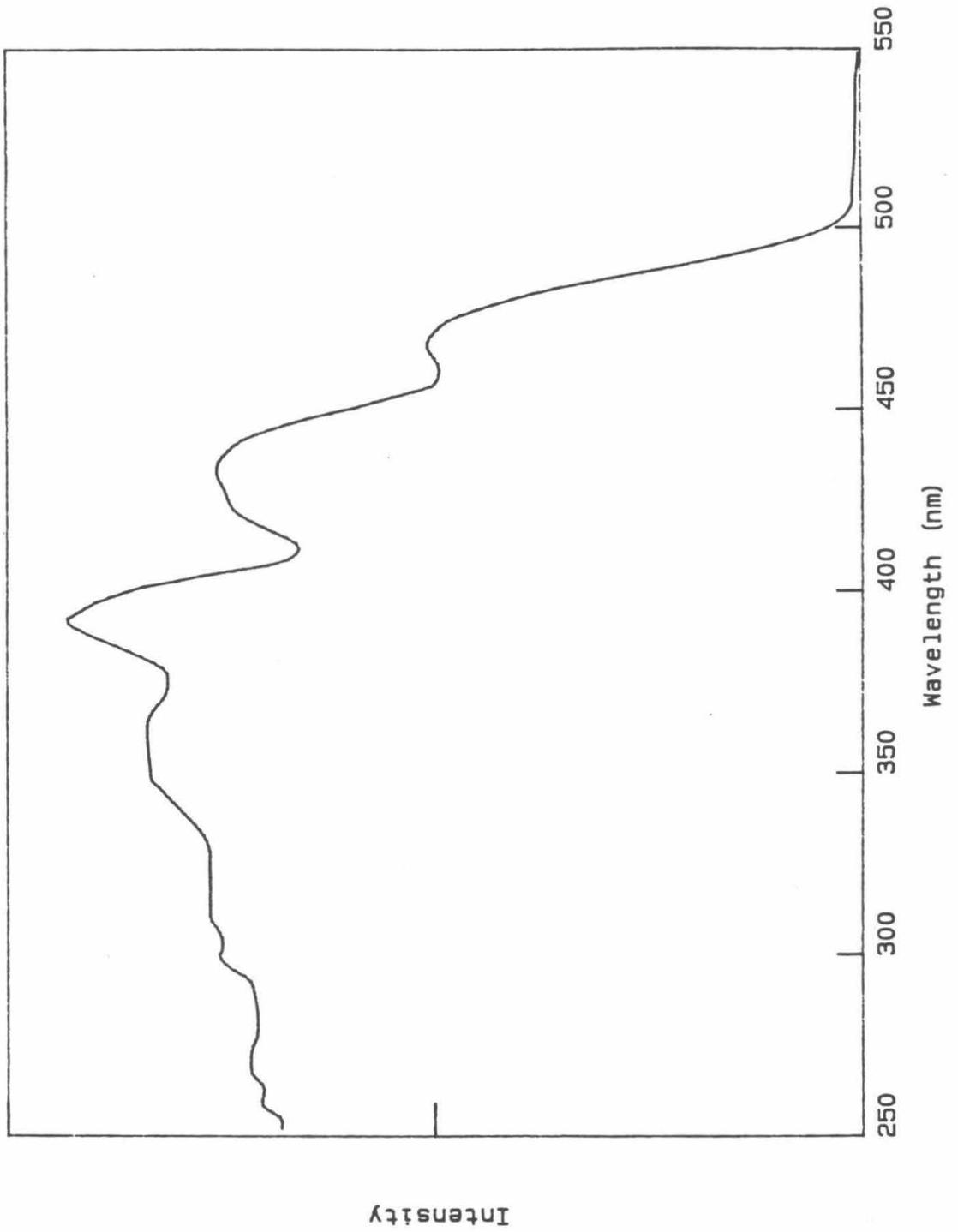
actual spectrum is shown in Figure 5.3. In this spectrum two absorption bands at lower energy are fairly well resolved with a shoulder at higher energy. Surprisingly, the absorption band shapes and extinction coefficients for the platinum(I) and palladium(I) dimers do not look as similar as might be expected for such closely related complexes. More than likely the observed absorption bands do not correspond to the same transitions in the platinum and palladium systems.

Luminescence from the platinum(I) dimers is observed only in the solid state and frozen solution. Similar luminescence is not observed in the analogous palladium system. It was necessary to confirm that the d^9-d^9 molecule present in solution was the species responsible for the observed luminescence, since highly luminescent d^8-d^8 and $d^{10}-d^{10}$ platinum complexes are known and may have been present in trace quantities. Comparison of the excitation spectra of $Pt_2(\mu-dppm)_2Cl_2 \cdot 0.5CH_2Cl_2$ and $Pt_2(\mu-dppm)_2Br_2 \cdot 0.5CH_2Cl_2$ with their absorption spectra verified that the observed luminescence was definitely arising from these d^9-d^9 dimers. (The low temperature solid emission and excitation spectra for $Pt_2(\mu-dppm)_2Cl_2 \cdot 0.5CH_2Cl_2$ and $Pt_2(\mu-dppm)_2Br_2 \cdot 0.5CH_2Cl_2$ are shown in Figures 5.4 and 5.5, respectively.)

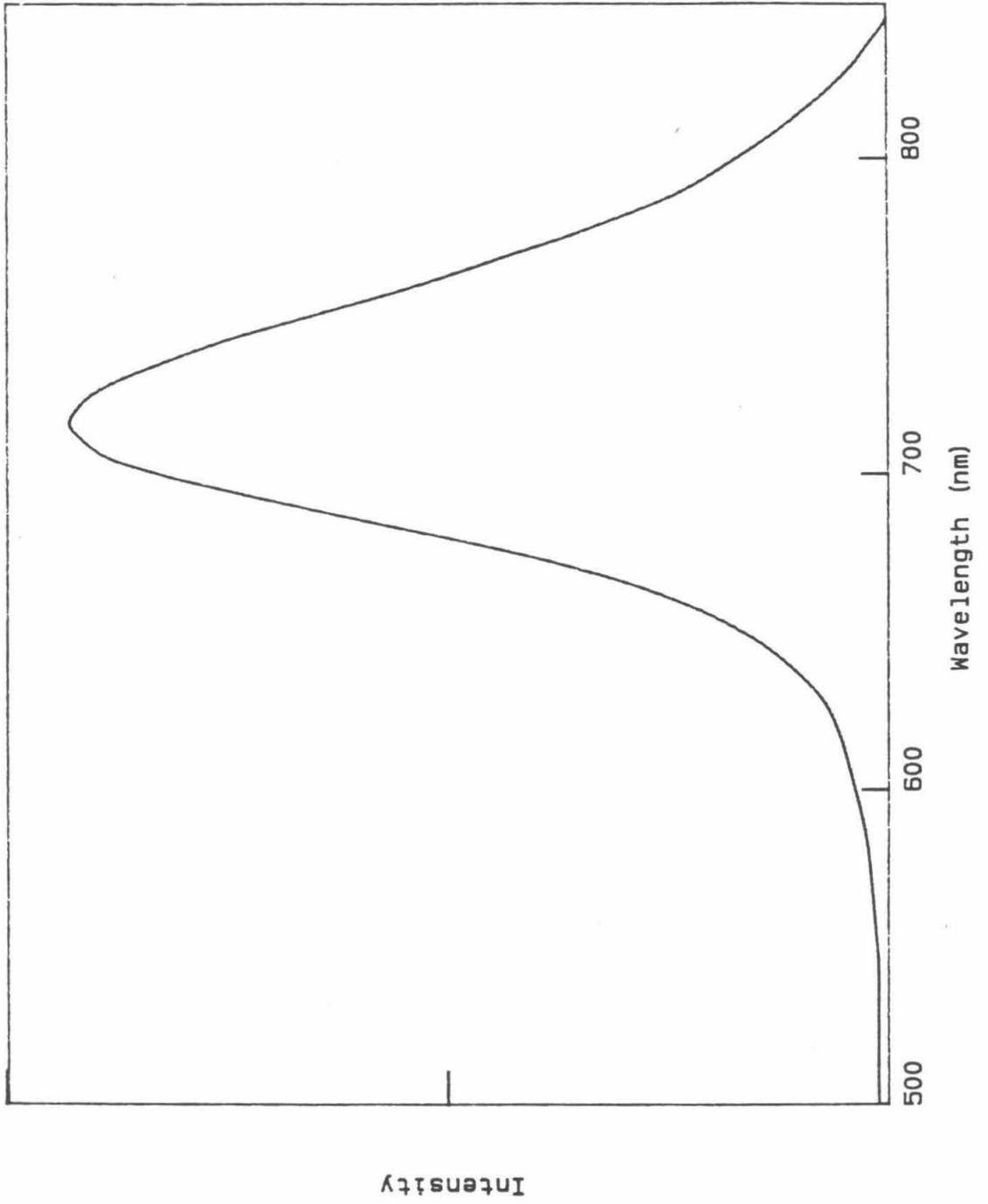
The low temperature solid excitation spectrum of $Pt_2(\mu-dppm)_2Cl_2$ revealed the presence of additional electronic transitions at ~ 380 and ~ 465 nm which were either completely masked or too weak to detect in solution. This observation prompted further investigation into the electronic absorption spectrum of $Pt_2(\mu-dppm)_2Cl_2$. At low temperature, 77 K, in a 1:1 CH_2Cl_2 :2-MeTHF glass (Figure 5.6), the absorption band at 380 nm begins to be resolved, but the absorption band at ~ 465 nm is

- Figure 5.4**
- (a) Corrected emission spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$, powder sample, 21 K. Emission maximum at 714 nm.
 - (b) Corrected excitation spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$, powder sample, 21 K. Absorption maxima at 465, 429, and 389 nm.





- Figure 5.5**
- (a) Corrected emission spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$, powder sample, 16 K. Emission maximum at 716 nm.
 - (b) Corrected excitation spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$, powder sample, 16 K. Absorption maxima at 478, 441, 390 nm.



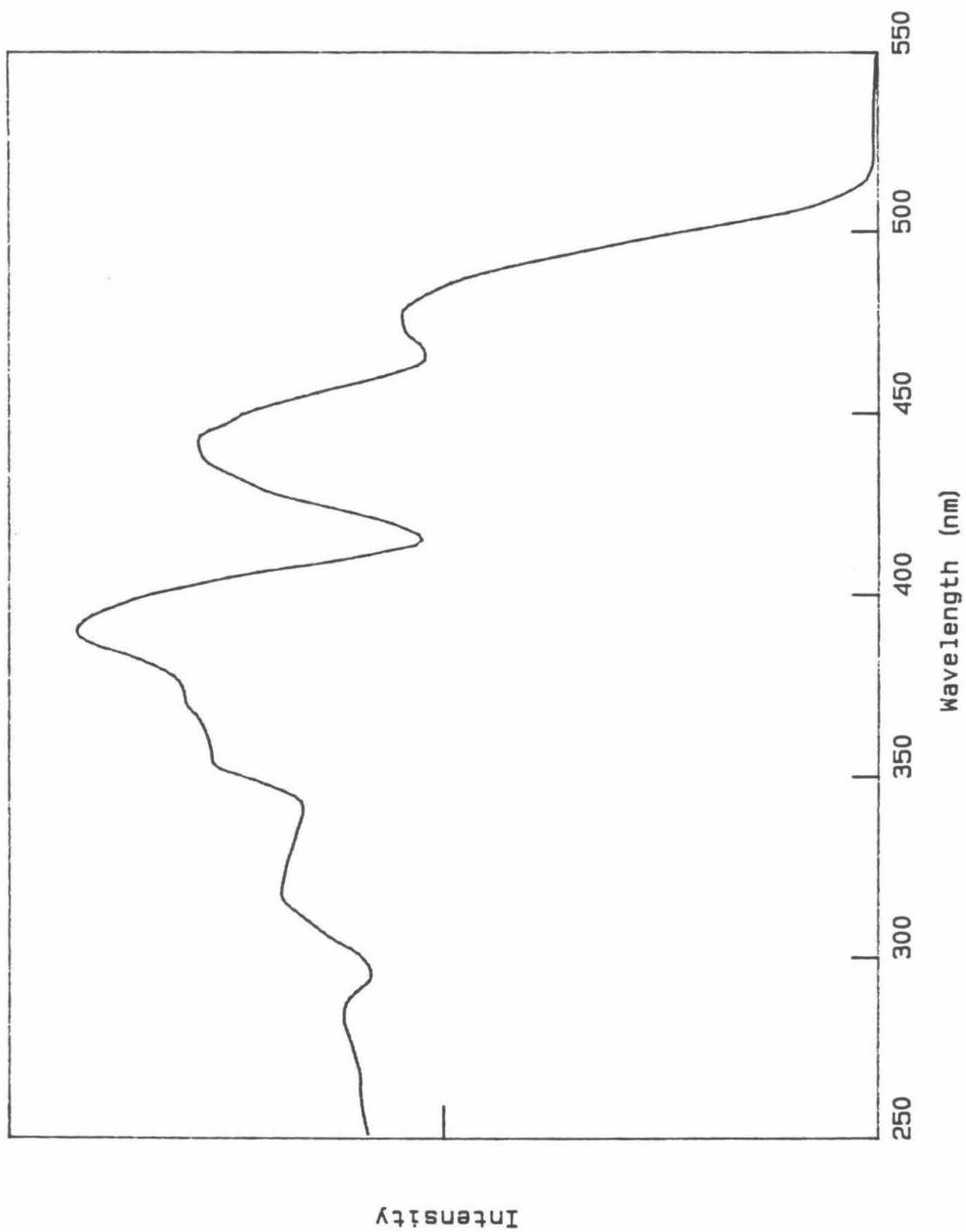
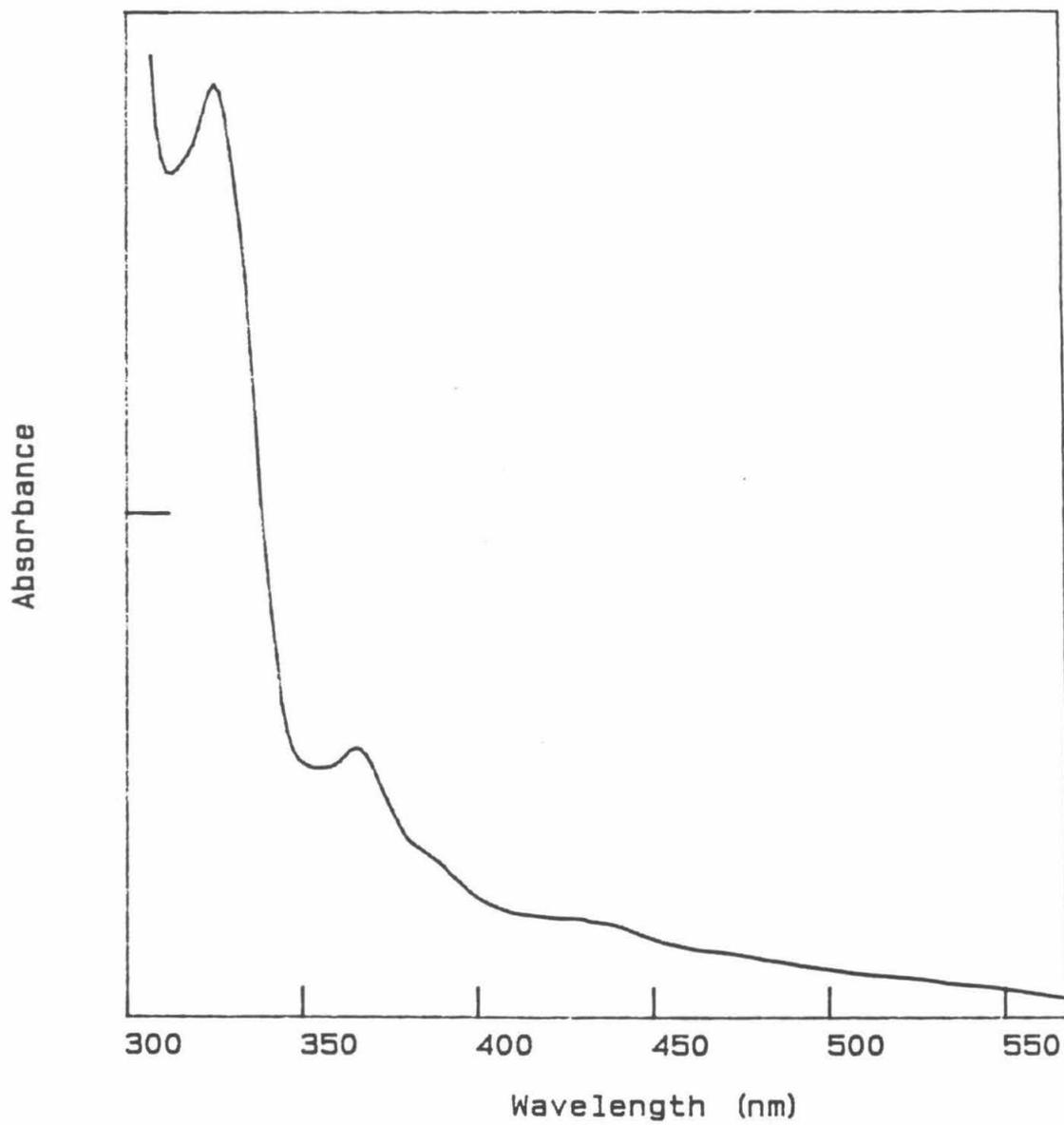


Figure 5.6 Electronic absorption spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ in a 1:1 dichloromethane:2-methyltetrahydrofuran glass (v/v) at 77 K.



still too weak to be detected. It appears that single crystal electronic absorption measurements will be required to observe the absorption bands at wavelengths longer than 450 nm. Thus, at present, it seems reasonable to estimate that the 465 nm absorption, observed in the solid state excitation spectrum, of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ is much less than $100 \text{ M}^{-1} \text{ cm}^{-1}$. (A similar argument may be made for the absorption band at 478 nm observed in the solid state excitation spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2$.)

The extinction coefficients of the two low energy absorption bands of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ at 427 nm ($\epsilon = 200 \text{ M}^{-1} \text{ cm}^{-1}$)²⁰ and 465 nm ($\epsilon < 100 \text{ M}^{-1} \text{ cm}^{-1}$) indicate that these transitions are most likely spin forbidden, orbitally allowed transitions such as $^3[\delta^*(d_{z^2}) \rightarrow \pi^*(py)]$ or $^3[\delta(d_{z^2}) \rightarrow \sigma^*(d_{x^2-y^2})]$ or spin allowed, orbitally forbidden transitions such as $^1[\delta^*(d_{z^2}) \rightarrow \sigma^*(d_{x^2-y^2})]$. The transitions lying to higher energy ($\lambda > 380 \text{ nm}$) all have extinction coefficients greater than $1000 \text{ M}^{-1} \text{ cm}^{-1}$ and can be assigned to a variety of spin and orbitally allowed transitions (see Table 5.1).

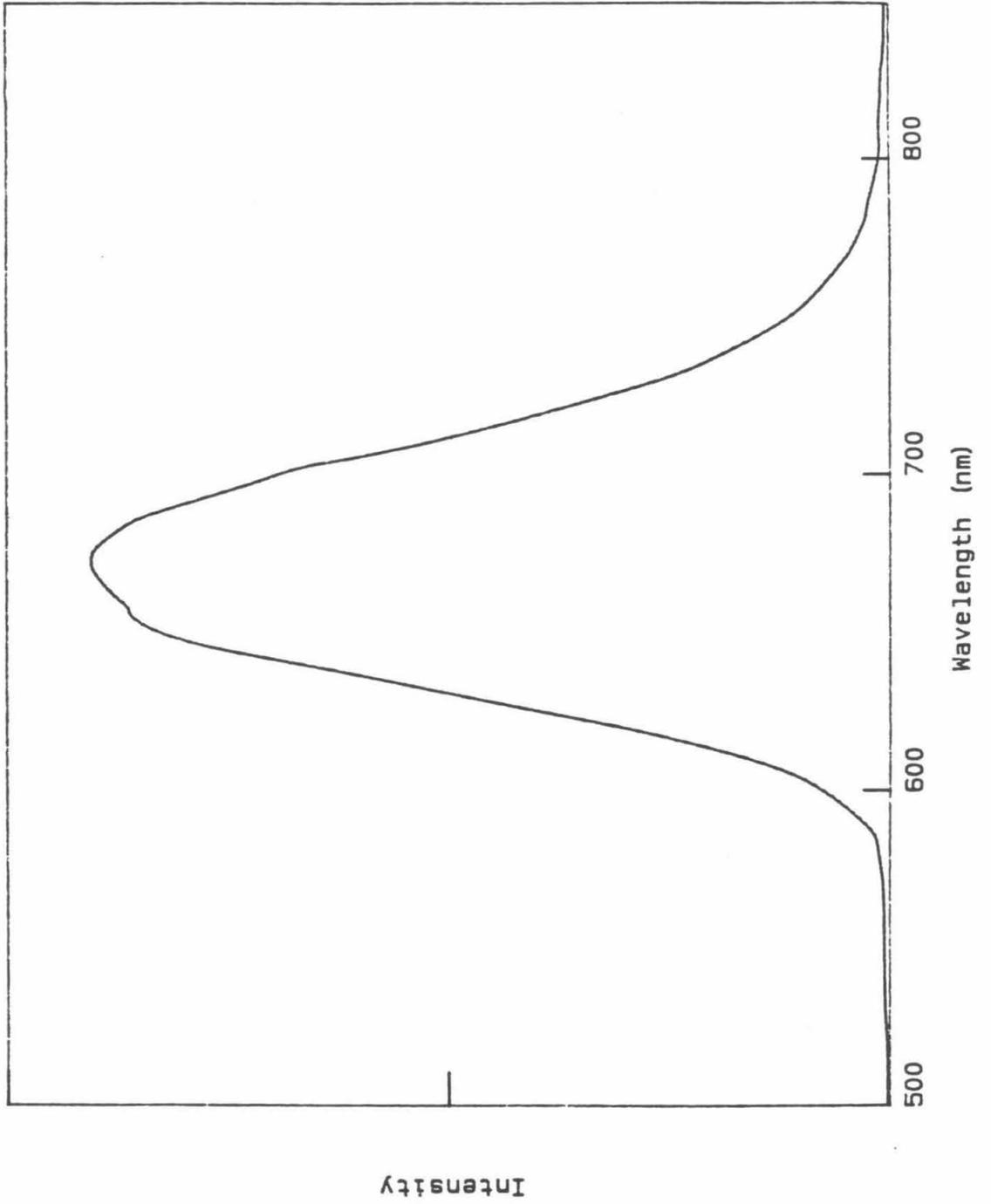
Surprisingly, the luminescence intensity and band maxima of solid samples were found to be dependent upon which solvate was present in the unit cell. Figure 5.7 shows the emission and excitation spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{C}_6\text{H}_6$. The luminescence intensity is much greater and the emission band maxima is significantly shifted to higher energy for the benzene solvate as compared to the dichloromethane solvate. On the other hand, the excitation spectrum for $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{C}_6\text{H}_6$ has nearly identical peak maxima to the excitation spectrum for $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$.

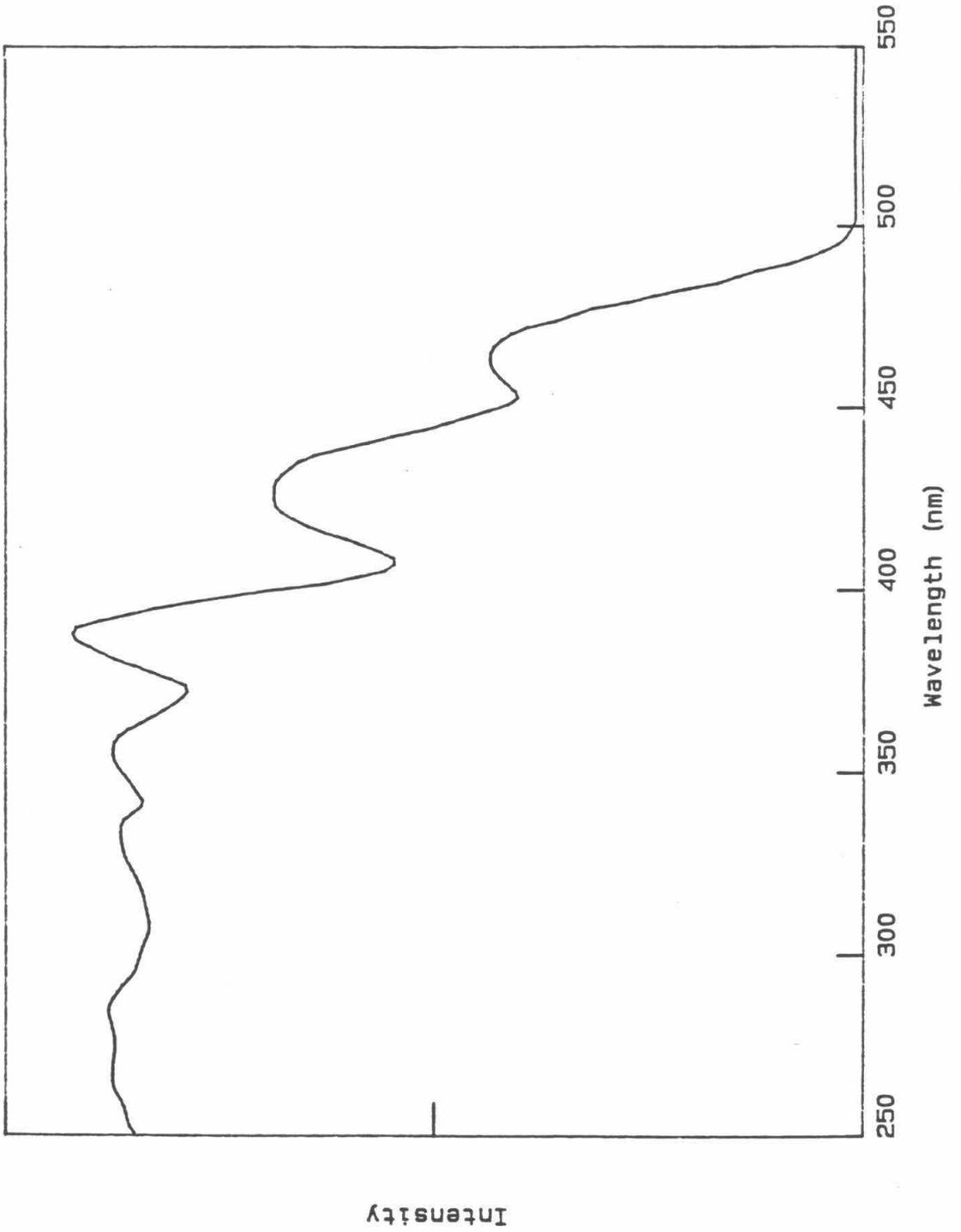
Table 5.1

Orbitally Allowed Transitions assuming D_{2h} symmetry

$\delta^* (d_{z^2})$	\rightarrow	$\pi^* (p_y)$	y allowed
$\delta (d_{z^2})$	\rightarrow	$\sigma^* (d_{x^2-y^2})$	x allowed
$\delta (d_{z^2})$	\rightarrow	$\pi (p_y)$	y allowed
$\pi^* (d_{xy})$	\rightarrow	$\sigma^* (d_{x^2-y^2})$	y allowed
$\pi^* (d_{xy})$	\rightarrow	$\pi (p_y)$	x allowed
$\pi (d_{xy})$	\rightarrow	$\pi^* (p_y)$	x allowed
$\pi^* (d_{xz})$	\rightarrow	$\sigma^* (d_{x^2-y^2})$	z allowed
$\delta^* (d_{yz})$	\rightarrow	$\pi^* (p_y)$	z allowed
$\delta (d_{yz})$	\rightarrow	$\pi (p_y)$	z allowed
$\sigma (d_{x^2-y^2})$	\rightarrow	$\sigma^* (d_{x^2-y^2})$	x allowed
$\sigma (d_{x^2-y^2})$	\rightarrow	$\pi (p_y)$	y allowed

- Figure 5.7**
- (a) Corrected emission spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{C}_6\text{H}_6$, powder sample, 20 K. Emission maximum at 672 nm.
 - (b) Corrected excitation spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{C}_6\text{H}_6$, powder sample, 20 K. Absorption maxima at 465, 428, 388 nm.





Differences in the solid state emission band maxima and lifetimes of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ complexes are listed in Table 5.2. The longest lifetime obtained is that of the benzene solvate, 900 ns; and the shortest lifetime is observed for the dichloromethane solvate, 130 ns.

The general trend of high-energy-emission/long-lifetime and low-energy-emission/short-lifetime is in accord with the energy gap law which relates the non-radiative rate out of an excited state with the energy of the excited state within a given chromophore.²¹ This law is based on the extent of vibrational coupling between the zeroeth vibrational level of the excited state, and the higher lying vibrational levels of the ground state.²² In the case of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{solvate}$, it appears that the solvate is somehow affecting the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ excited state resulting in an emission-energy/emission-lifetime correlation similar to that predicted by the energy gap law.

These results indicate that the solvate molecule interacts with the excited state in some fashion. One possibility is that this interaction is dependent upon the polarity of the solvate molecule, such that a more polar solvent decreases the emission energy and lifetime. Another possibility is that the π system of the solvate interacts with the unfilled σ^* ($d_{x^2-y^2}$) and $\pi\pi^*$ (p_y) orbitals of the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ molecule resulting in an increase of the excited state emission energy and lifetime. Unfortunately, the investigation of these suggestions is not straightforward since these complexes do not luminesce in fluid solution.

Additional evidence for solvate interaction with the metal complex is obtained from glass emission and excitation spectra of $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$ in a 1:1 dichloromethane/2-methyltetrahydrofuran

Table 5.2

Emission Maxima and Lifetimes of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{solvate}$

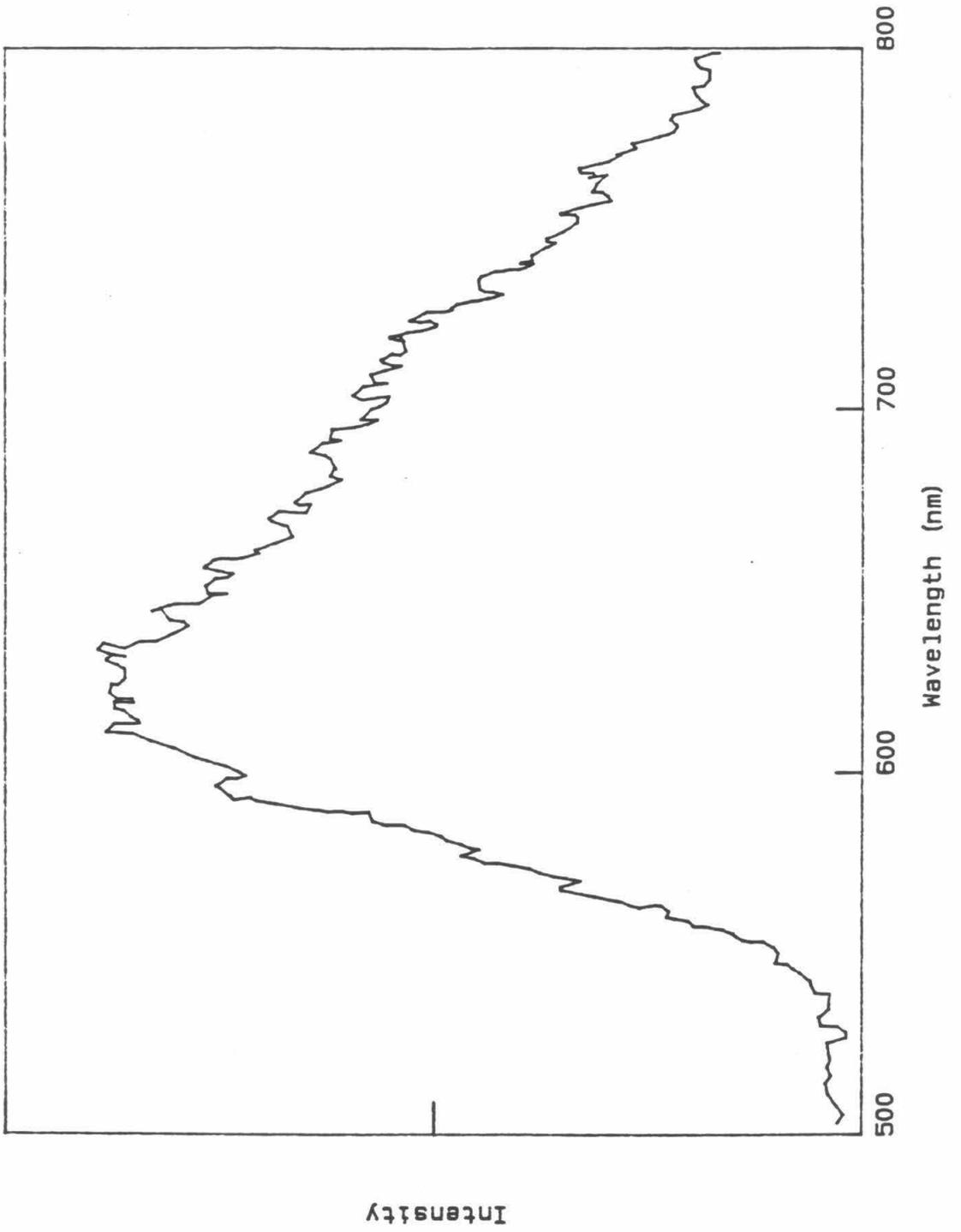
Solvate	Emission Maximum (nm, 77 K)	Lifetime (ns, 300 K)
$\cdot\text{C}_6\text{H}_6$	667	870
$\cdot\text{C}_6\text{H}_5\text{Cl}$	674	470
$\cdot\text{CHCl}_3$	683	120
$\cdot\text{CH}_2\text{Cl}_2$	709	71

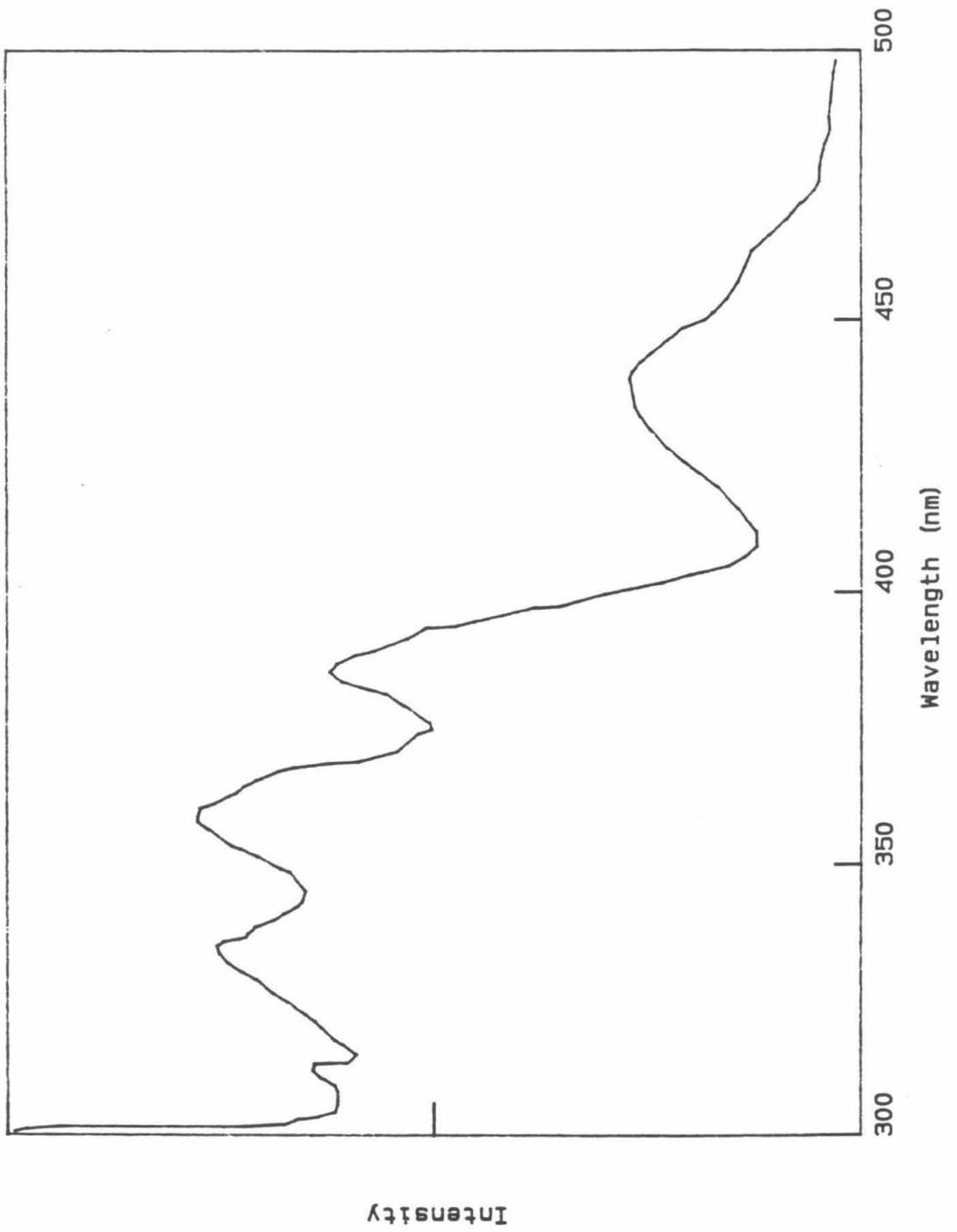
(v/v) mixed solvent system (see Figure 5.8). The emission spectrum is comprised of two bands, one of which has its maximum at ~626 nm and the other which manifests itself as a shoulder to lower energy, λ_{max} ~705 nm. The lower energy shoulder may be residual dichloromethane interacting with the $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2$ molecule, since $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$ has its emission maximum at ~716 nm. The higher energy emission band may be ascribed to the interaction of 2-methyl-tetrahydrofuran with the $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2$ excited state. The excitation spectrum of the $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2$ glass, on the other hand, exhibits excitation bands with identical peak maxima to those recorded for the solid excitation spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$. Apparently, the solvate interacts strongly with the lowest energy excited state but not with any of the higher energy states observed in the absorption and excitation spectra. (Similar emission maxima shifts were observed for the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ complex in this mixed solvent system.)

Considering the sensitivity of the emissive excited state to a single solvate molecule in the crystal, it is not surprising that no luminescence is observed in fluid solution. In fluid solution the relaxation of molecular rigidity and the increased frequency of solvent collision with the $\text{Pt}_2(\mu\text{-dppm})_2\text{X}_2$ molecule results in more deactivation pathways accessible to the excited state.

In order to gain some information about the primary deactivation pathway available to the excited state in the solid, and also to more clearly chart the effect of temperature on the luminescent lifetime, variable temperature studies on $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{C}_6\text{H}_6$ were performed.²³ The results show that the lifetime of the excited state is indeed

- Figure 5.8**
- (a) Corrected emission spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$ in a 1:1 dichloromethane:2-methyltetrahydrofuran glass (v/v) at 77 K. Emission maximum at 626 nm with a shoulder at 705 nm.
- (b) Corrected excitation spectrum of $\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$ in a 1:1 dichloromethane:2-methyltetrahydrofuran glass (v/v) at 77 K. Absorption maxima at 438, 388, 360, 335 nm.





temperature dependent, increasing from 900 ns at 300 K to 27.3 μ s at 77 K. A similarly dramatic increase in the luminescent intensity was observed between 300 and 77 K implying that the variation of the lifetime with temperature is primarily due to variations in the non-radiative rate constant, k_{nr} . Therefore, $k_{obs} \approx k_{nr}$, and the Arrhenius expression may be used in this study. Plotting $\ln k_{obs}$ vs. $1/T$ and fitting the linear portion of these data to the Arrhenius expression:

$$k_{obs} = A \exp(-E_a/RT)$$

(where $k_{obs} = 1/\tau_{obs}$, A is the pre-exponential factor, and E_a is the activation energy for the thermally accessible, non-radiative pathway leading to the deactivation of the excited state) yields values of $E_a = 1450 \text{ cm}^{-1}$ and $A = 1 \times 10^9 \text{ sec}^{-1}$. The data are listed in Table 5.3 and plotted in Figure 5.9. These results suggest that the major deactivation pathway accessible to the excited state lies $\sim 1450 \text{ cm}^{-1}$ to higher energy; however, since there is no vibrational mode readily assignable to this energy, the interpretation of this result is nebulous.

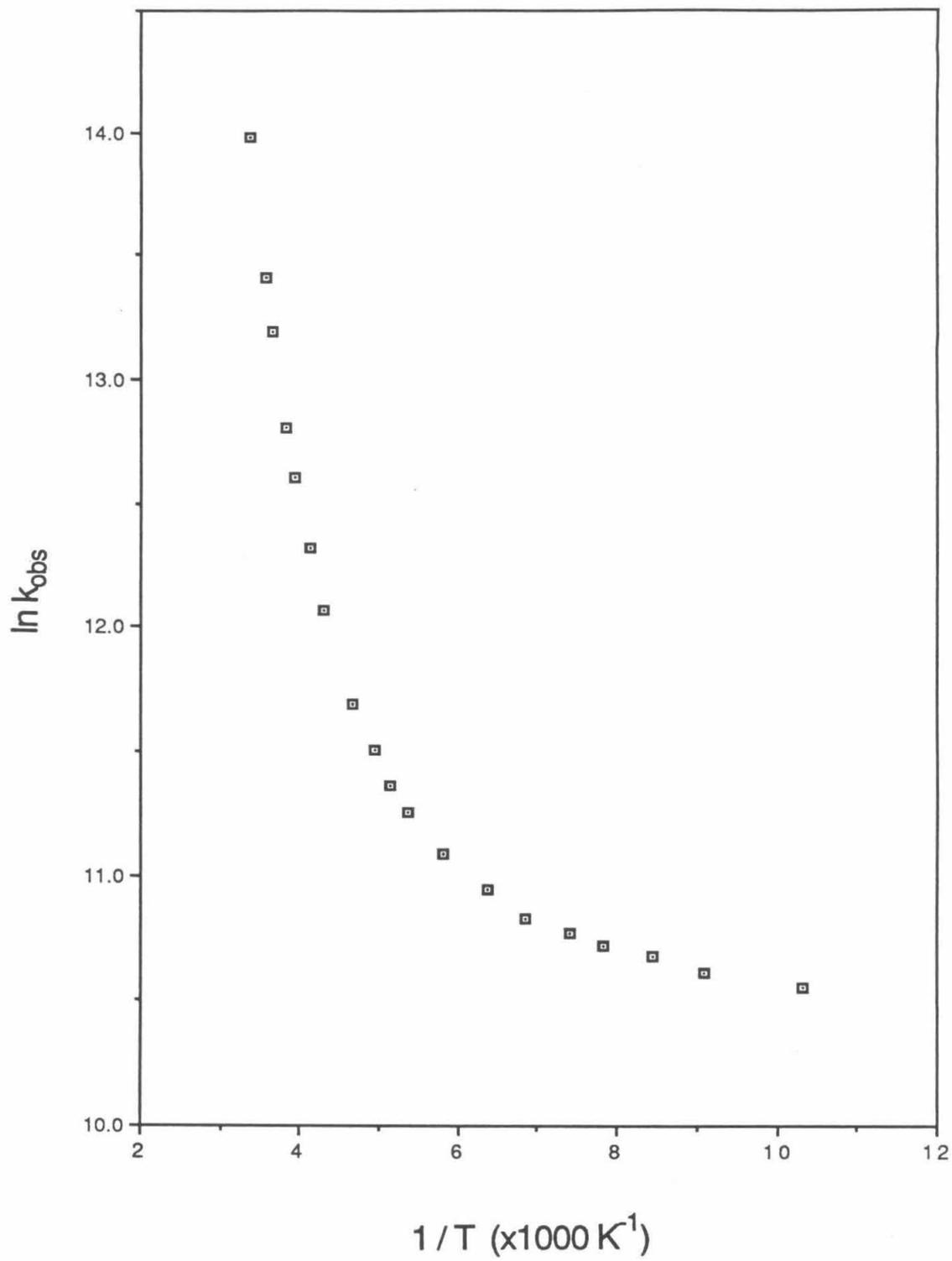
Several cationic phosphine and amine complexes of the form $[\text{Pt}_2(\mu\text{-dppm})_2\text{L}_2]^{2+}$ were synthesized and studied. Their electronic absorption and emission spectra are presented in Figures 5.10, 5.11, and 5.12. The absorption spectra of the three complexes are nearly identical; two resolved bands at $\sim 400 \text{ nm}$ and $\sim 340 \text{ nm}$ with molar extinctions in the range of $5000 - 25,000 \text{ M}^{-1} \text{ cm}^{-1}$. As observed for the neutral halide complexes, the phosphine complexes only luminesce in the solid state and frozen solution, not in fluid solution. Of the three phosphine complexes studied, the PPh_3 adduct exhibits the most intense luminescence with a emission maximum at 583 nm , while the PPhMe_2 adduct is the

Table 5.3

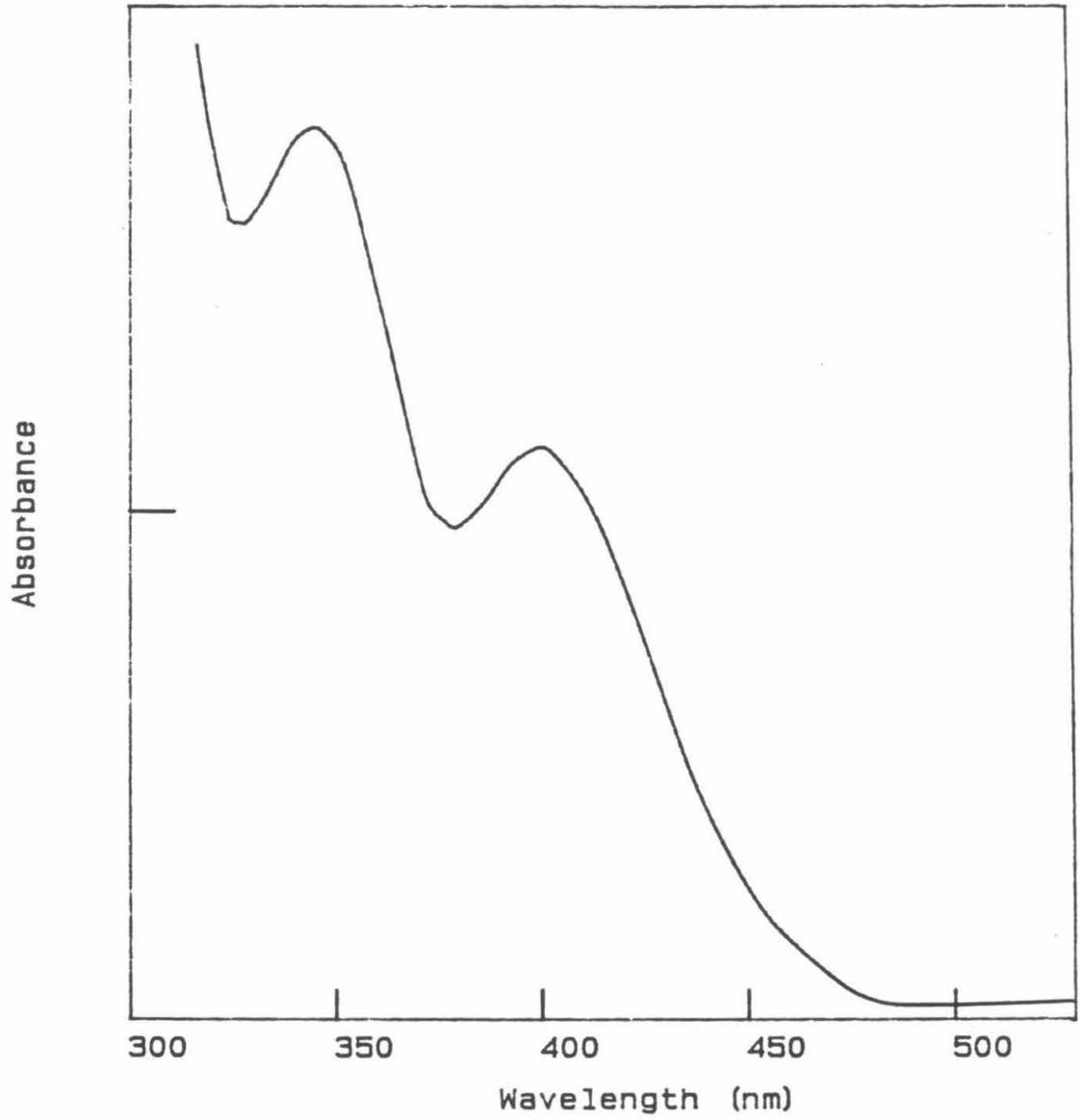
Pt₂(μ-dppm)₂Cl₂·C₆H₆ Variable Temperature Data

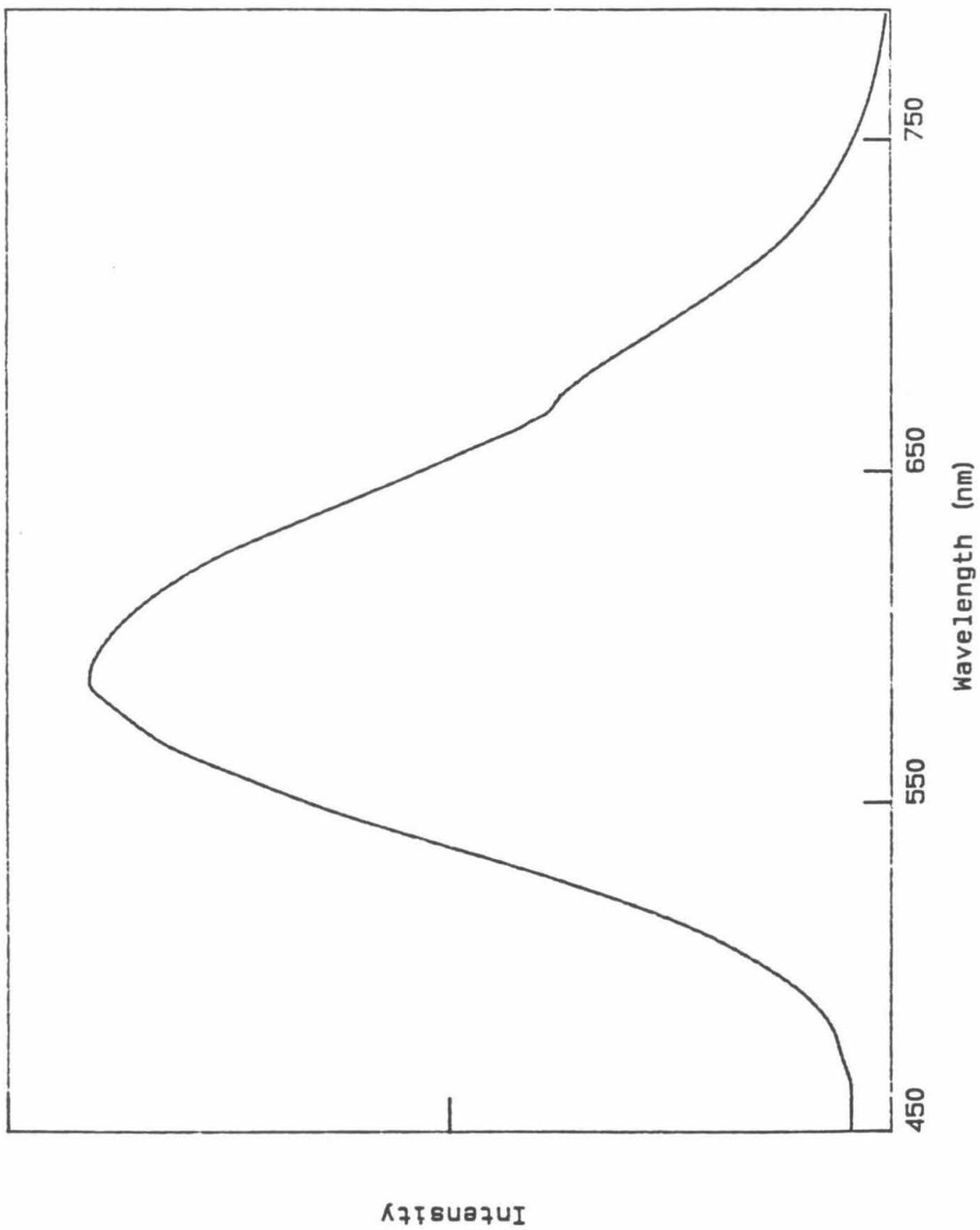
τ_{obs} (μs)	k_{obs} (sec^{-1})	$\ln k_{\text{obs}}$	T (K)	1/T ($\times 1000$ K)
0.85	1176624	13.98	296	3.38
1.50	667999	13.41	279	3.58
1.87	535407	13.19	273	3.66
2.73	365656	12.81	261	3.83
3.33	300492	12.61	254	3.94
4.47	223620	12.32	243	4.12
5.74	174094	12.07	233	4.29
8.41	118844	11.69	214	4.67
10.01	99898	11.51	203	4.93
11.51	86872	11.37	195	5.13
12.87	77695	11.26	187	5.35
15.28	65445	11.09	172	5.81
17.61	56791	10.95	157	6.37
19.74	50657	10.83	146	6.85
21.13	47320	10.77	135	7.41
22.13	45187	10.72	128	7.84
23.12	43255	10.68	119	8.44
24.70	40483	10.61	110	9.09
26.27	38062	10.55	97	10.31

Figure 5.9 Arrhenius plot for $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{C}_6\text{H}_6$ (τ measured at 670 nm).

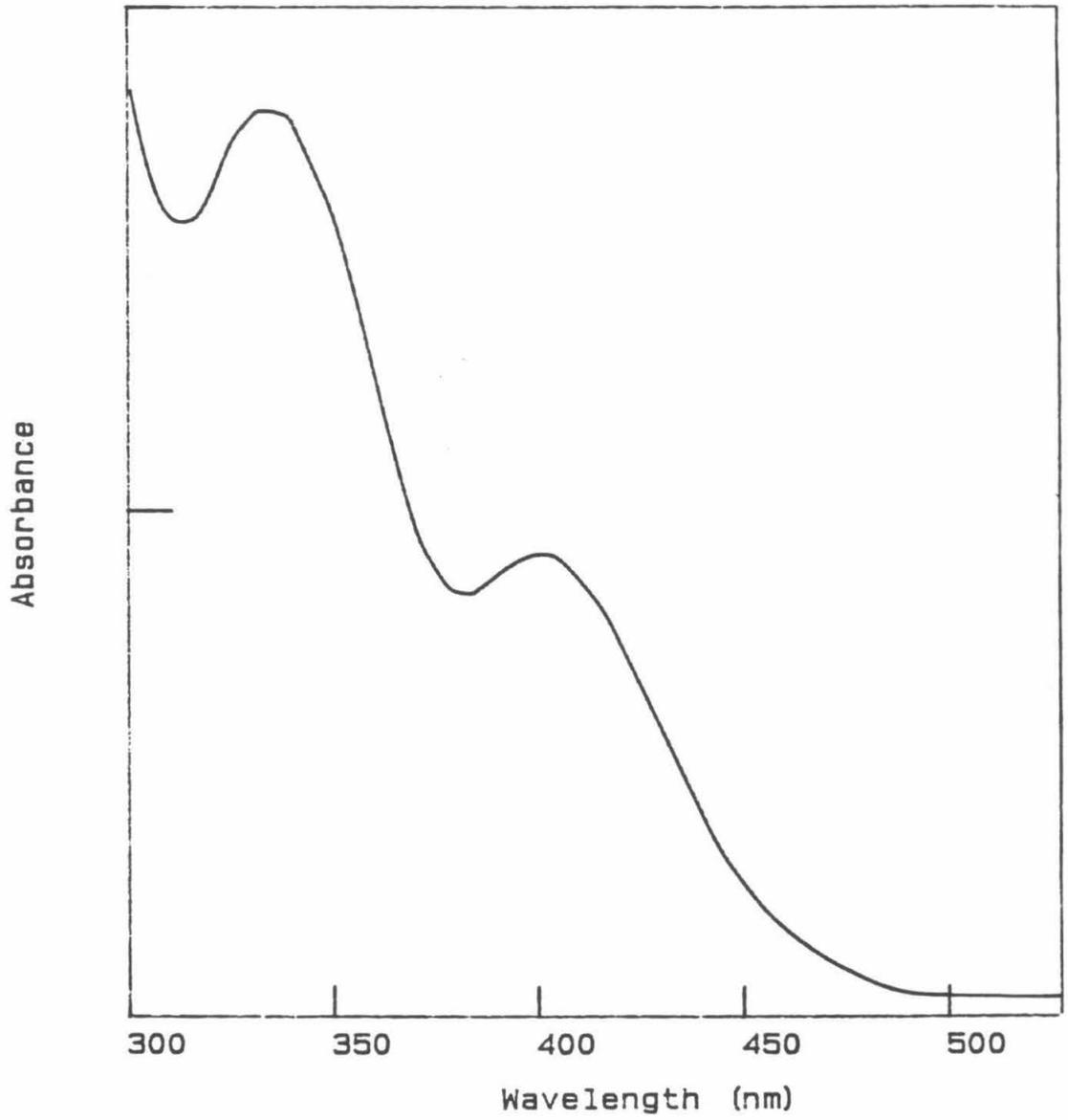


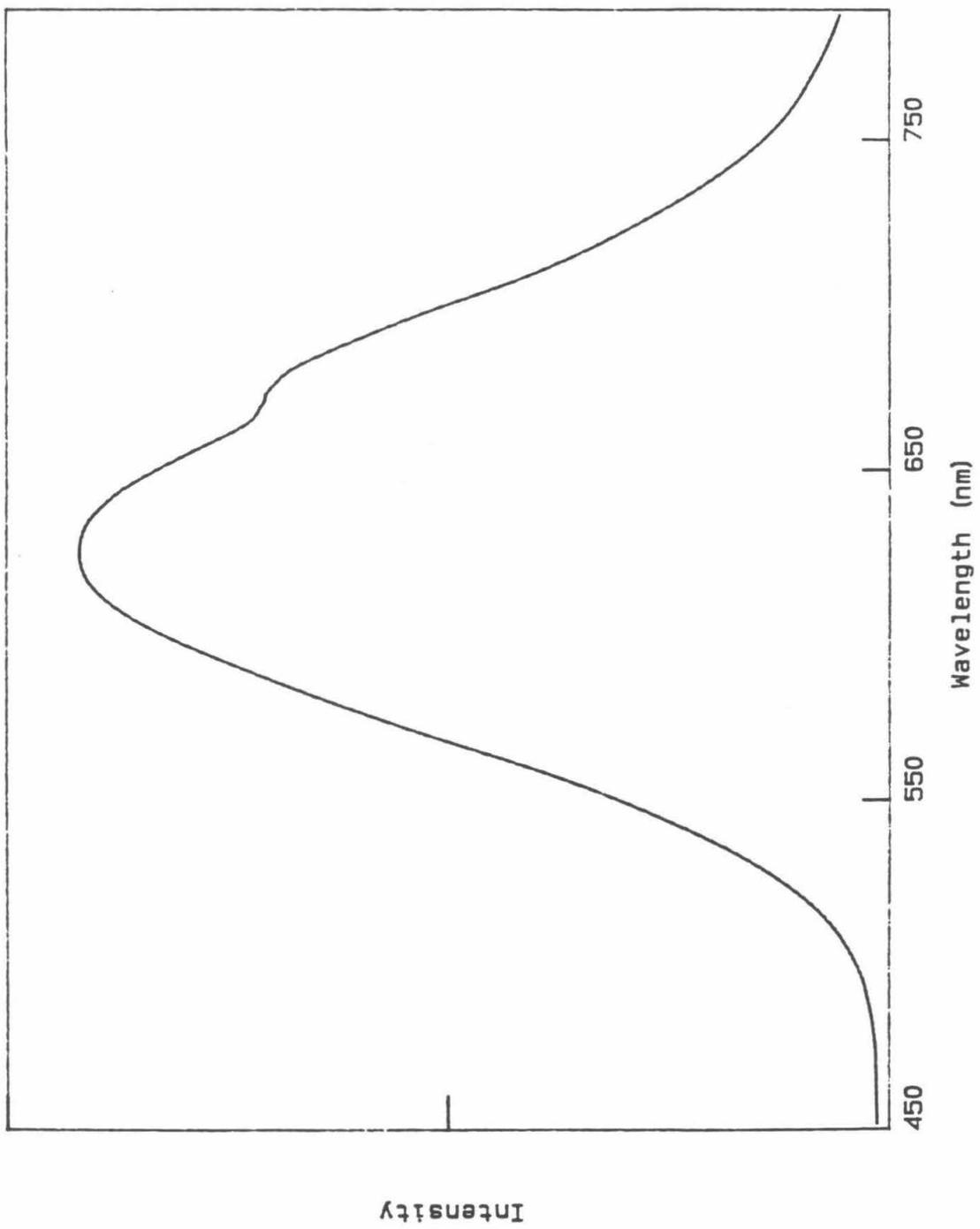
- Figure 5.10**
- (a) Electronic absorption spectrum of $[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_3)_2](\text{PF}_6)_2$ in ethanol (300 K). Absorption maxima at 404 and 350 nm.
- (b) Uncorrected emission spectrum of $[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_3)_2](\text{PF}_6)_2$, powder sample, 300 K. Emission maximum at 583 nm. (A grating anomaly occurs at ~667 nm.)



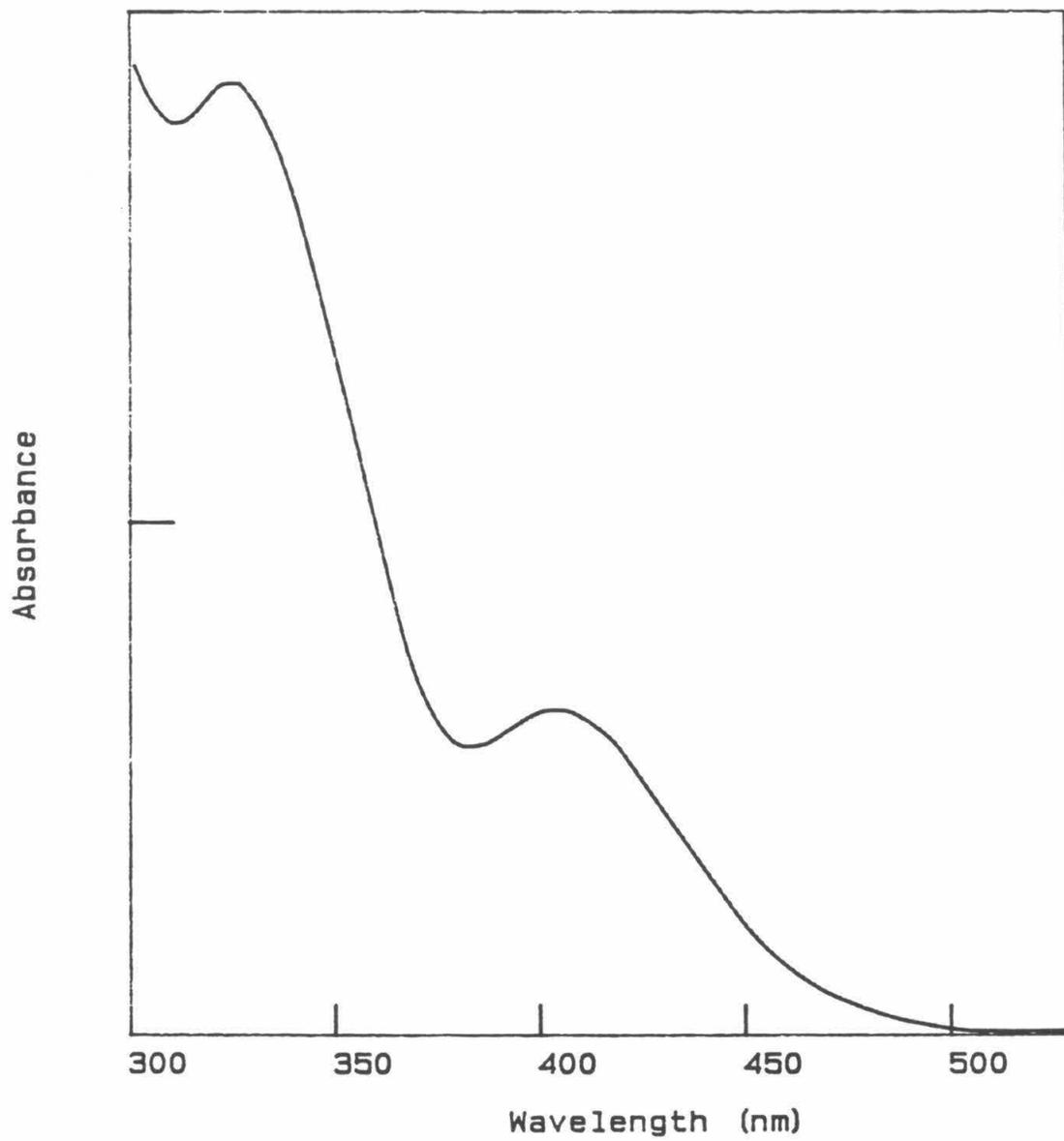


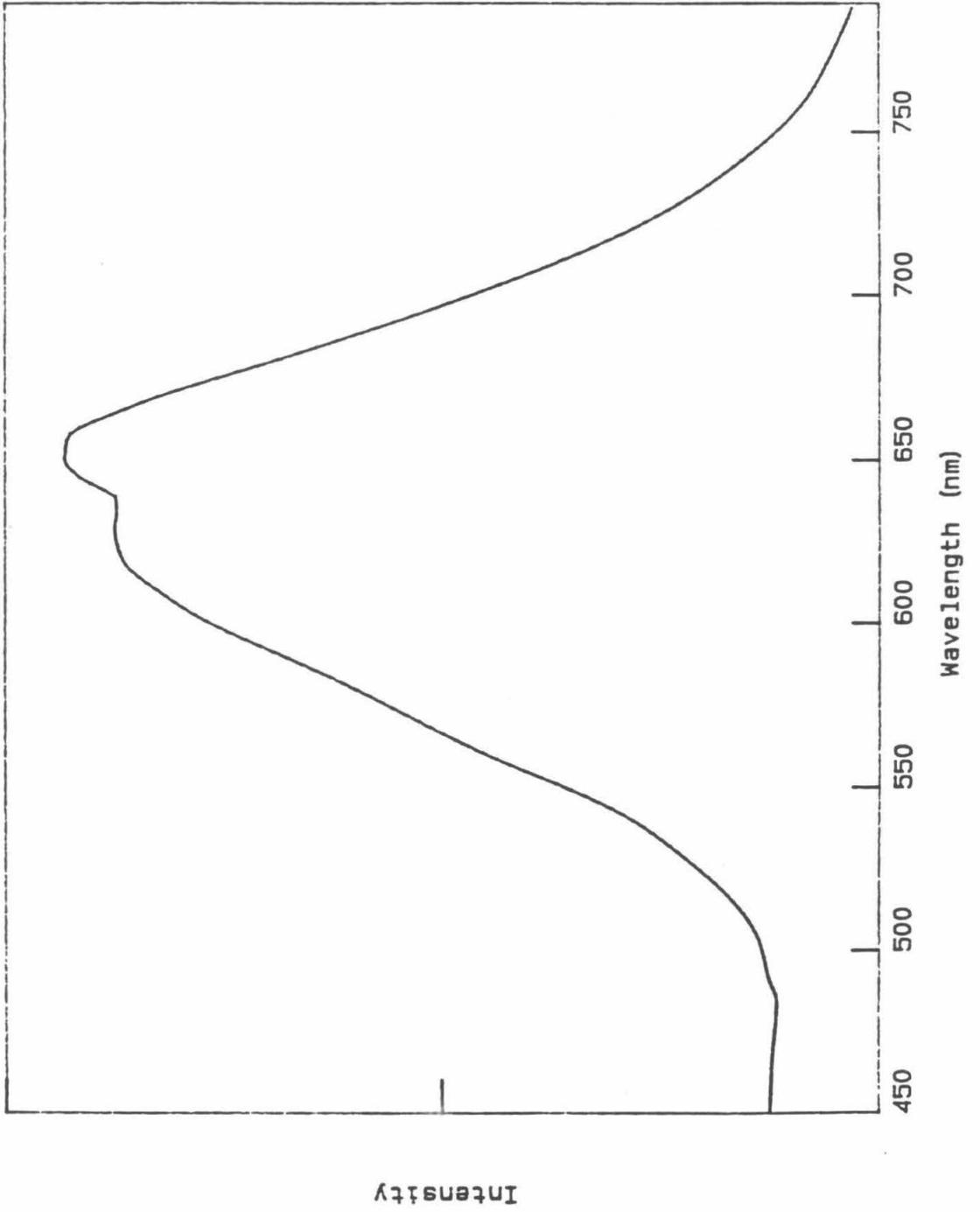
- Figure 5.11**
- (a) Electronic absorption spectrum of $[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_2\text{Me})_2](\text{PF}_6)_2$ in acetonitrile (300 K). Absorption maxima at 400 and 335 nm.
- (b) Uncorrected emission spectrum of $[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_2\text{Me})_2](\text{PF}_6)_2$, powder sample, 300 K. Emission maximum at 624 nm. (A grating anomaly occurs at ~667 nm.)





- Figure 5.12**
- (a) Electronic absorption spectrum of $[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPhMe}_2)_2](\text{PF}_6)_2$ in acetonitrile (300 K). Absorption maxima at 404 and 330 nm.
- (b) Uncorrected emission spectrum of $[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPhMe}_2)_2](\text{PF}_6)_2$, powder sample, 300 K. Emission maximum at ~675 nm. (A grating anomaly occurs at ~667 nm.)





weakest with its emission maximum at 675 nm at 300 K. Similar emission profiles were obtained for glass samples, run at 77 K, where the phosphine complex was dissolved in a 5:2 ethanol/methanol (v/v) mixed solvent system.

Attempts to measure the luminescent lifetime of the phosphine complexes were unsuccessful at both 300 and 77 K in the solid state. On the other hand, glass samples (5:2 ethanol:methanol) of the phosphine complexes gave luminescent lifetimes at 77 K on the order of 20 μ s. The luminescent lifetimes measured for the three complexes were within experimental error, suggesting that they were the same or similar molecular species. Although these results are rather confusing, in light of the phosphine solubilizing ability of ethanol, the possibility exist that all the complexes were $\text{Pt}_2(\mu\text{-dppm})_2(\text{CH}_3\text{CH}_2\text{OH})_2^{2+}$ or $\text{Pt}_2(\mu\text{-dppm})_2(\text{CH}_3\text{OH})_2^{2+}$. These results indicate that axial ligand dissociation may be a facile deactivation pathway of the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ excited state in solution.

A list of electronic absorption data of several platinum(I) and palladium(I) complexes is presented in Table 5.4.^{8,20,24,25} Several similarities between the electronic absorption spectra of Pt(I) and Pd(I) dimers can be seen immediately. Both systems exhibit a multitude of electronic transitions, none of which is distinct or well resolved from the other. The extinction coefficients for the high energy transitions are quite large, characteristic of spin and orbitally allowed transitions. Difficulties arise when trying to compare the low energy electronic transitions of the Pt halide(X) molecule with those of either the Pt phosphine(L) or Pd halide(X) complexes. From comparison

Table 5.4

Absorption Band Maxima of Some Pt(I) and Pd(I) Dimers[†]Platinum(I) Dimers

$\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$	465 (<100), 427 (200), 380 (~2000), 360 (2600), 323(7800)
$\text{Pt}_2(\mu\text{-dppm})_2\text{Br}_2$	478, 441, 390, 360, 323
$[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_3)_2](\text{PF}_6)_2$	404 (15500), 350 (23000)
$[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_2\text{Me})_2](\text{PF}_6)_2$	400 (7480), 335 (13400)
$[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPhMe}_2)_2](\text{PF}_6)_2$	404 (5260), 330 (14000)
$[\text{Pt}_2(\mu\text{-dppm})_2(\text{SnCl}_3)_2$	450sh(15700), 438(16800), 360(24600), 303 (26700), 276sh (30800)

Palladium(I) Dimers

$\text{Pd}_2(\mu\text{-dppm})_2\text{Cl}_2$	416 (7550), 347 (16800), 293 (25900)
$\text{Pd}_2(\mu\text{-dppm})_2\text{Br}_2$	428 (10600), 364 (17500), 301(23100), 285sh (21300), 258sh (29400)
$\text{Pd}_2(\mu\text{-dppm})_2\text{I}_2$	488 (13000), 439 (11900), 394(11900), 313 (20000), 280sh (23900), 262sh(28500)
$\text{Pd}_2(\mu\text{-dppm})_2(\text{SCN})_2$	426 (14600), 303(21200), 259sh(30400)
$\text{Pd}_2(\mu\text{-dppm})_2(\text{NCO})_2$	394(7390), 326sh(13500), 290sh(36300)
$\text{Pd}_2(\mu\text{-dppm})_2(\text{N}_3)_2$	428 (19900), 305sh (19700), 287sh(25000), 265sh (28100)
$\text{Pd}_2(\mu\text{-dppm})_2(\text{NO}_2)_2$	434 (11500), 362 (9300), 316sh (13800), 294sh (22000)
$\text{Pd}_2(\mu\text{-dppm})_2(\text{SnCl}_3)\text{Cl}$	495(14200), 422 (8640), 340 (18200), 294sh (18100), 260sh (31,700)
$\text{Pd}_2(\mu\text{-dppm})_2(\text{SnCl}_3)_2$	536 (21200), 500sh (17900), 417sh(6120), 353(19000), 261sh(35000)

[†]Some of the data taken from references 8, 20, 24, and 25.

of the molar extinction coefficients, the lowest energy electronic absorption bands observed for $\text{Pt}_2(\mu\text{-dppm})_2\text{X}_2$ are attributable to transitions different from those observed in the analogous $\text{Pd}_2(\mu\text{-dppm})_2\text{X}_2$ and $\text{Pt}_2(\mu\text{-dppm})_2\text{L}_2^{2+}$ complexes. Also, the temperature dependence of the emissive excited state lifetime of $\text{Pt}_2(\mu\text{-dppm})_2\text{X}_2$ is significantly different from that of $\text{Pt}_2(\mu\text{-dppm})_2\text{L}_2^{2+}$.

Finally, for $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$, no intensity overlap between the lowest energy absorption band and the emission band is observed. This large "window" may be attributed to a significant Stokes shift, or it may be indicative of the presence of a very weak absorption band which is, as of yet, unobserved. The emission spectra of the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ molecule show significant environment dependence, whereas, the corresponding excitation spectra are all identical to one another, showing no environment dependence. This apparent contradiction may be best explained by the presence of a lower energy absorption, too weak to be observed, which is directly related to the observed emission and is environment sensitive. Thus, the higher energy absorption bands observed in the excitation spectra, which are not environment sensitive, are not directly related to the emission band; instead, they merely communicate by various deactivation pathways with the lower energy state responsible for the observed luminescence.

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23. As the temperature was lowered from 300 to 77 K all of the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{solvate}$ lifetimes increased. In a few cases biphasic decays were observed. This biphasic decay may be due to the presence of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ in different solvate environments. If, for example, one environment allows for greater solvent interaction with the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ excited state than another environment, a superposition of these different luminescent lifetimes may be observed, resulting in a net multiphasic decay. As for $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot \text{C}_6\text{H}_6$, this complex exhibited a monophasic decay throughout the entire temperature range studied.
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CHAPTER 6**Energy Transfer Studies of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$**

INTRODUCTION

An understanding of the nature of the excited state is necessary in order to intelligently discuss the observed excited state reaction chemistry. Traditional techniques used to elucidate excited state information are not amenable to this $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ system since (i) these complexes do not luminesce in fluid solution at room temperature and (ii) absorption and/or emission band fine structure are not observed at low temperature. Instead, indirect methods for gaining the required information must be sought.

From data presented in Chapter 5, it is believed that the absorption band corresponding to the observed emission is not observed in solution electronic absorption or solid state excitation spectra. Both excitation spectra of the benzene and dichloromethane solvates of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ show an $\sim 3000\text{ cm}^{-1}$ window where no other transitions between the lowest absorption band and the emission band are observed (see Figures 5.4ab and 5.7ab). (A Stokes shift of this type or magnitude makes little physical sense; therefore, these high energy absorptions have been ruled out as being directly related to the emissive excited state.) The absorption band corresponding to the emissive excited state is assumed to have an extremely small molar extinction coefficient and is, logically, attributed to a singlet-triplet transition. Besides this argument, there is no other experimental evidence supporting the spin assignment of the emissive excited state. In energy transfer experiments, the conservation of spin is quite important. Investigating the energy transfer rates between $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ and other known triplet excited states may provide direct evidence for the triplet spin assignment of the emissive excited state.

From emission and lifetime studies of the solvate effect on powder $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ samples, it has been determined that the excited state is environment sensitive. Exactly what the excited state energy of the complex in solution is unclear since direct measurement is impossible; however, indirect determination of its approximate energy is possible by conducting energy transfer experiments aimed at determining the energy at which the rate of transfer from donor to acceptor begins to drop off. Since $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ does not luminesce in fluid solution its role must be that of acceptor. A series of molecules which do luminesce in fluid solution and have excited state lifetimes greater than several hundred nanoseconds are employed as donors; in this experiment $\text{W}_6\text{X}_8\text{Y}_6^{2-}$ (X, Y = Cl, Br, I), ReO_2L_4^+ (L = pyridine and its derivatives), and $\text{Ru}(\text{bpy})_3^{2+}$ (bpy = 2,2'-bipyridine) were used.

EXPERIMENTAL

Materials

All solvents, unless otherwise noted, were reagent grade and used as received without further purification. Dichloromethane (Burdick and Jackson, UV grade), used in the energy transfer quenching experiments, was dried and stored in glass solvent pots as described in the experimental section of Chapter 5.

Donors. Hexanuclear clusters, $(\text{TBA})_2[\text{W}_6\text{I}_{14}]$, $(\text{TBA})_2[\text{W}_6\text{I}_8\text{Cl}_6]$, $(\text{TBA})_2[\text{W}_6\text{Br}_8\text{I}_6]$, $(\text{TBA})_2[\text{W}_6\text{Cl}_8\text{I}_6]$, $(\text{TBA})_2[\text{W}_6\text{Cl}_8\text{Br}_6]$, $(\text{TBA})_2[\text{W}_6\text{Cl}_{14}]$ were prepared by Dr. Thomas Zietlow and used as received.¹ $(\text{TBA})_2[\text{W}_6\text{I}_8\text{Cl}_6]$ required recrystallization from dichloromethane and petroleum ether prior to use. Rhenium dioxo complexes, $[\text{ReO}_2(4\text{-picoline})_4]\text{BPh}_4$, $[\text{ReO}_2(\text{pyridine})_4]\text{BPh}_4$, and $[\text{ReO}_2(\text{d}_5\text{-pyridine})_4]\text{BPh}_4$ were prepared by Dr. Jay Winkler and used as received.² $[\text{ReO}_2(\text{isoquinoline})_4]\text{PF}_6$ was prepared by Mr. John Brewer and used as received.³ $[\text{Ru}(2,2'\text{-bipyridine})_3](\text{PF}_6)_2$ was prepared by Dr. Michael Hopkins and used as received.⁴

Acceptors. $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$ was prepared as described in the experimental section of Chapter 5 and recrystallized from dichloromethane and iso-octane. Anthracene (Aldrich, 98%) was recrystallized from a hot 2:1:1 (v/v/v) acetone:ethanol:water mixture.

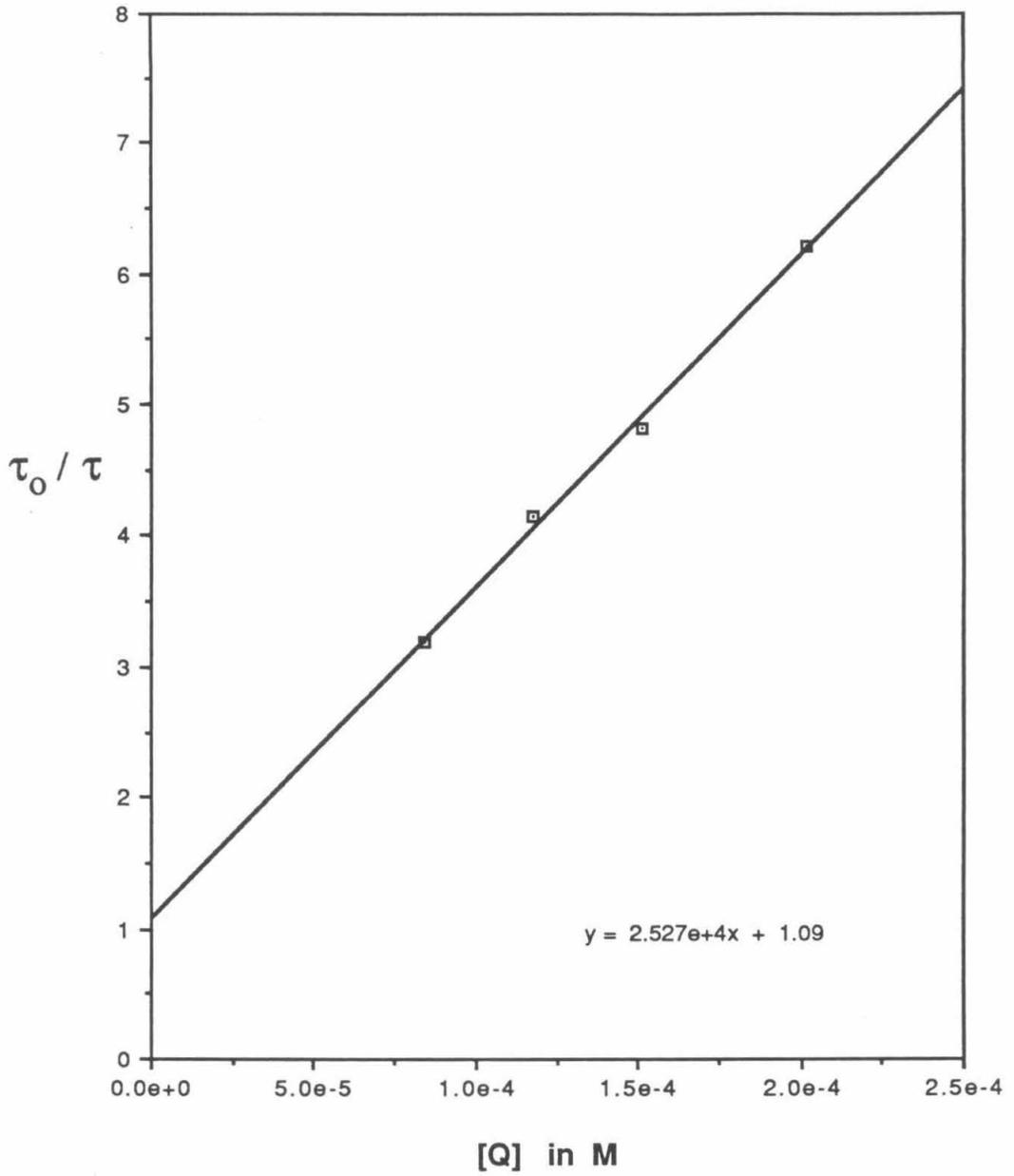
Physical Measurements

Stern-Volmer Quenching Procedures. A Nd:YAG pulsed laser system, previously described,⁵ was used to measure the luminescent lifetimes as a function of quencher concentration. Energy transfer rate constants were calculated using Stern-Volmer analysis as described in the results

and discussion section of Chapter 2. Samples for all energy transfer experiments were prepared by placing the donor (~1 mg) into a quenching cell and evacuating it on a high vacuum line (limiting pressure $< 10^{-3}$ Torr). (The quenching cell is described in the experimental section of Chapter 2.) Four to six mL of dichloromethane was distilled from the storage flask into the quenching cell under vacuum; the exact volume of dichloromethane used was recorded. Acceptor solutions were added following the same procedures for quencher addition outlined in the experimental section of Chapter 2. Figure 6.1 illustrates the energy transfer quenching of $(\text{TBA})_2[\text{W}_6\text{Br}_8\text{I}_6]$ with $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$.

In the ideal energy transfer experiment only the donor molecule should absorb the laser excitation; the acceptor should be unaffected by the laser excitation and remain in its ground state. The majority of the reported energy transfer experiments, however, were carried out using the laser's second harmonic, 355 nm, as the excitation wavelength, a wavelength at which the acceptor, $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$, has an appreciable absorption. Therefore, to insure that the acceptor's excited state was not involved in initiating the energy transfer process, several measurements were made using the laser's first harmonic, 532 nm, as the excitation wavelength; ($\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ has no absorption at this wavelength). Donors, $\text{Ru}(\text{bpy})_3^{2+}$ and $\text{ReO}_2(\text{py})_4^+$ which do have appreciable absorption at 532 nm were used in this control experiment. Nearly identical energy transfer rate constants were calculated for the data collected using 532 nm (no $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ absorption) and 355 nm (significant $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ absorption) excitation wavelengths. Thus, any excited state $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2^*$ molecules generated using 355 nm excitation appear to be

Figure 6.1 Stern-Volmer plot of the energy transfer quenching of $(\text{TBA})_2[\text{W}_6\text{Br}_8\text{I}_6]$ with $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ in dichloromethane.



uninvolved in the energy transfer process.

Emission Spectroscopy. Attempts to observe the luminescence from excited $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ were conducted on the emission instrument built at Caltech and described previously.⁶ An excitation wavelength, 547 nm, was chosen in order to excite only the donor, $\text{ReO}_2(\text{d}_5\text{-py})_4^+$. Emission spectra were recorded between 585 and 825 nm. Samples were placed in NMR tubes equipped with 14/20 ground glass joints and capped with a septum. First, emission spectra of frozen dichloromethane solutions containing only $\text{ReO}_2(\text{d}_5\text{-py})_4^+$ (1×10^{-3} M) were measured. Next, emission spectra of frozen dichloromethane solutions containing 2×10^{-3} M $\text{ReO}_2(\text{d}_5\text{-py})_4^+$ and 1×10^{-2} M $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ were measured. Comparison between these emission spectra revealed subtle changes in the emission band profile of the ReO_2L_4^+ chromophore (not attributable to concentration changes).

RESULTS and DISCUSSION

There are several mechanisms (other than the intrinsic radiative and non-radiative processes) by which an excited state molecule, D^* , can decay to its ground state, D : (1) energy transfer, (2) electron transfer, (3) a chemical reaction (other than electron transfer), (4) spin-catalyzed deactivation, (5) external heavy-atom effect, and (6) complex (excimer and exciplex) formation.⁷

Energy transfer is by far the most common deactivation pathway. Its basic requirement is that the acceptor molecule, A , possess an excited state, A^* , in which the energy required for the $A \rightarrow A^*$ transition is approximately equal to or less than the $D^* \rightarrow D$ transition energy.⁷⁻⁹

The second prevalent deactivation pathway is electron transfer between the excited state donor and a ground state acceptor, such that either (i) $D^* + A \rightarrow D^+ + A^-$ or (ii) $D^* + A \rightarrow D^- + A^+$ is a favorable reaction.^{7,10,11} In the first case (i), the donor must have an accessible oxidized state and D^* must be of high enough energy such that it is energetically favorable to give up an electron from its excited state and reduce the acceptor. Similarly, the same type of requirements exist for the second case (ii), in which the donor must have an accessible reduced state and D^* must be of high enough energy to oxidize the acceptor.

Chemical reactions between coordination complexes which do not involve electron transfer but rather some type of ligand loss or substitution are rare, since they are second order reactions which are usually unable to compete on the timescale of the donor's excited state lifetime.⁷ If, however, such reactions were occurring, they would be

readily detected by changes in the electronic absorption spectrum which are not attributable to the donor or acceptor species.

The latter three processes are very rare, with very little, if any, evidence in support of them occurring between inorganic donors and acceptors.⁷ Therefore, it is supposed that these processes are far too inefficient to compete with the more facile energy and electron transfer deactivation pathways.

Since no changes in the electronic absorption spectrum (aside from changes in acceptor concentration) were noted during the course of the quenching experiment, excited state deactivation by a chemical process (pathway 3) was ruled out. Although the acceptor and all donor molecules exhibit some type of oxidation wave more positive than -1 V vs. SCE, neither acceptor nor donor molecules (other than $\text{Ru}(\text{bpy})_3^{2+}$) have a reduction potential less than +1.5 V vs. SCE. Therefore, deactivation by way of an electron transfer mechanism is not favorable.

The best way to prove energy transfer as being the operative deactivation pathway is to observe the luminescence from the excited acceptor molecule. Since $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$, the acceptor, does not luminesce in fluid solution at room temperature, it is difficult to prove that energy transfer is the deactivation pathway. However, energy transfer mechanisms have been shown to occur even in frozen solutions according to the Perrin formula.⁸ Excitation at 547 nm of a frozen dichloromethane solution containing $\text{ReO}_2(\text{d}_5\text{-py})_4^+$ and $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ did show subtle changes in the emission band profile of the ReO_2L_4^+ chromophore which may be due to emission from excited $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$. This information alone is not definite proof of the energy transfer mechanism; however, combined with the unlikelihood that an

electron transfer or chemical reaction is occurring, electronic energy transfer from the donor to $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ is the most logical deactivation pathway available.

In fluid solution, there are two electronic mechanisms by which electronic energy transfer may occur: coulombic interaction (dipole-dipole mechanism) and exchange interaction (collisional mechanism). Detailed theoretical analysis of these mechanisms is available elsewhere.^{12,13} However, the main difference between the two is that energy transfer by the coulombic interaction depends upon the oscillator strength of the $\text{D}^* \rightarrow \text{D}$ and, more importantly, $\text{A} \rightarrow \text{A}^*$ (A = acceptor, D = donor) transitions, whereas, the energy transfer by the exchange interaction is independent of these terms.⁸ Thus, singlet-singlet transitions $[\text{D}(\text{S}_1) + \text{A}(\text{S}_0) \rightarrow \text{D}(\text{S}_0) + \text{A}(\text{S}_1)]$ are spin allowed by both coulombic and exchange interactions; triplet-triplet transitions $[\text{D}(\text{T}_1) + \text{A}(\text{S}_0) \rightarrow \text{D}(\text{S}_0) + \text{A}(\text{T}_1)]$ are spin forbidden by the coulombic interaction and spin allowed the exchange interaction; and both triplet-singlet $[\text{D}(\text{T}_1) + \text{A}(\text{S}_0) \rightarrow \text{D}(\text{S}_0) + \text{A}(\text{S}_1)]$ and singlet-triplet $[\text{D}(\text{S}_1) + \text{A}(\text{S}_0) \rightarrow \text{D}(\text{S}_0) + \text{A}(\text{S}_1)]$ are spin forbidden by both interactions (and rarely observed).

The majority of donors in this study are either tungsten hexanuclear dianionic clusters or rhenium dioxo monocationic monomers. Although the hexanuclear clusters are significantly larger in size than the rhenium monomers, they are both octahedral-like complexes. Also, although the hexanuclear clusters are tungsten(II) d^4 and the monomers are rhenium(V) d^2 molecules, the lowest energy absorption and emission in both systems arises from $\text{d} \rightarrow \text{d}$ transitions.^{1,2} It has been experi-

mentally observed in organic systems that triplet-triplet energy transfer rates are not very sensitive to the type of orbitals involved.⁸ Therefore, the difference in donor type, W vs. Re, is presumed to be of little consequence in this study.

The electronic energy transfer data between the various donors ($W_6X_8Y_6^{2-}$, $ReO_2L_4^+$, and $Ru(bpy)_3^{2+}$) and the acceptor ($Pt_2(\mu-dppm)_2Cl_2$) are presented in Table 6.1 and plotted in Figure 6.2.¹⁴ Since electronic energy transfer is observed from donor molecules, known to have triplet emissive excited states, to the acceptor, further evidence supporting the triplet spin label of the $Pt_2(\mu-dppm)_2Cl_2$ excited state is gained.

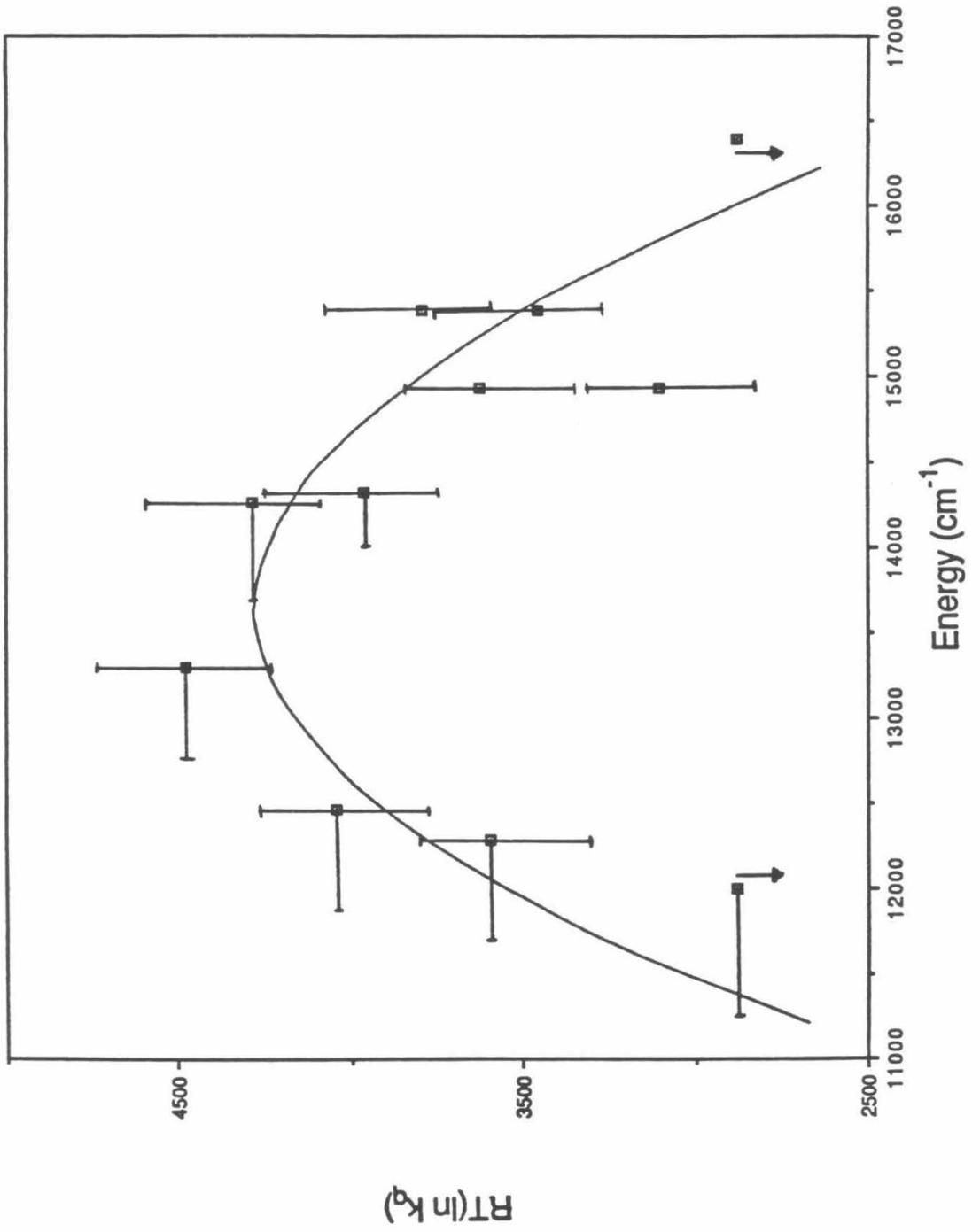
The highest energy transfer rate of $2.1 \times 10^9 \text{ M}^{-1} \text{ sec}^{-1}$ is observed with $(TBA)_2[W_6Br_8I_6]$ ($\lambda_{max} = 752 \text{ nm}$). It is at this approximate energy that facile energy transfer between the donor and acceptor is believed to be occurring. Using the equation, $k_{dif} (\text{M}^{-1} \text{ sec}^{-1}) = (2 \times 10^5) \cdot T/\eta$ where $T = 300 \text{ K}$ (temperature) and $\eta = 0.004 \text{ P}$ (viscosity), the diffusion controlled rate constant (k_{dif}) in dichloromethane is calculated to be $\sim 1.5 \times 10^{10} \text{ M}^{-1} \text{ sec}^{-1}$. In triplet-triplet energy transfer experiments, exothermic energy transfer is not expected to occur with energy transfer rate constants, k_{et} , greater than k_{dif} . In fact, in many cases, upper limits for energy transfer quenching rate constants have been observed to be several orders of magnitude lower than the k_{dif} when one or both donor/acceptor molecules is a transition metal complex.⁷ Activation barriers to energy transfer (similar to those observed in bimolecular photochemical reactions), although not required from a theoretical standpoint, may exist which keep the rate constants

Table 6.1

Energy Transfer Data Using $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ as the Acceptor

Donor	Emission Maximum		k_q	RT/lnk_q
	nm	cm^{-1}		
$\text{Ru}(\text{bpy})_3^{2+}$	610	16393	$< 1 \times 10^6$	< 2880
$\text{ReO}_2(\text{py})_4^+$	650	15385	8.1×10^7	3796.6
$\text{ReO}_2(\text{py-d}_5)_4^+$	650	15385	1.6×10^7	3458.4
$\text{ReO}_2(4\text{-pic})_4^+$	670	14925	2.9×10^6	3102.3
$\text{ReO}_2(\text{isoquin})_4^+$	670	14925	3.6×10^7	3627.5
$\text{W}_6\text{I}_{14}^{2-}$	698	14327	1.8×10^8	3963.0
$\text{W}_6\text{I}_8\text{Cl}_6^{2-}$	701	14265	8.5×10^8	4286.8
$\text{W}_6\text{Br}_8\text{I}_6^{2-}$	752	13298	2.1×10^9	4475.2
$\text{W}_6\text{Cl}_8\text{I}_6^{2-}$	802	12469	2.6×10^8	4039.7
$\text{W}_6\text{Cl}_8\text{Br}_6^{2-}$	814	12285	3.1×10^7	3596.2
$\text{W}_6\text{Cl}_{14}^{2-}$	833	12005	$< 1 \times 10^6$	< 2880

Figure 6.2 Plot of the energy transfer data between donors ($W_6X_8Y_6^{2-}$, $ReO_2L_4^+$, $Ru(bpy)_3^{2+}$) and the acceptor, $Pt_2(\mu-dppm)_2Cl_2$. (For an explanation of x axis error bars, please see reference 14.)



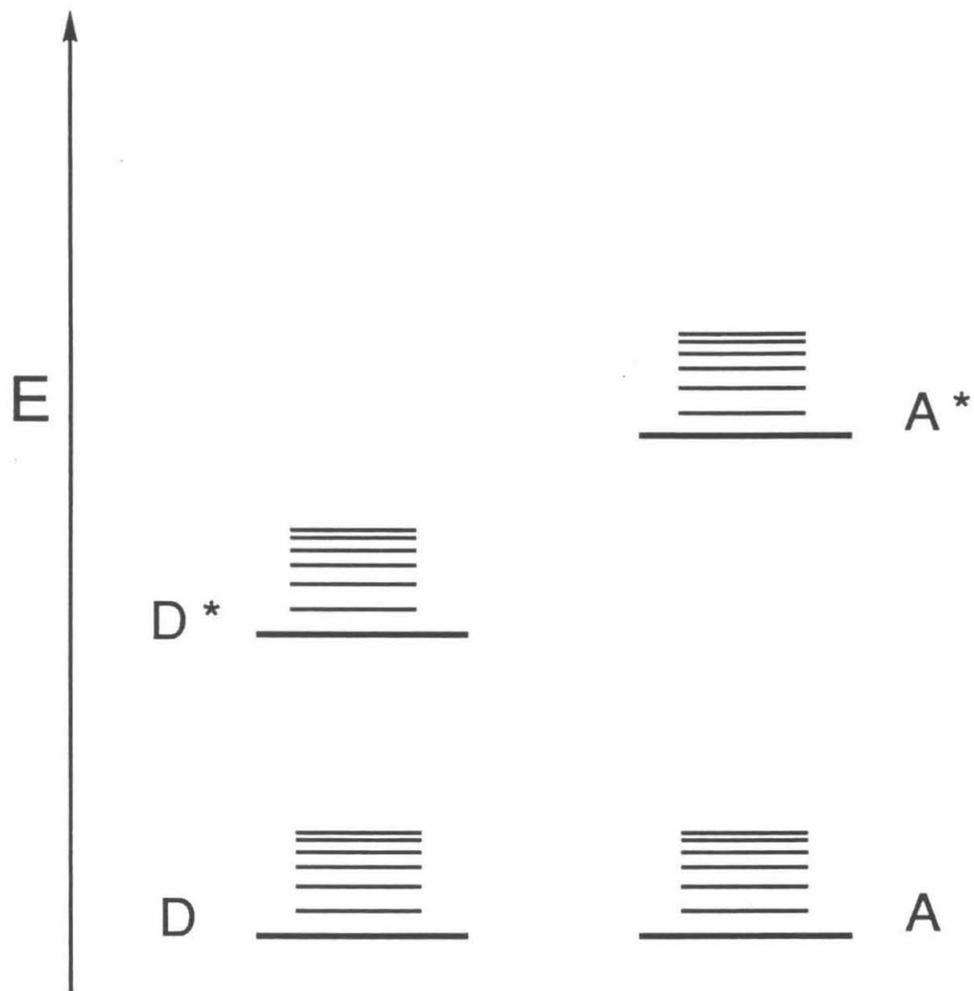
of the most facile energy transfer reactions below k_{dif} .⁸

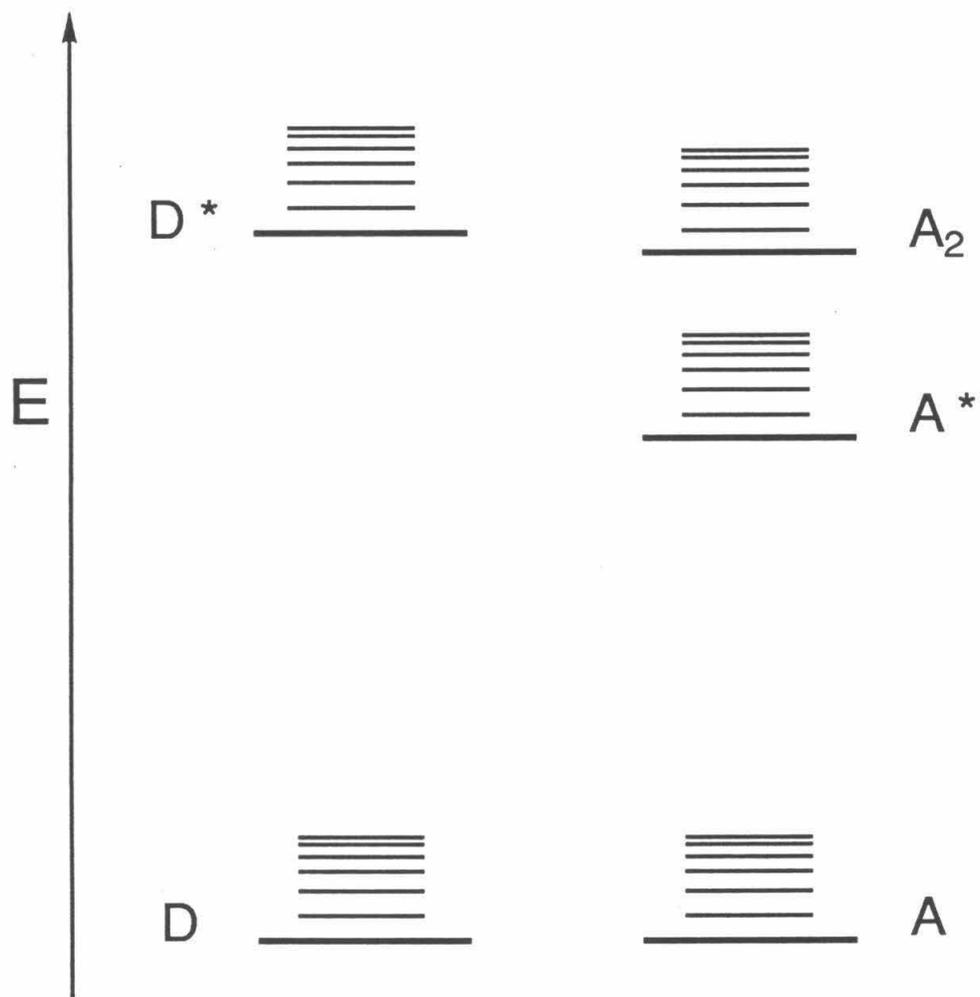
When the acceptor's excited state energy level surpasses the donor's energy level, and the energy transfer becomes endothermic by several kilocalories per mole, energy transfer quenching is essentially shut down (see Figure 6.3a). This expected drop off in energy transfer rate constants is clearly observed in Figure 6.2. An estimate of the acceptor's excited state energy is obtained from the point at which the quenching rate begins to drop off as the energy transfer reaction becomes endothermic. Thus, the estimated energy of the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ excited state in dichloromethane at room temperature is $\sim 13000 \text{ cm}^{-1}$.

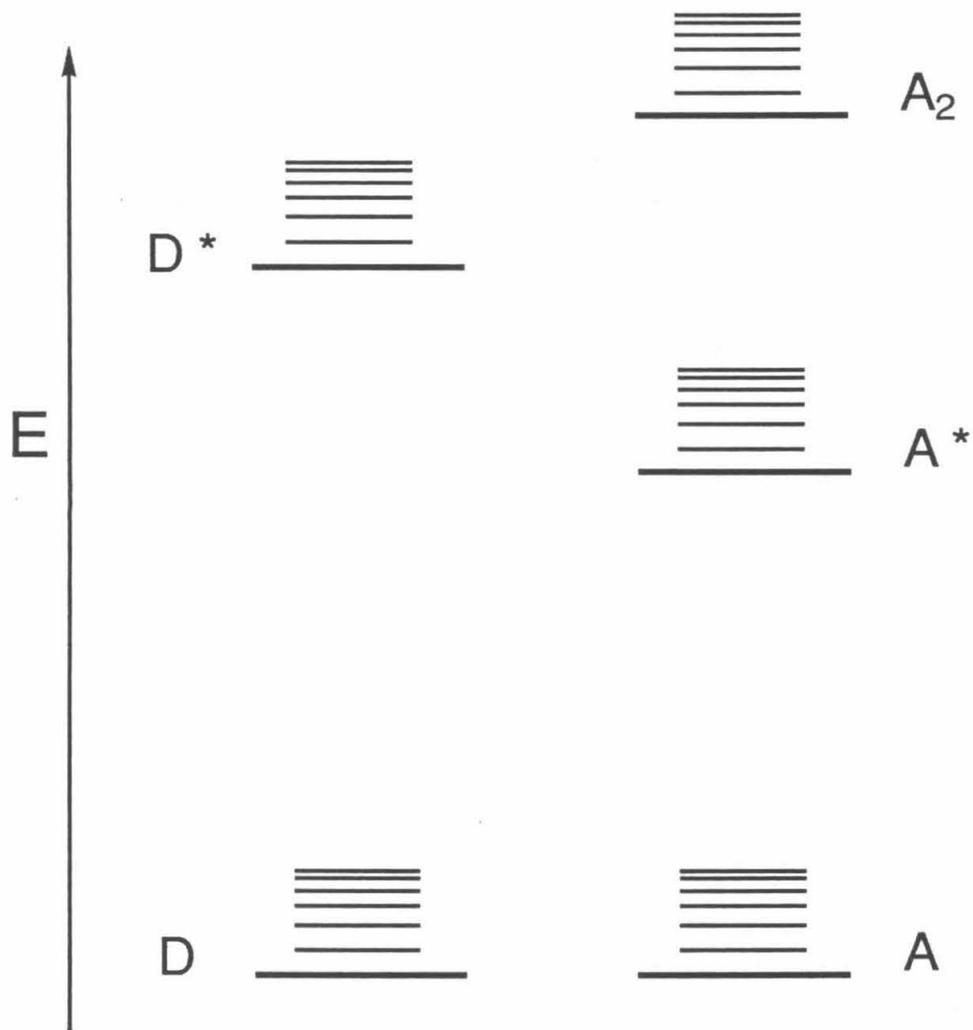
A similar drop off in transfer rates is observed for reactions in which the donor's excited state energy level is significantly greater than the acceptor's energy level (exothermic energy transfer reaction). Energy transfer theory,⁹ which closely parallels Marcus theory for electron transfer, predicts a drop off of energy transfer rate constants in this the "inverted region".^{10,11,17,18} In most cases, this theoretically predicted decrease in energy transfer rates is not observed; instead as the energy level difference between the donor and the acceptor becomes more exothermic, the energy transfer rate constants tend to saturate at some value, usually near k_{dif} .⁸

The most straightforward explanation for the saturation of the energy transfer rate constants as the system becomes more exothermic is depicted in Figure 6.3b. Typically, the acceptor molecule has several other higher lying electronic excited states. As the energy transfer reaction becomes more exothermic the donor's triplet level may overlap with these higher lying acceptor states; thus, the donor's energy may be transferred into these states and cascade down through established

- Figure 6.3**
- (a) Donor and acceptor energy level schematic for endothermic energy transfer.
 - (b) Donor and acceptor energy level schematic for exothermic energy transfer. (Inverted region is not observed.)
 - (c) Donor and acceptor energy level schematic for exothermic energy transfer. (Inverted region is observed.)







deactivation pathways to the lowest triplet excited state of the acceptor. When this occurs, the anticipated inverted region is not observed. If, however, the situation arises in which the acceptor's additional electronic levels lie much higher in energy and are unable to accept the donor's excited state energy (see figure 6.3c), then the energy transfer rate constants should follow theory and the predicted decrease should be observed. It is this latter scenario which is proposed to be operative in the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ system.

The energy transfer inverted region has been observed by a few researches in fluid solution, most notably by Engel and coworkers.^{19,20} Their studies involved singlet-singlet energy transfer from various organic donors to azoalkane acceptors. In these systems, the exchange interaction was found to be the dominant energy transfer mechanism, and, a similar "window" or energy gap between the lowest electronic level and the next higher lying level was observed for the azoalkane acceptors.

In order to assert that the observed drop off in energy transfer rate is due to inherent properties of the acceptor molecule, $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$, and not the nature of the donor molecules, information about energy transfer experiments of the donors with anthracene as the acceptor^{21,22} was gathered and is presented in Table 6.2 and Figure 6.4. Clearly these data show the usual saturation of rates as the energy transfer reaction becomes more exothermic. Therefore, the marked decrease in energy transfer rate constants in the inverted region is attributable to the inherent characteristic of the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ molecule.

Table 6.2

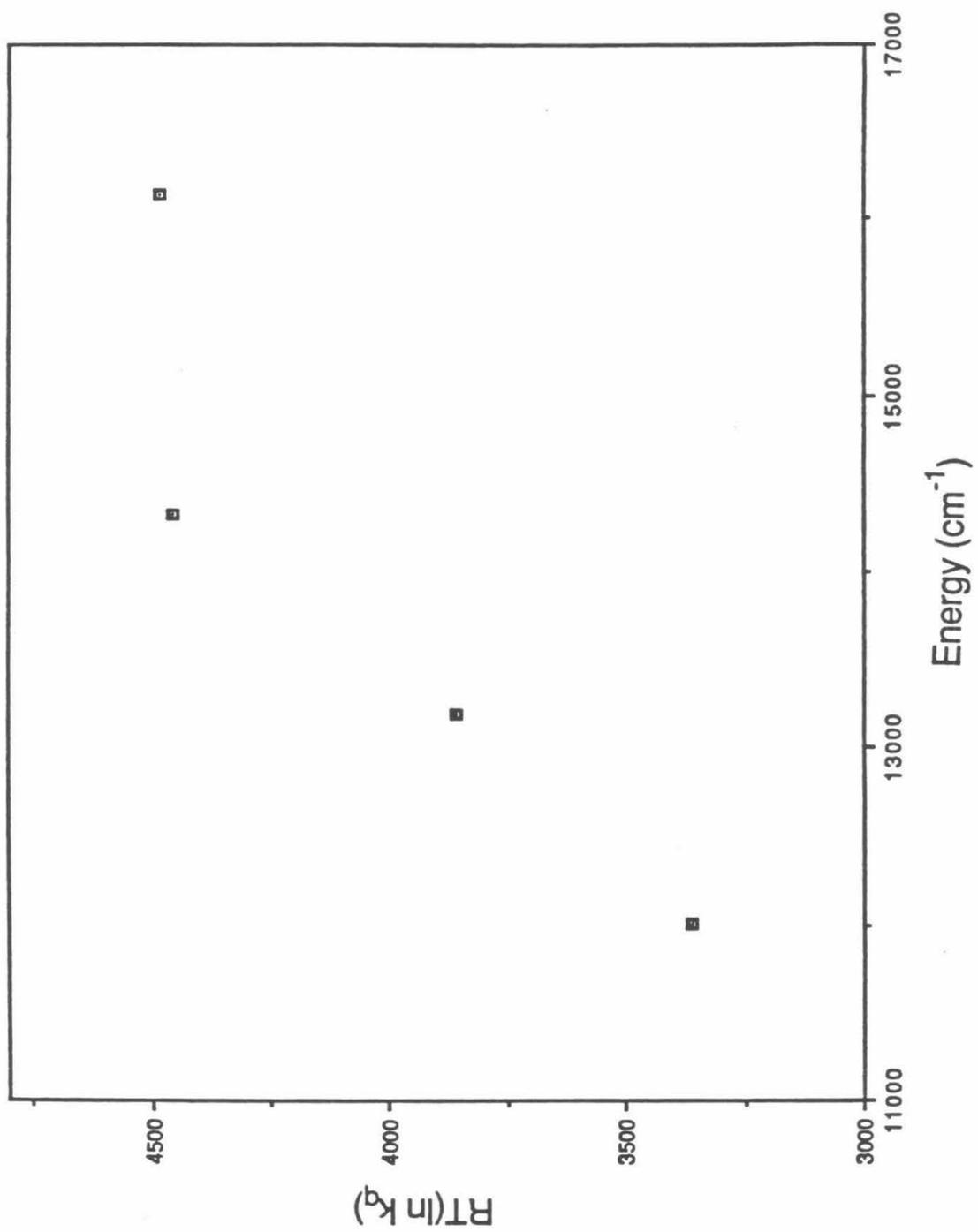
Energy Transfer Data Using Anthracene as the Acceptor[†]

Donor	Emission Maximum		k_q	$RT/\ln k_q$
	nm	cm^{-1}		
$\text{Ru}(\text{bpy})_3^{2+}$	610	16393	2.2×10^9	4485.0
$\text{W}_6\text{I}_{14}^{2-}$	698	14327	1.9×10^9	4454.4
$\text{W}_6\text{Br}_{14}^{2-}$	752	13298	1.1×10^8	3860.4
$\text{W}_6\text{Cl}_{14}^{2-}$	833	12005	1.0×10^7	3360.5

[†]Data taken from references 21 and 22.

Figure 6.4

Plot of the energy transfer data between donors ($W_6X_8Y_6^{2-}$, $Ru(bpy)_3^{2+}$) and the acceptor, anthracene. (Data taken from references 21 and 22.)



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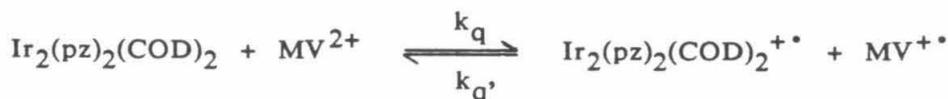
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CHAPTER 7**Electron Transfer Studies of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ and Its Analogs**

INTRODUCTION

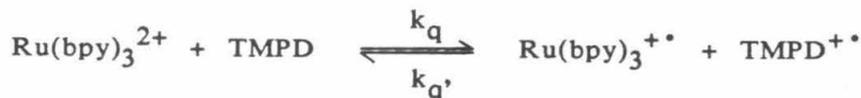
Oxidation and reduction reactions are common in inorganic chemistry; thousands of examples are known. All redox reactions occur by either an inner-sphere (the complexes share one ligand in their coordination environment) or an outer-sphere (the complexes have separated coordination environments) electron transfer mechanism.¹ The electron transfer mechanism is also one of the more common excited state deactivation processes observed in inorganic photochemistry.²⁻⁴ Rates of inner-sphere electron transfer are frequently limited by the rate of ligand substitution and are unlikely to be competitive with excited state decay rates.² Therefore, most of the excited state redox reactions occur by the outer-sphere electron transfer mechanism. Two examples of this type of reactivity are shown below.^{5,6}



$$\lambda_{em} = 684 \text{ nm}; \tau = 250 \text{ ns}$$

$$k_q = 8.7 \times 10^9 \text{ M}^{-1} \text{ sec}^{-1}; k_{q'} = 1.4 \times 10^9 \text{ M}^{-1} \text{ sec}^{-1}$$

and



$$\lambda_{em} = 610 \text{ nm}; \tau = 600 \text{ ns}$$

$$k_q = 1.2 \times 10^{10} \text{ M}^{-1} \text{ sec}^{-1}; k_{q'} = 2.8 \times 10^{10} \text{ M}^{-1} \text{ sec}^{-1}$$

(where pz = pyrazine, COD = cyclooctadiene, MV²⁺ = methylviologen, bpy = 2,2'-bipyridine, and TMPD = N,N,N',N'-tetramethyl-1,4-phenylenediamine). These transition metal complexes luminesce in room temperature, fluid solution and have excited state lifetimes on the order of 10⁻⁷ sec.

A few examples of transition metal complexes which do not luminesce

in fluid solution, but which have long-lived excited states (in fluid solution) are known. Keep in mind, that the lack of fluid solution luminescence does not rule out the possibility of excited state electron transfer reactivity. But rather, since luminescence is used as a tool for monitoring excited state electron transfer reactions, complexes which do not luminesce rarely have their excited state electron transfer reactivity investigated.

This chapter describes experiments designed to investigate the potential thermal and photochemical electron transfer reactivity of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$, which does not luminesce in fluid solution, and its analogs.

EXPERIMENTAL

Materials

All solvents, unless otherwise noted, were reagent grade and used without further purification. The dichloromethane (Baker) used in the electrochemical experiments was dried over calcium hydride and freshly distilled prior to use. The dichloromethane (Burdick and Jackson, UV grade) used in the electronic absorption and flash photolysis studies was also dried over calcium hydride and stored over Linde 3A molecular sieves in a storage flask as described in the experimental section of Chapter 2. The acetonitrile (Burdick and Jackson, UV grade) used in the flash photolysis studies was dried over calcium hydride and stored over alumina in a storage flask as outlined in the experimental section of Chapter 2.

The platinum complexes, $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 \cdot 0.5\text{CH}_2\text{Cl}_2$ and $[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_2\text{Me})_2^{2+}](\text{PF}_6)_2$, were prepared as described in the experimental section of Chapter 5.

Tetracyanoethylene, TCNE, (Aldrich, 98%) was recrystallized twice from chlorobenzene and sublimated under vacuum. White crystalline chunks were obtained.

Tetracyanoquinodimethane, TCNQ, (Aldrich, 98%) was recrystallized twice from chlorobenzene and dried *in vacuo*. This complex darkens after several weeks in air; therefore, recrystallization and drying was carried out immediately prior to the EPR experiments.

Methylviologen hexafluorophosphate, MV^{2+} , and tetrabutylammonium hexafluorophosphate, TBAPF_6 , were prepared and purified as outlined in the experimental section of Chapter 2.

Physical Measurements

Electrochemical Measurements. All cyclic voltammograms were measured in dichloromethane using the same equipment and operating procedures as detailed in the experimental section of Chapter 2. Several sets of electrochemical experiments were done. First, cyclic voltammograms were measured of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ and $[\text{Pt}_2(\mu\text{-dppm})_2(\text{L})_2(\text{PF}_6)_2]$ where $\text{L} = \text{PPh}_3, \text{PPh}_2\text{Me},$ and PPhMe_2 . Next, bulk electrolysis of a dichloromethane solution of 1×10^{-3} M TCNE and 0.1 M TBAPF_6 was performed. And, finally, cyclic voltammograms monitoring the progress of the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 + \text{TCNE}$ thermal reaction were measured in conjunction with EPR experiments. These samples consisted of an N_2 purged dichloromethane solution of 1×10^{-3} M $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$, 1×10^{-3} M TCNE, and 0.1 M TBAPF_6 .

Electronic Absorption Spectroscopy. Spectra were recorded using the Shimadzu UV-260 spectrophotometer. The $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 + \text{TCNE}$ dark reaction was run in a special photometric cell which consisted of a central square precision quartz cuvette (1 cm pathlength), flanked on either side by a pyrex round bulb (2 in total), which were separated from the cuvette by teflon vacuum valves. A third teflon vacuum valve atop the photometric cell was used to keep the entire vessel under vacuum. The platinum complex and TCNE were placed in separate bulbs and the entire cell was evacuated on a high vacuum line (limiting pressure $<10^{-3}$ Torr). Under vacuum, ~ 4 mL of dichloromethane was distilled from the solvent pot into the bulb containing the platinum complex. This solution was transferred to the cuvette and an absorption spectrum was measured. Subsequently, some TCNE from the second bulb was added to the platinum solution in the cuvette and the progress of the ensuing

dark reaction was monitored by successive absorption scans.

Electron Paramagnetic Resonance Spectroscopy. All EPR spectra were measured on a Varian E-line Century Series spectrometer at 5 mW power. EPR spectra run in conjunction with electrochemical measurements were prepared under N_2 by syringing 0.5 mL of the electrochemical solution (described under the electrochemical measurements) into a N_2 purged EPR tube capped with a septum. All other EPR samples were prepared by mixing 1:1 dichloromethane solutions of 1×10^{-3} M or 5×10^{-4} M $Pt_2(\mu\text{-dppm})_2Cl_2$ with 1×10^{-3} M or 5×10^{-4} M TCNE (v/v) together in a graduated cylinder and syringing 0.5 mL into a quartz EPR tube. Most of these samples were prepared in air; however, several $Pt_2(\mu\text{-dppm})_2Cl_2$ + TCNE samples were prepared under N_2 or degassed completely. No significant differences between the EPR spectra of samples prepared in air, under nitrogen, or in degassed solutions were observed. Similar procedures were used for preparing $Pt_2(\mu\text{-dppm})_2(PPh_2Me)_2^{2+}$ + TCNE and $Pt_2(\mu\text{-dppm})_2Cl_2$ + TCNQ samples.

Microsecond Flash Photolysis Experiments. Flash photolysis studies were conducted on an apparatus constructed at Caltech and described elsewhere.^{7,8} Flash lamp excitation wavelengths greater than 350 nm were used. Absorption changes were monitored at 360 and 605 nm. A 1:1 $CH_2Cl_2:CH_3CN$ (v/v) mixed solvent containing $\sim 1 \times 10^{-4}$ M $Pt_2(\mu\text{-dppm})_2Cl_2$ and 5×10^{-3} M MV^{2+} was prepared in a flash photolysis cell consisting of a cylindrical pyrex compartment (15 cm pathlength), a 50 mL pyrex bulb, and a square cuvette (1 cm pathlength), each oriented 90° from one another. A set of control experiments to check for MV^{2+} and solvent contamination were done. The transient absorption at 605 nm of a CH_3CN solution of 5×10^{-3} M MV^{2+} and a 1:1 $CH_2Cl_2:CH_3CN$ (v/v) mixed

solvent of 5×10^{-3} M MV^{2+} were measured; however, in both cases no $MV^{+ \bullet}$ was detected.

Steady State Photolysis Procedures. A 1000 W high pressure Hg/Xe arc lamp equipped with Corning cutoff filters ($\lambda > 350$ nm) was used for the irradiations. Samples were prepared by placing $Pt_2(\mu\text{-dppm})_2Cl_2$ (~1 mg) in a two compartment spectrophotometric cell equipped with a teflon vacuum valve and evacuating on a high pressure vacuum line (limiting pressure $< 10^{-3}$ Torr). Dichloromethane was distilled from the storage flask into the spectrophotometric cell under vacuum. The dichloromethane solutions of $Pt_2(\mu\text{-dppm})_2Cl_2$ (various concentrations) were photolyzed over a period of 3 days. No electronic absorption spectral changes were observed over the course of the photolysis; hence, $Pt_2(\mu\text{-dppm})_2Cl_2$ is considered to be photo-stable in dichloromethane.

RESULTS and DISCUSSION

The redox properties of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ and $\text{Pt}_2(\mu\text{-dppm})_2\text{L}_2$ have been investigated and the results are presented in Table 7.1. The $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ complex exhibits an irreversible oxidation +0.85 V vs. SCE. Similar studies performed by Nemra *et al.*⁹ on $\text{Pd}_2(\mu\text{-dppm})_2\text{Cl}_2$ and $\text{PtPd}(\mu\text{-dppm})_2\text{Cl}_2$ in 1,2-dichloroethane show irreversible oxidations at +0.98 V and +0.96 V vs. SCE, respectively. In the Pd(I) and mixed Pd(I)Pt(I) dimers, the oxidation occurs by way of a one electron process; this same process is believed to occur in the platinum system. No return peak on the reverse is observed in the cyclic voltammogram at scans rates up to 500 mV/sec; therefore, the Pt(I,II) species produced must be fairly unstable, possibly reacting with the solvent.

No reduction wave was observed as far out as -1.5 V vs. SCE for the platinum(I) dimer. Irreversible two electron reductions were observed for $\text{Pd}_2(\mu\text{-dppm})_2\text{Cl}_2$ and $\text{PtPd}(\mu\text{-dppm})_2\text{Cl}_2$ at -1.26 V and -1.70 V vs. SCE, respectively.⁹ From these results, a similar two electron irreversible reduction of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ is expected to occur at a potential more negative than -1.70 V vs. SCE. Thus, for the most part, any electron transfer chemistry observed with the platinum(I) dimers will most likely involve a one electron oxidation of the complex.

Thermal and photochemical reactions designed to probe the electron transfer reactivity of these platinum complexes with various organic acceptors were undertaken. Tetracyanoethylene, TCNE, is very easily reduced, with a one electron reduction potential of +0.15 V vs. SCE.¹⁰ Owing to the pronounced energy difference between neutral TCNE and TCNE^{2-} , the $\text{TCNE}^{\cdot -}$ has a high solution stability, making it an ex-

Table 7.1

Oxidation Potentials of $[\text{Pt}_2(\mu\text{-dppm})_2\text{Y}_2]^{n+}$

Complex	$-E_{1/2}$ vs. SCE
$\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$	+0.85 V
$[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_3)_2](\text{PF}_6)_2$	+1.3 V
$[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_2\text{Me})_2](\text{PF}_6)_2$	+1.3 V
$[\text{Pt}_2(\mu\text{-dppm})_2(\text{PPhMe}_2)_2](\text{PF}_6)_2$	+1.35 V

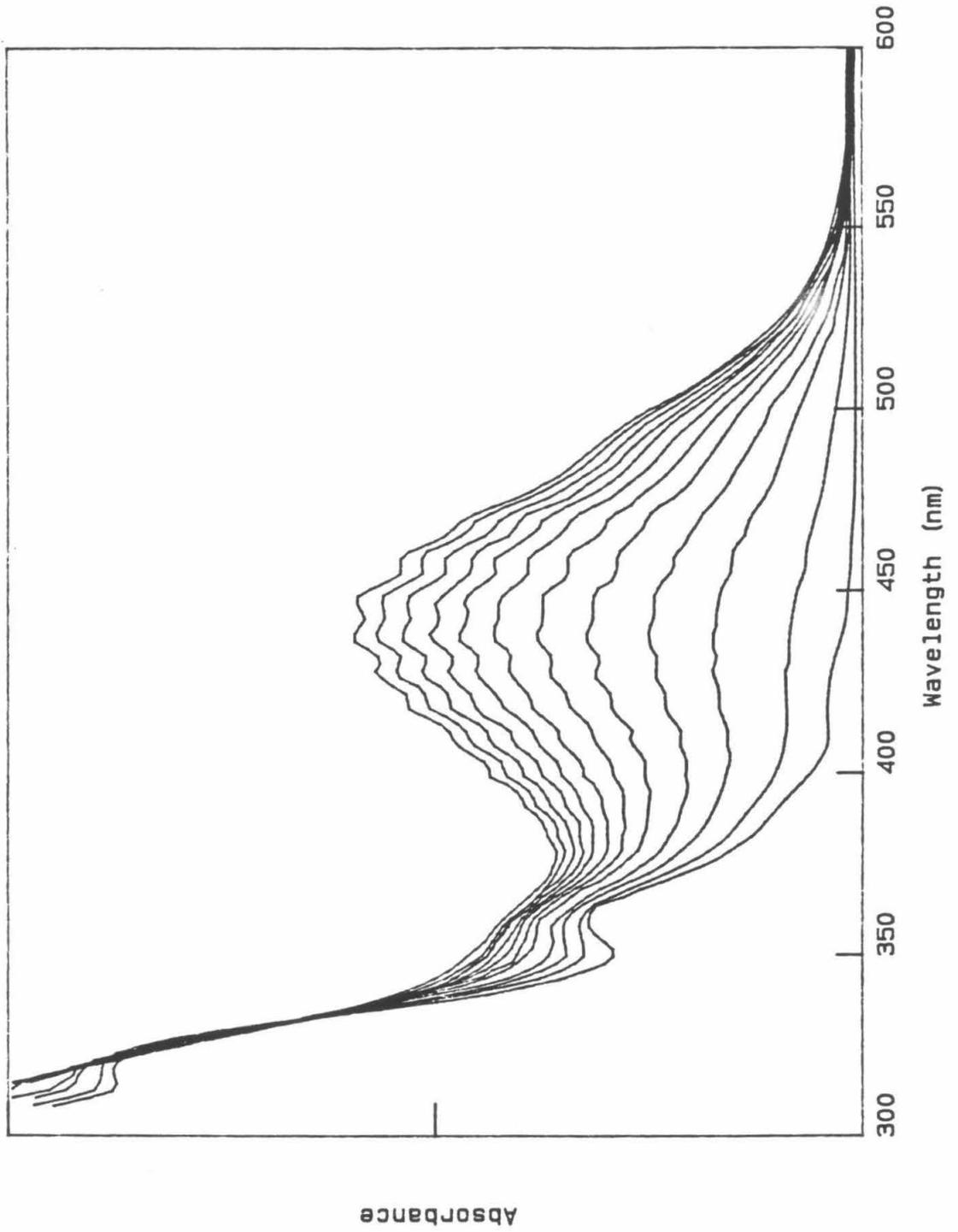
cellent one electron acceptor.¹¹ Upon mixing $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}$ with TCNE in deaerated dichloromethane, $\text{TCNE}^{\cdot-}$ was immediately produced, even in reactions where room light was rigorously excluded. Successive scans of the electronic absorption spectrum of the reaction mixture (Figure 7.1) clearly show an absorption band with the vibrational fine structure characteristic of the $\text{TCNE}^{\cdot-}$ growing in at ~ 440 nm.

Although ground state electron transfer from the platinum complex to TCNE should not occur based on ground state redox potentials, some type of adduct formation may occur which facilitates the electron transfer. Considering that several examples of small molecules adding thermally across the metal-metal bond of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ are known to occur, the prospect for a $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2\text{-TCNE}$ adduct forming prior to the electron transfer step is quite reasonable. Further oxidation of the platinum complex to produce a stable $d^8\text{-}d^8$ platinum dimer containing a tetracyanoethane bridge does not seem to occur due to the high stability of the $\text{TCNE}^{\cdot-}$.

Evidence for the proposed $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2^{\cdot+} / \text{TCNE}^{\cdot-}$ adduct is obtained from EPR spectra. The EPR spectrum of the electrochemically generated $\text{TCNE}^{\cdot-}$ in dichloromethane (tetrabutylammonium hexafluorophosphate, TBAPF_6 , present as electrolyte) and spectrum of the $\text{TCNE}^{\cdot-}$ generated by reduction with $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ in dichloromethane (no TBAPF_6 present) have several similarities, but, the overall shapes of the spectra are different.

The EPR spectrum of the electrochemically generated $\text{TCNE}^{\cdot-}$ is shown in Figure 7.2a. The EPR lines are significantly broadened from spectra previously reported for the $\text{TCNE}^{\cdot-}$.¹² This broadening is due to self-exchange between neutral TCNE and the radical facilitated by the

Figure 7.1 Electronic absorption spectrum of the thermal electron transfer reaction between $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ and TCNE in dichloromethane (room light excluded).



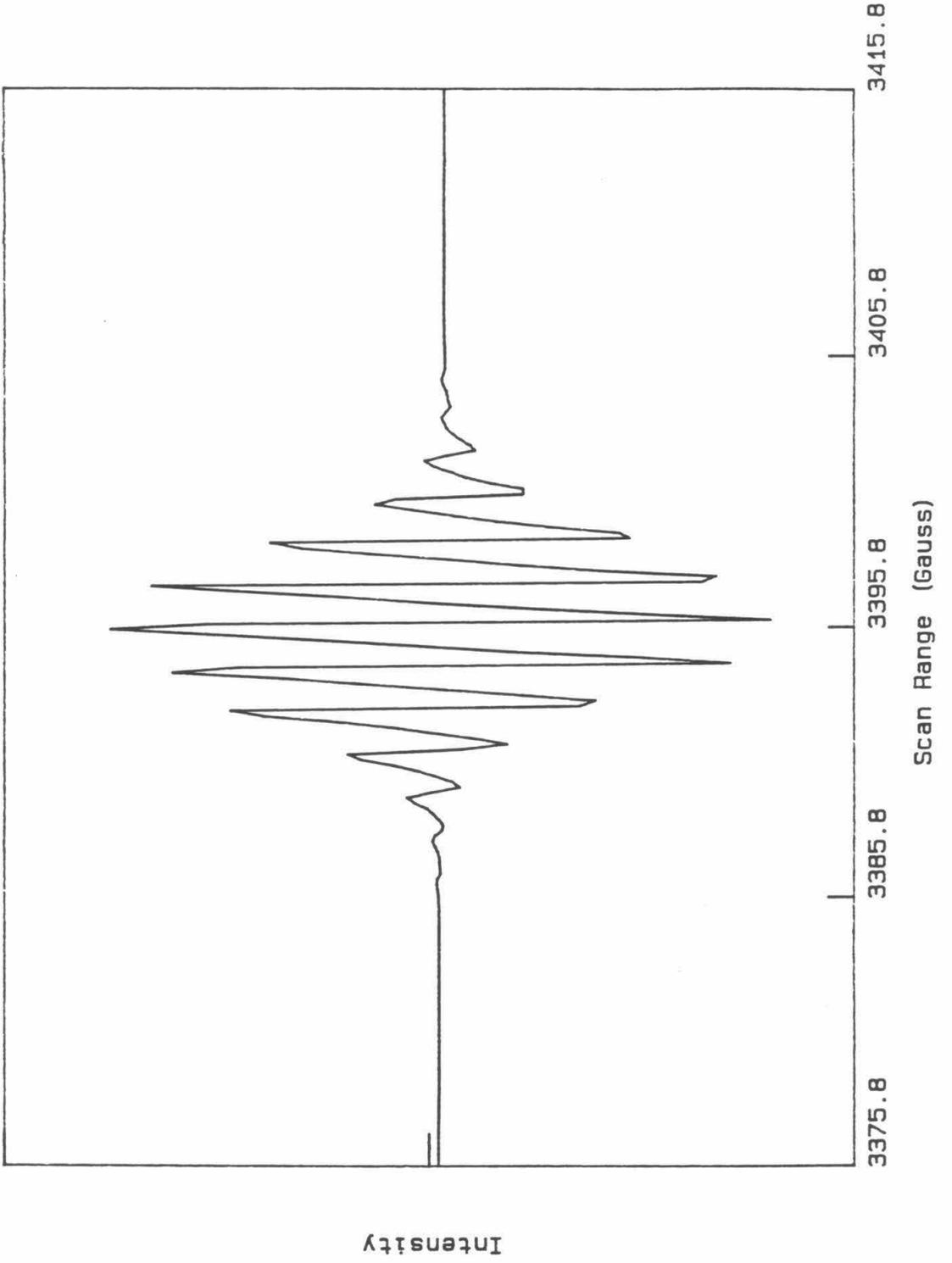
presence of the supporting electrolyte, TBAPF₆ (required for bulk electrolysis).¹³ The uncoordinated TCNE^{-•} spectrum exhibits a total of 13 lines. The nine predominant lines (center of the spectrum) are expected for the four equivalent ¹⁴N (I=1) nuclei; whereas, the smaller lines found at the beginning and end of the spectrum are best attributed to ¹³C nuclei.¹² The spectrum has a g value of 2.001 and the spacing is 1.56 Hz; both results are in excellent agreement with previously reported results for TCNE^{-•}.¹²

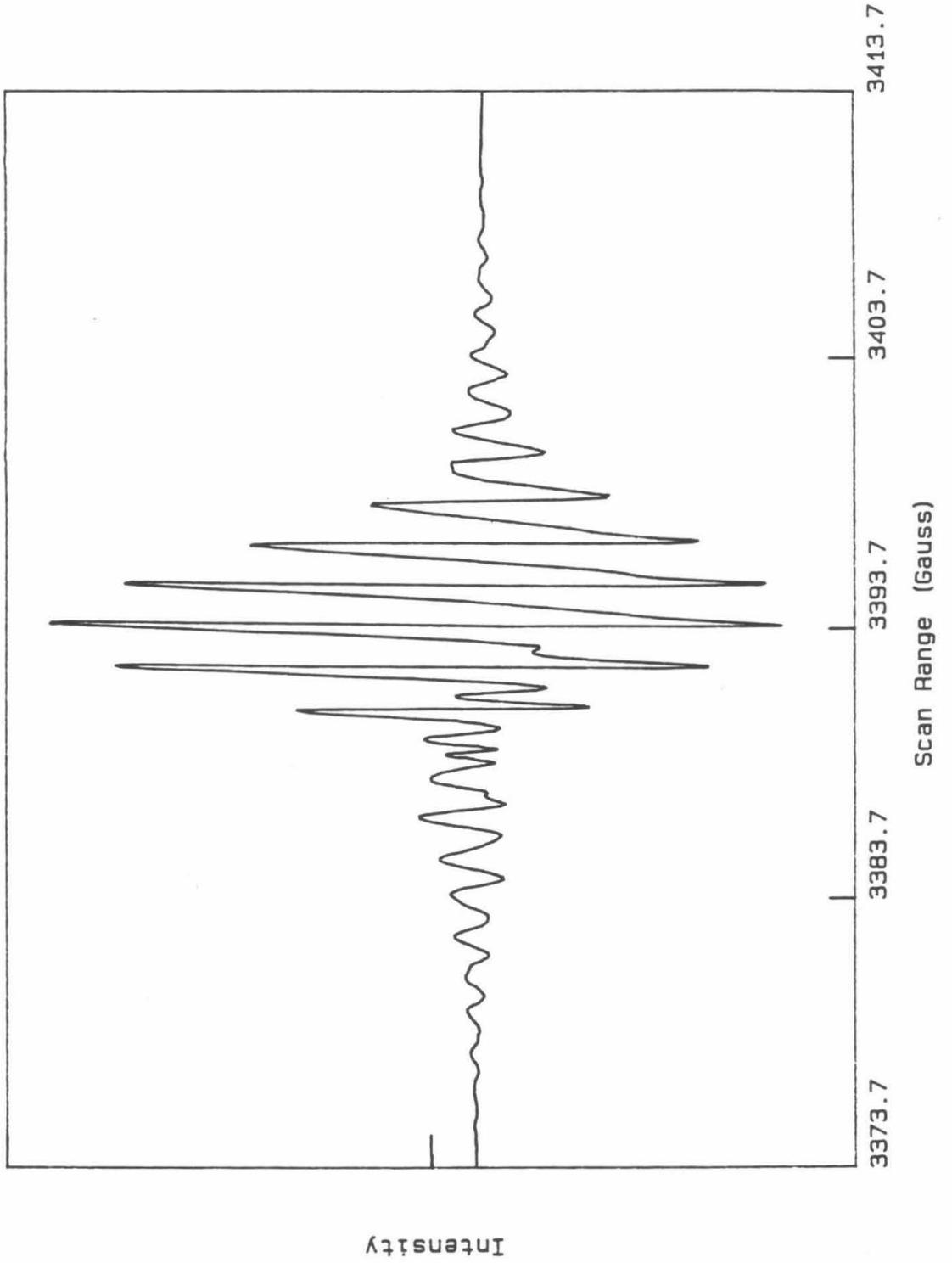
As for the EPR spectrum obtained when TCNE is reduced with Pt₂(μ-dppm)₂Cl₂, a similar g value of 2.000 and a range of line spacings 1.52 to 1.56 Hz is observed, see Figure 7.2b. From these results it is evident that the radical generated from this reaction does indeed resemble that of TCNE^{-•}. However, in this spectrum, ~27 lines are observable. Since the total number of lines observed far exceeds that expected for uncoordinated TCNE^{-•} radical, it is presumed that the electron primarily localized on the TCNE also feels the presence of other nuclei in the vicinity such as ³¹P (I = 1/2) and ¹⁹⁵Pt (I = 1/2). It is the observation of all these additional EPR lines which supports the supposition that some type of Pt₂(μ-dppm)₂Cl₂-TCNE adduct formation occurs to facilitate the electron transfer.

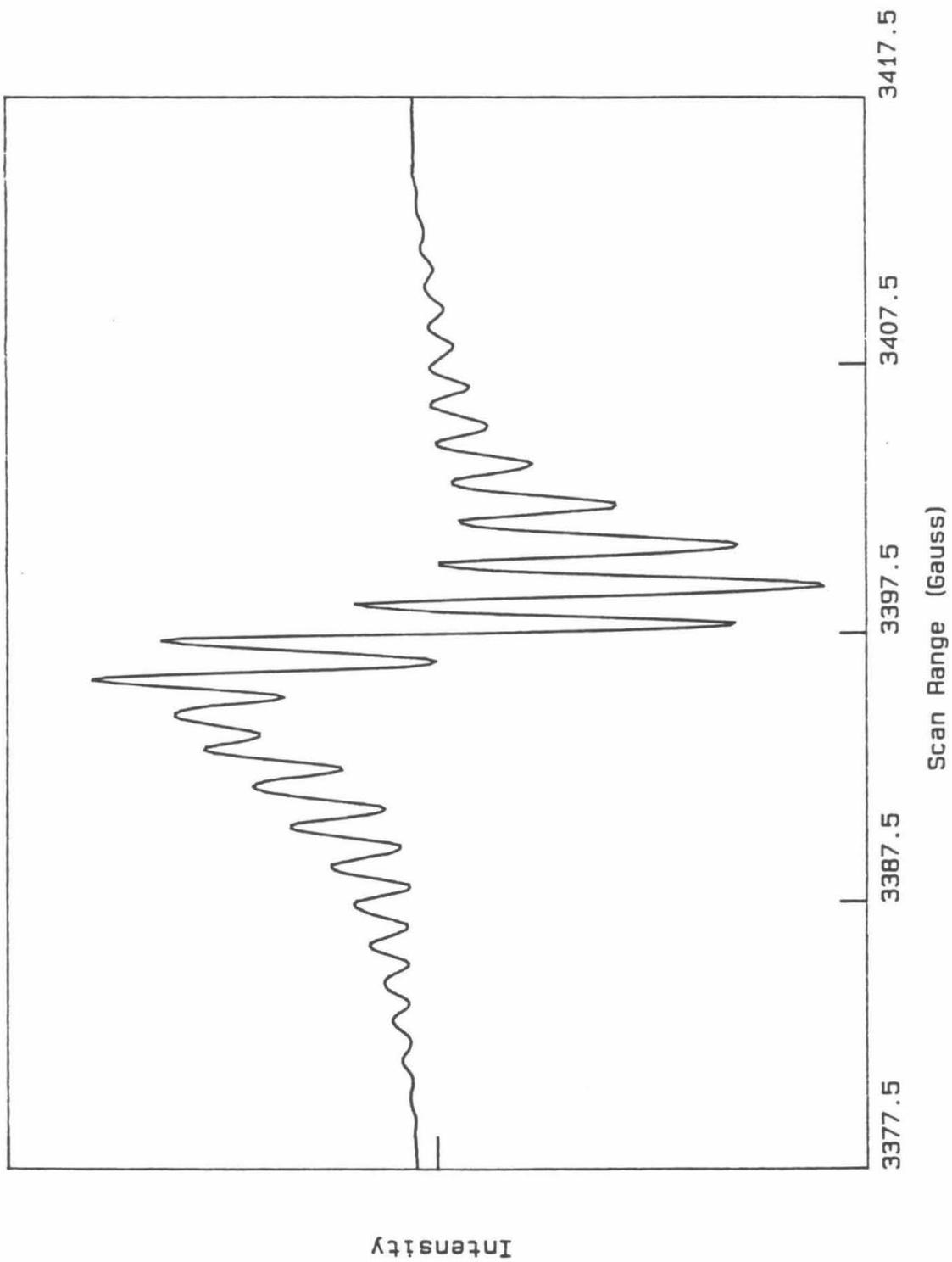
An EPR signal attributable to the oxidized platinum species was not observed. Since it is not uncommon for transition metal complexes to have spin lattice relaxation times which are too short to allow for signal detection at room temperature,¹⁴ the absence of the oxidized platinum complex's signal is not unusual.

The overall shape of the EPR spectrum obtaining from the Pt₂(μ-dppm)₂Cl₂ + TCNE reaction differs from reaction mixture to reaction

- Figure 7.2**
- (a) EPR spectrum of electrochemically generated $\text{TCNE}^{\cdot-}$ in dichloromethane (0.1 M TBAPF_6 , 300 K).
 - (b) EPR spectrum of the thermal electron transfer reaction between $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ and TCNE in dichloromethane (300 K).
 - (c) A second EPR spectrum of the thermal electron transfer reaction between $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ and TCNE in dichloromethane (300 K).







mixture, although about the same number of EPR lines are observed and the g value and range of line spacings are identical in all cases. (Examples of these differences may be seen by comparing Figures 7.2b and 7.2c.) The difference in shape, along with the general complexity of the EPR spectrum itself, seems to indicate that the reaction mixture may contain several structural isomers of the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2\text{-TCNE}$ adduct.

Similar thermal electron transfer experiments between $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ and 7,7,8,8-tetracyanoquinodimethane, TCNQ were also conducted. TCNQ, like TCNE, is easy to reduce, $E_{1/2} = +0.127$ V vs. SCE,¹⁵ but unlike TCNE, TCNQ is a much bulkier electron acceptor and may not allow similar adduct formation to occur. Upon mixing the platinum complex with TCNQ a very complicated radical signal was observed. However, since the EPR spectrum of uncoordinated $\text{TCNQ}^{\cdot-}$ is reported to contain 43 lines, information regarding $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2\text{-TCNQ}$ adduct formation was not forthcoming.

Surprisingly, reaction of $\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_2\text{Me})_2^{2+}$ with TCNE did not result in the formation of any radical species. There are two explanations for this lack of reactivity: (i) the phosphine groups in the axial positions of $\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_2\text{Me})_2^{2+}$ constricts access to the metal-metal bond, thus making adduct formation much more difficult than in the halide analog, and (ii) $\text{Pt}_2(\mu\text{-dppm})_2(\text{PPh}_2\text{Me})_2^{2+}$ is -45 mV more difficult to oxidize than $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$.

In order to obtain additional information about the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 + \text{TCNE}$ thermal reaction, its progress was simultaneously monitored by cyclic voltammetry and EPR over a 14 hour period. An N_2 purged

dichloromethane solution of 1×10^{-3} M $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$, 1×10^{-3} M TCNE, and 0.1 M TBAPF_6 was prepared; 11 mL of the solutions was placed in the electrochemical cell and 0.5 mL was placed in a septum-sealed quartz EPR tube.

At first only the TCNE reduction and $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ oxidation waves were observed in the cyclic voltammogram, but, within a few minutes a new quasi-reversible wave at $\sim +0.3$ V vs. SCE and an irreversible oxidation wave at $\sim +1.2$ V vs. SCE appeared. After about 2 and a half hours the quasi-reversible wave at $\sim +0.3$ V vs. SCE disappeared. And after 14 hours, the reduction wave of TCNE was substantially diminished (only 10% of the original signal remained) and only the irreversible oxidation wave at $+1.2$ V vs. SCE was observed.

The initial EPR spectrum obtained resembled the EPR spectra previously obtained for the thermal $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 + \text{TCNE}$ reaction (without electrolyte). However, as the reaction progressed, some of the EPR lines began to disappear and all the remaining lines began to sharpen. By the end of the experiment (14 hours later), the EPR spectrum consisted of 13 rather sharp lines, remarkably resembling uncoordinated $\text{TCNE}^{\cdot -}$. The results of these experiments are puzzling, especially in light of the fact that the EPR spectra of dichloromethane solutions of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2 + \text{TCNE}$ without electrolyte present do not show this unusual behavior with time. Naturally, these results raise questions about the role of the TBAPF_6 in this reaction.

One possibility is that the electrolyte facilitates the further oxidation of the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2^{\cdot +} / \text{TCNE}^{\cdot -}$ adduct to produce a $d^8\text{-}d^8$ platinum dimer containing a tetracyanoethane bridge. Since the reaction was run at a 1:1 stoichiometry and only ~ 90 % of the TCNE was

consumed during the course of the reaction, possible platinum dimer decomposition may have occurred resulting in uncoordinated TCNE existing in its radical state. With this in mind, the transient quasi-reversible wave at $\sim +0.3$ V vs. SCE may be attributable to the $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2^{+\bullet}$ / $\text{TCNE}^{-\bullet}$ adduct, the irreversible oxidation at ~ 1.2 V vs. SCE may be attributable to the $d^8\text{-}d^8$ platinum dimer (which is EPR silent), and the residual $\text{TCNE}^{-\bullet}$ (whose oxidation wave is observed in the final cyclic voltammogram) remains as the only species observed in the final EPR spectrum.

Further electron transfer experiments between $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ and methylviologen hexafluorophosphate, MV^{2+} , were investigated; the reduction potential of MV^{2+} is -0.45 V vs. SCE.⁶ Unlike in the TCNE experiments, no ground state electron transfer between the platinum complex and MV^{2+} was observed. Flash photolysis of MV^{2+} in the presence of $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$, however, resulted in the production of $\text{MV}^{+\bullet}$. Although a complete transient absorption spectrum of the $\text{MV}^{+\bullet}$ was not collected, an increase of the absorptions at 605 and 360 nm were observed and the characteristic blue color of the $\text{MV}^{+\bullet}$ radical was detected after several repetitive flashes.

Since a mixed solvent system, 1:1 (v/v) $\text{CH}_2\text{Cl}_2:\text{CH}_3\text{CN}$ was used, several control experiments were done to rule out the possibility of MV^{2+} or solvent contamination; see insert in Figure 7.3a. Monitoring at 605 nm, the transient decay signal shown in Figure 7.3a was obtained. These data were evaluated according to the bimolecular rate equation:

$$1/\Delta C = 1/C_t - 1/C_o = kt$$

(where C_t is the concentration of the absorbing species at some time t , and C_0 is the initial concentration). Since the flash photolysis data obtained are in absorption units and $A = \epsilon bC$, the bimolecular rate expression used to calculate the rate constant for back electron transfer, k_b , is modified as:

$$b\Delta\epsilon/\Delta A = b\epsilon_t/A_t - b\epsilon_0/A_0 = k_b t$$

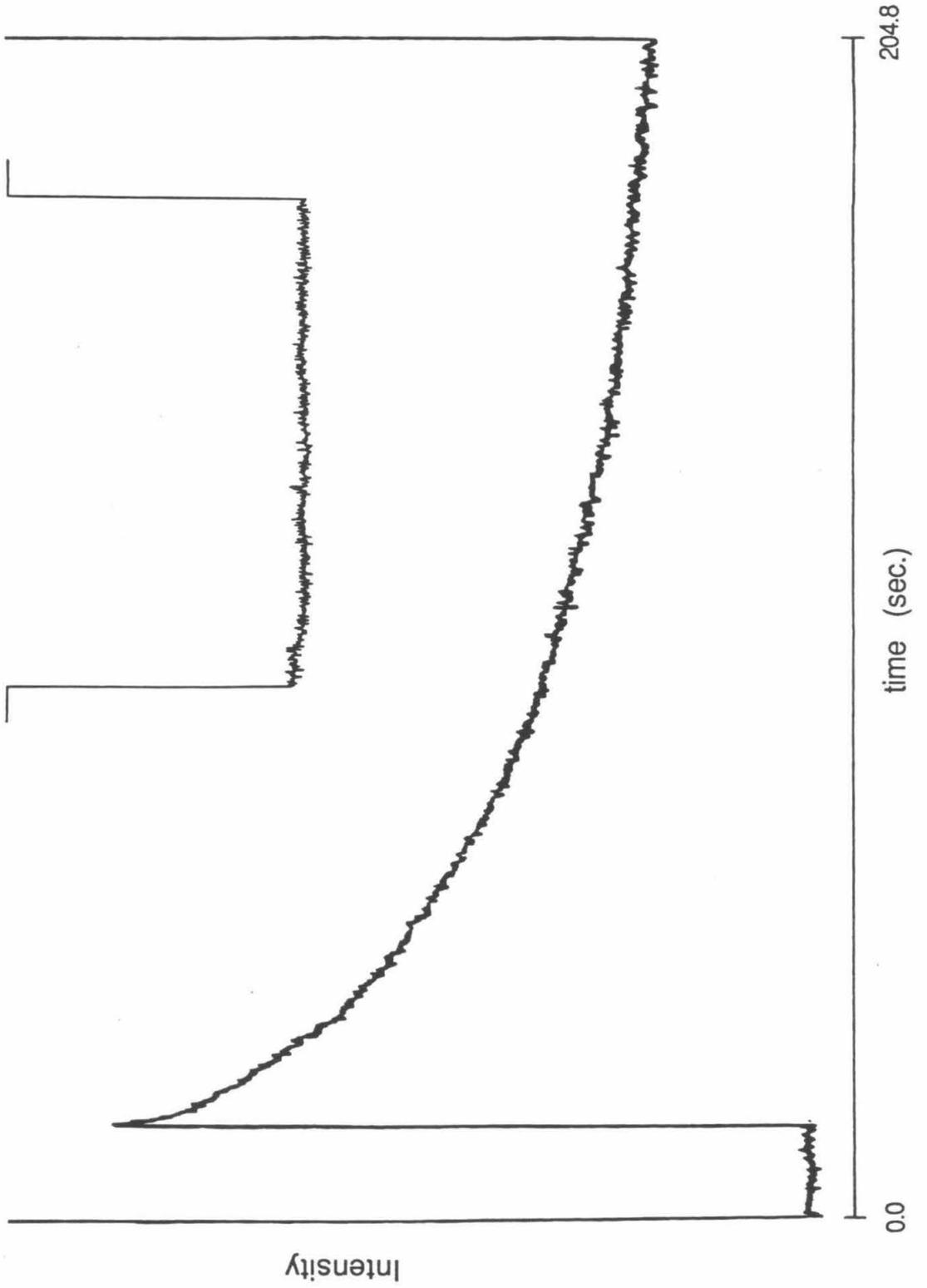
or

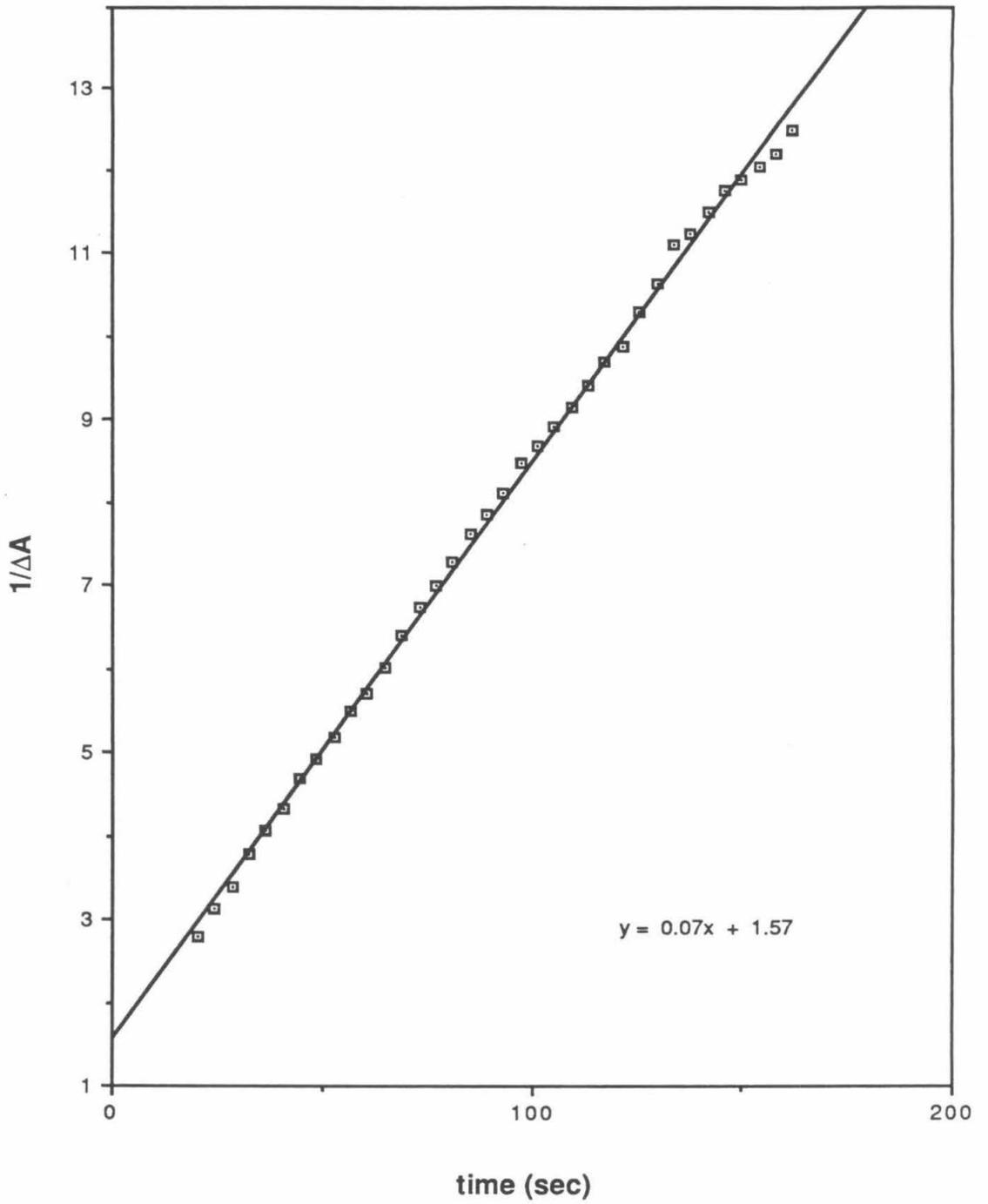
$$1/\Delta A = (k_b/b\Delta\epsilon)t.$$

The plot of $1/\Delta A$ vs. t is shown in Figure 7.3b. Since most of the absorption at 605 nm is due to the $MV^{+\bullet}$, $\Delta\epsilon$ is approximated as $\Delta\epsilon_{MV^{+\bullet}}$, whose value has been experimentally determined by Kosower *et al.*¹⁶ as $13,000 \text{ M}^{-1} \text{ cm}^{-1}$. The cell pathlength, b , is 15 cm and the slope of the $1/\Delta A$ vs. t plot is 0.07 sec^{-1} . Therefore, the back electron transfer rate constant is calculated to be $1.5 (\pm 0.5) \times 10^4 \text{ M}^{-1} \text{ sec}^{-1}$. For stability, the oxidized platinum dimer may coordinate solvent, namely CH_3CN , which could account for the surprisingly slow back electron transfer rate. There is, however, no direct evidence for this proposed solvent coordination.

In conclusion, these experiments have shown that both ground state and excited state electron transfer from $\text{Pt}_2(\mu\text{-dppm})_2\text{Cl}_2$ are possible even though the platinum complex does not luminesce in fluid solution.

- Figure 7.3**
- (a) Flash photolysis trace of the back electron transfer from $MV^{+\bullet}$ to the oxidized $Pt_2(\mu\text{-dppm})_2Cl_2$ in 1:1 dichloromethane:acetonitrile (v/v). [Insert: Flash photolysis trace of MV^{2+} in 1:1 dichloromethane:acetonitrile (v/v).] Absorption monitored at 605 nm.
- (b) Plot of $1/\Delta A$ vs. t (sec) for the back electron transfer from $MV^{+\bullet}$ to the oxidized $Pt_2(\mu\text{-dppm})_2Cl_2$.





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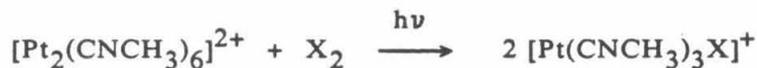
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CHAPTER 8**Preliminary Investigations of $[\text{Pt}_2(\text{CNCH}_3)_6](\text{BF}_4)_2$**

INTRODUCTION

Another class of d^9-d^9 platinum and palladium dimers exist which are unbridged, being supported only by a single metal-metal bond. These complexes, unlike the dppm bridged molecules, have a staggered geometry (see Figure 8.1).¹

The thermal and photochemistry of these unbridged complexes are dominated by homolytic metal-metal bond cleavage reactions producing $15e^-$ radical species which rapidly react with an available substrate or the solvent to form d^8 monomeric complexes. For example:



and



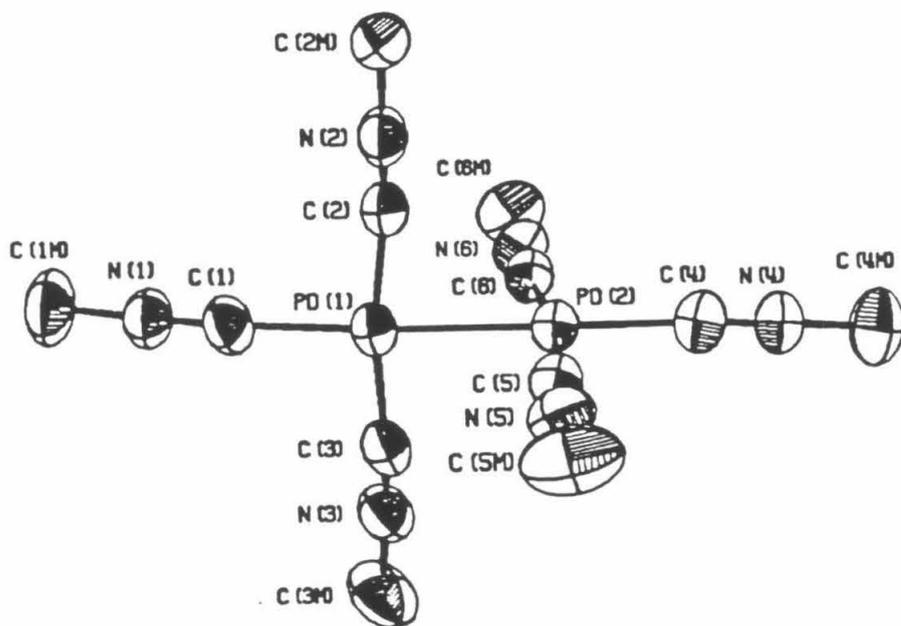
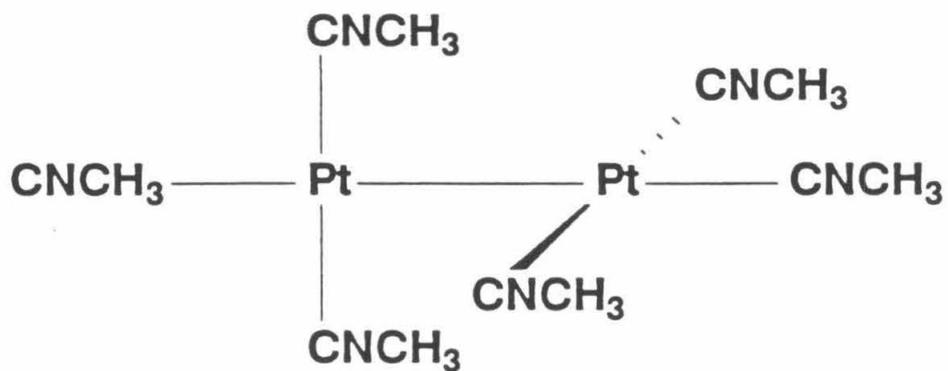
(Such reactions have been investigated by several authors, most notably Balch *et al.*² and Kubiak and coworkers.^{3,4})

Similar single metal-metal bond cleavage reactions are observed with $\text{Mn}_2(\text{CO})_{10}$, which, upon photo-excitation, undergoes homolysis to produce $17e^-$ radical species.^{5,6} Also, reactions of this type are seen with the dppm bridged palladium(I) and platinum(I) dimers which oxidatively add small molecules resulting in net metal-metal bond cleavage;⁷⁻¹³ however, these complexes do not dissociate into monomers due to the presence of the dppm bridges which maintain the dimeric framework.

The $\text{Pt}_2(\text{CNCH}_3)_6^{2+}$ molecule shows great potential as a precursor for synthesizing new d^9-d^9 and d^8-d^8 A-frame dimers. Direct substitution of four CNCH_3 groups with two bridging ligands would afford new d^9-d^9

dimers; whereas, substitution of three CNCH_3 groups by two bridging ligands would produce d^8-d^8 A-frame complexes containing a bridging CNCH_3 group. Preliminary experiments along these lines are described in this chapter.

Figure 8.1 Structure of $[\text{Pt}_2(\text{CNCH}_3)_6](\text{BF}_4)_2$ and ORTEP of $[\text{Pd}_2(\text{CNCH}_3)_6](\text{BF}_4)_2$. (ORTEP taken from reference 1.)



EXPERIMENTAL

Materials

All solvents, unless otherwise noted, were reagent grade and used without further purification.

$[\text{Pt}_2(\text{CNCH}_3)_6](\text{BF}_4)_2$ was prepared by published procedures.² To obtain the best results, it is important that the water solution containing K_2PtCl_4 and NaBF_4 be stirred for 30-45 minutes to attain the maximum NaBF_4 saturation.

Bis(diphenylphosphino)methane, dppm, (Aldrich, 98%) and bis(dimethylphosphino)methane, dmpm, (Aldrich, 98 %) were used as received. 1,3-Diisocyanopropane, $\text{CN}(\text{CH}_2)_3\text{NC}$, was prepared following published procedures,¹⁴ was purified by column chromatography (Alumina, Woelm, Grade A), and gave a clean NMR spectrum.

Physical Measurements

Nuclear Magnetic Resonance. ^1H NMR spectra were recorded at 400 MHz on a JEOL GX-400 spectrometer, at 90 MHz on a Jeol FX90-Q spectrometer, and at 90 MHz on a Varian EM-390 spectrometer. ^1H chemical shifts are reported in ppm (δ) vs. CD_3CN (δ 1.98), CDCl_3 (δ 7.24), CD_2Cl_2 (δ 5.32), and $(\text{CD}_3)_2\text{CO}$ (δ 2.04). ^{31}P NMR spectra were recorded at 36.28 MHz on a JEOL FX90-Q spectrometer. The ^{31}P chemical shifts are reported in ppm (δ) vs. H_3PO_4 external standard. All NMR spectra were run at ambient temperature.

Electronic Absorption Spectroscopy. Absorption spectra were recorded on a Shimadzu 260-UV spectrophotometer.

Emission Spectroscopy and Lifetime Measurements. Emission spectra were measured using an emission instrument constructed at Caltech and

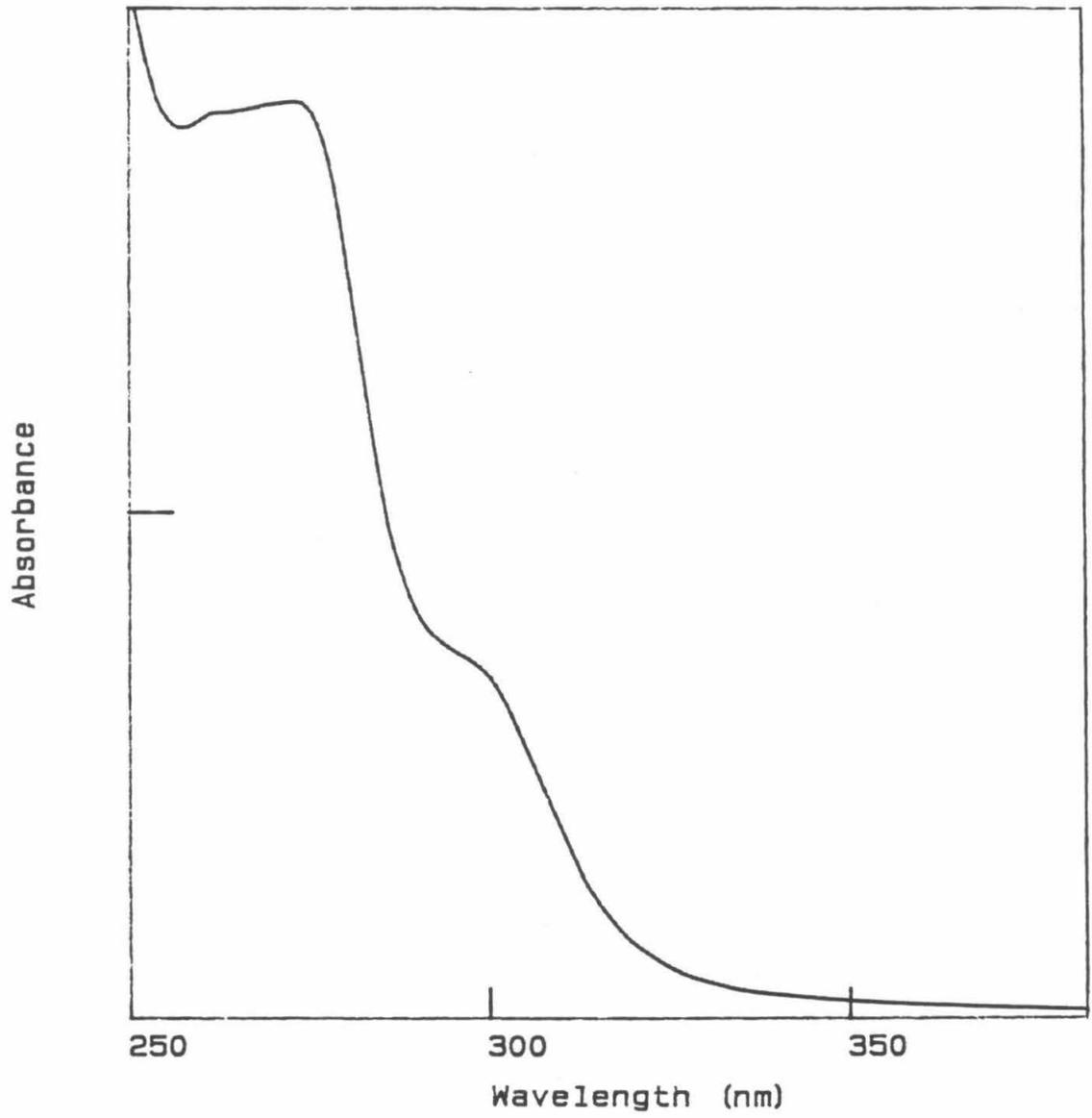
described previously.¹⁵ Emission lifetimes were measured using a Nd:YAG laser system described elsewhere.¹⁶ Low temperature, 77 K, spectra were collected using a liquid-nitrogen-filled quartz finger dewar. Solid samples were placed in an NMR tube and sealed with a septum and parafilm. Acetonitrile solutions of $[\text{Pt}_2(\mu\text{-CN}(\text{CH}_2)_3\text{NC})_2(\mu\text{-CNCH}_3)(\text{CNCH}_3)_2](\text{BF}_4)_2$ were placed in a two-compartment photometric cell equipped with a teflon vacuum valve and degassed (a minimum of 6 freeze/pump/thaw cycles) on a high vacuum line (limiting pressure $< 10^{-3}$ Torr).

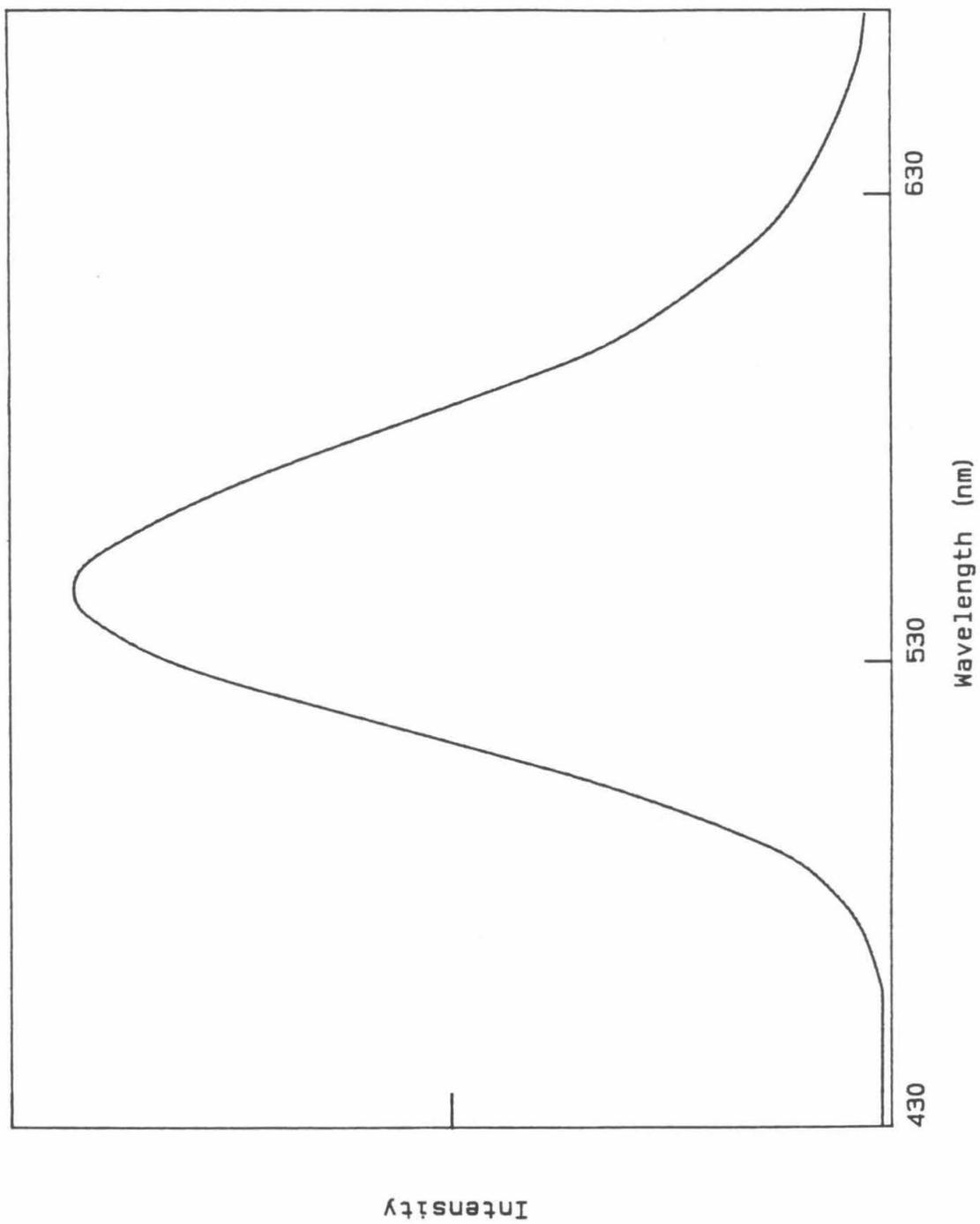
RESULTS and DISCUSSION

The electronic absorption spectrum of $[\text{Pt}_2(\text{CNCH}_3)_6](\text{BF}_4)_2$ reveals a shoulder at 296 nm and a larger broad absorption band maximizing at 272 nm (Figure 8.2a). Reinking *et al.*⁴ report the molar extinction of the 296 nm absorption as $4350 \text{ cm}^{-1} \text{ M}^{-1}$ in acetonitrile. This complex does not luminesce in fluid solution or the solid state at room temperature; however, upon cooling the sample (powder) down to 77 K, yellow emission is observed ($\lambda_{\text{max}} = 542 \text{ nm}$); see Figure 8.2b. This result implies that some facile deactivation pathway exist at room temperature. Since these complexes are known to readily undergo thermal and photochemical homolytic bond cleavage, the rupture of the metal-metal bond may be one of the temperature dependent deactivation pathways.

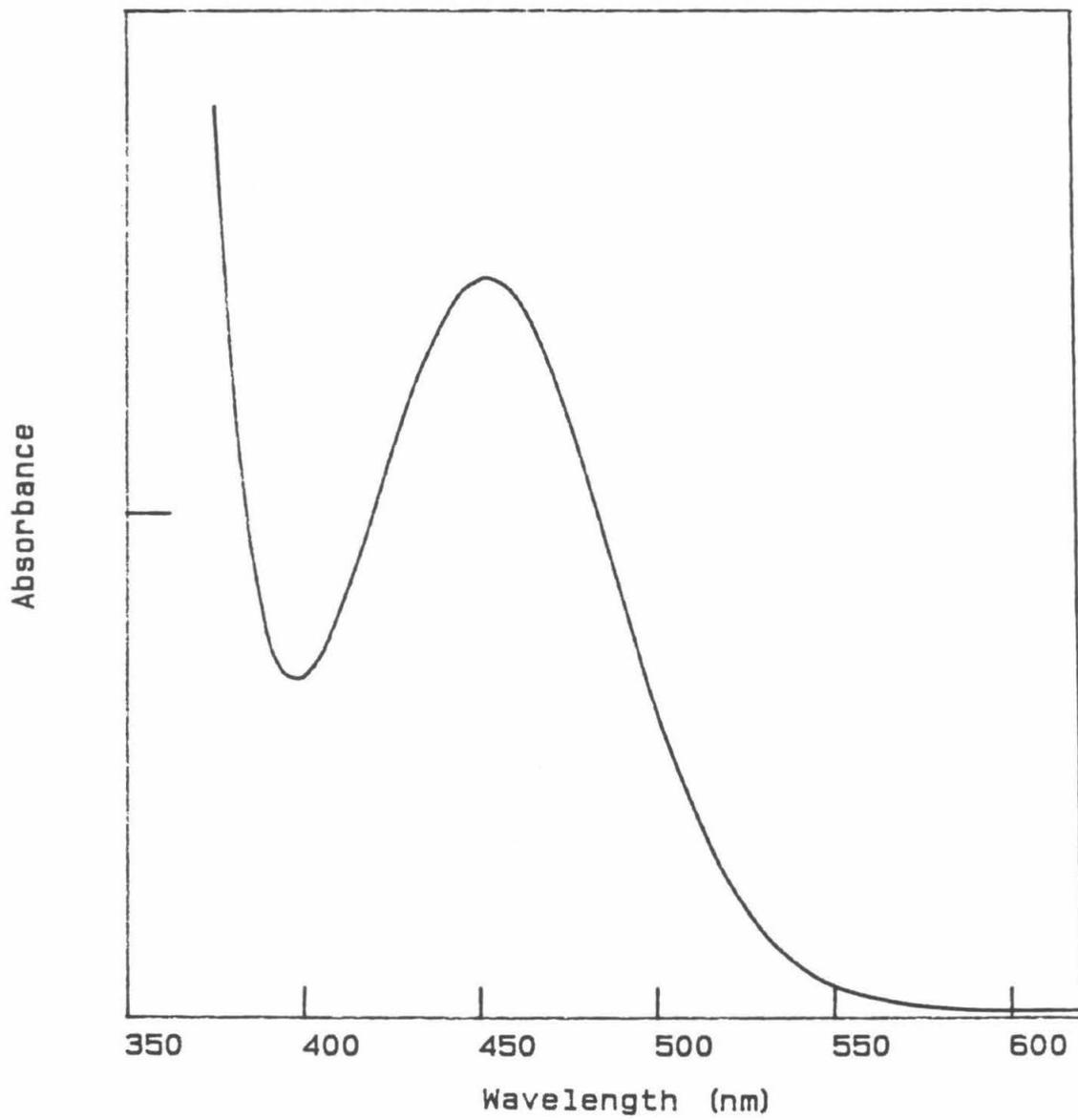
In order to synthesize another class of bridged d^9-d^9 dimers, a colorless acetonitrile solution of $\text{Pt}_2(\text{CNCH}_3)_6^{2+}$ was added to a colorless acetonitrile solution of dppm (1:2 stoichiometry). The reaction was immediate as evidenced by the resulting bright orange solution. The absorption spectrum of this solution shows a new absorption at 452 nm; the absorption spectrum is presented in Figure 8.3a. The solvent was evaporated and an orange solid was obtained. Free CNCH_3 was detected. NMR analysis of the isolated complex clearly revealed the presence of the dppm ligand bridging two platinum centers $\delta = 5.57 \text{ ppm}$; see Figure 8.4. Bands attributable to methylisocyanide were located at $\delta = 2.83$ and 2.50 ppm . Several smaller peaks, unassignable to the starting material or the proposed $\text{Pt}_2(\mu\text{-dppm})_2(\text{CNCH}_3)_n^{2+}$ complex were observed. These peaks are presumed to be decomposition products and d^8 monomeric $\text{Pt}(\text{dppm})(\text{CNCH}_3)_2^{2+}$,¹⁷ from comparison with NMR data of sim-

- Figure 8.2**
- (a) Electronic absorption spectrum of $[\text{Pt}_2(\text{CNCH}_3)_6](\text{BF}_4)_2$ in acetonitrile (300 K). Absorption maxima at 296 (shoulder) and 272 nm.
- (b) Uncorrected emission spectrum of $[\text{Pt}_2(\text{CNCH}_3)_6](\text{BF}_4)_2$, powder sample, 77 K. Emission maximum at 542 nm.





- Figure 8.3**
- (a) Electronic absorption spectrum of $[\text{Pt}_2(\mu\text{-dppm})_2(\mu\text{-CNCH}_3)(\text{CNCH}_3)_2](\text{BF}_4)_2$ in acetonitrile (300 K). Absorption maximum at 452 nm.
- (b) Uncorrected emission spectrum of $[\text{Pt}_2(\mu\text{-dppm})_2(\mu\text{-CNCH}_3)(\text{CNCH}_3)_2](\text{BF}_4)_2$, powder sample, 77 K. Emission maximum at 604 nm. (A grating anomaly occurs at ~667 nm.)



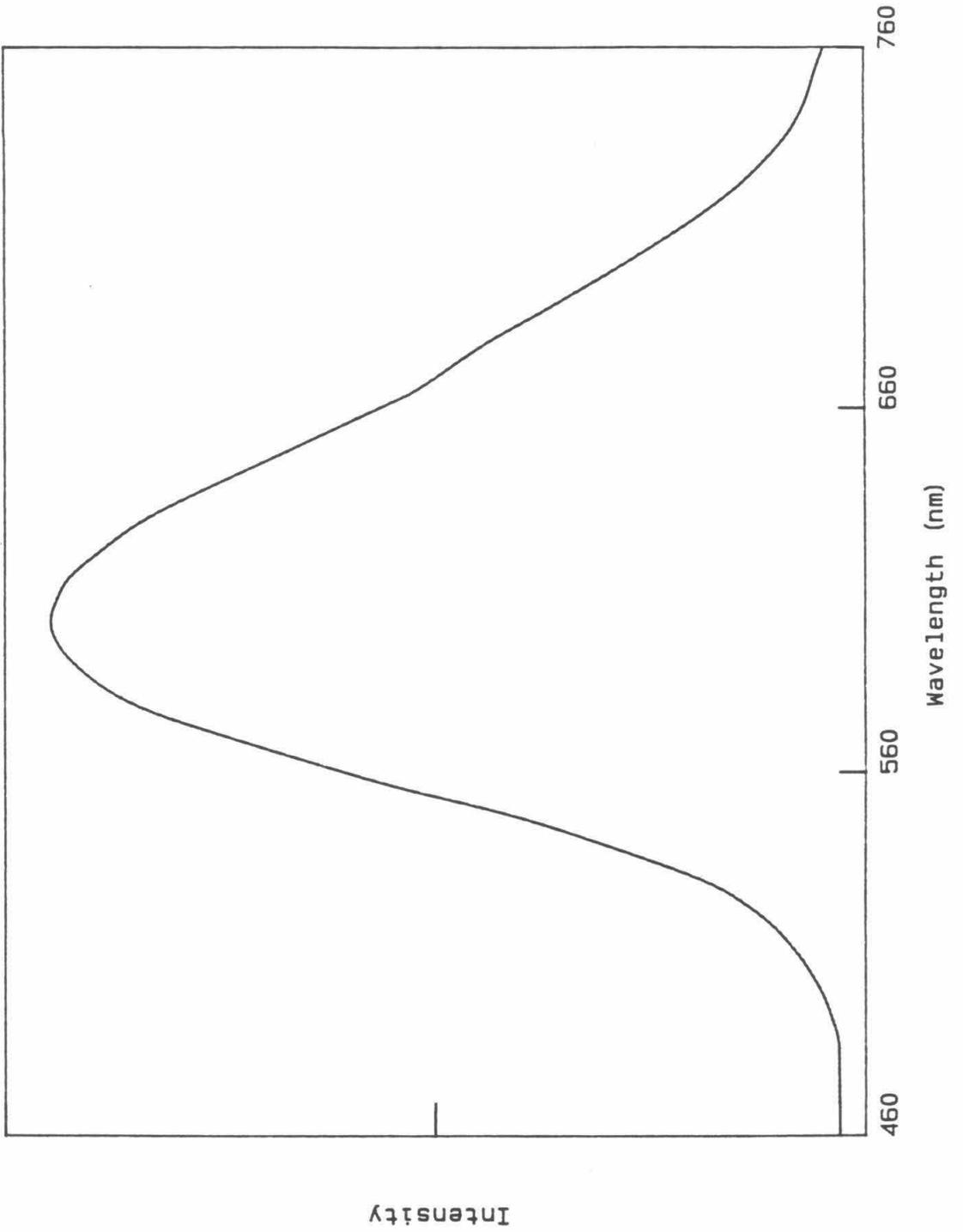
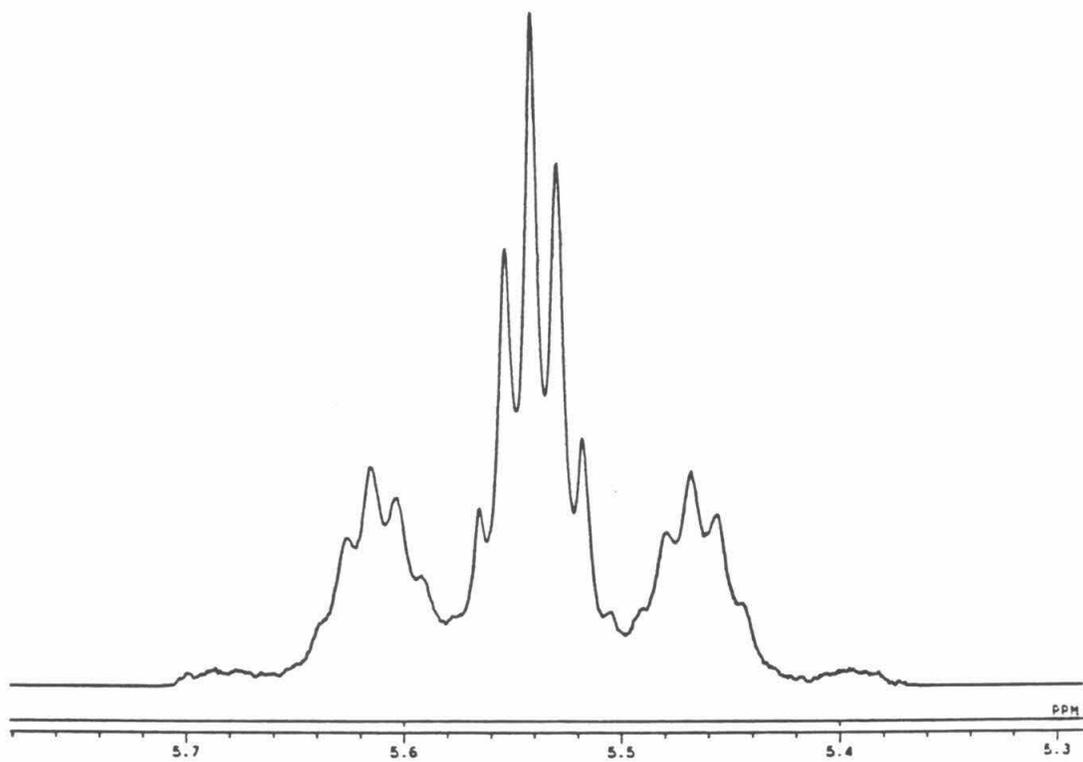


Figure 8.4 ^1H NMR spectrum of the methylene protons in the bridging dppm ligand of $[\text{Pt}_2(\mu\text{-dppm})_2(\mu\text{-CNCH}_3)(\text{CNCH}_3)_2](\text{BF}_4)_2$ in d_6 -acetone. $^3\text{J}(\text{PtH})$ coupling 58.4 Hz; $^2\text{J}(\text{PH})$ coupling 9.3 Hz.



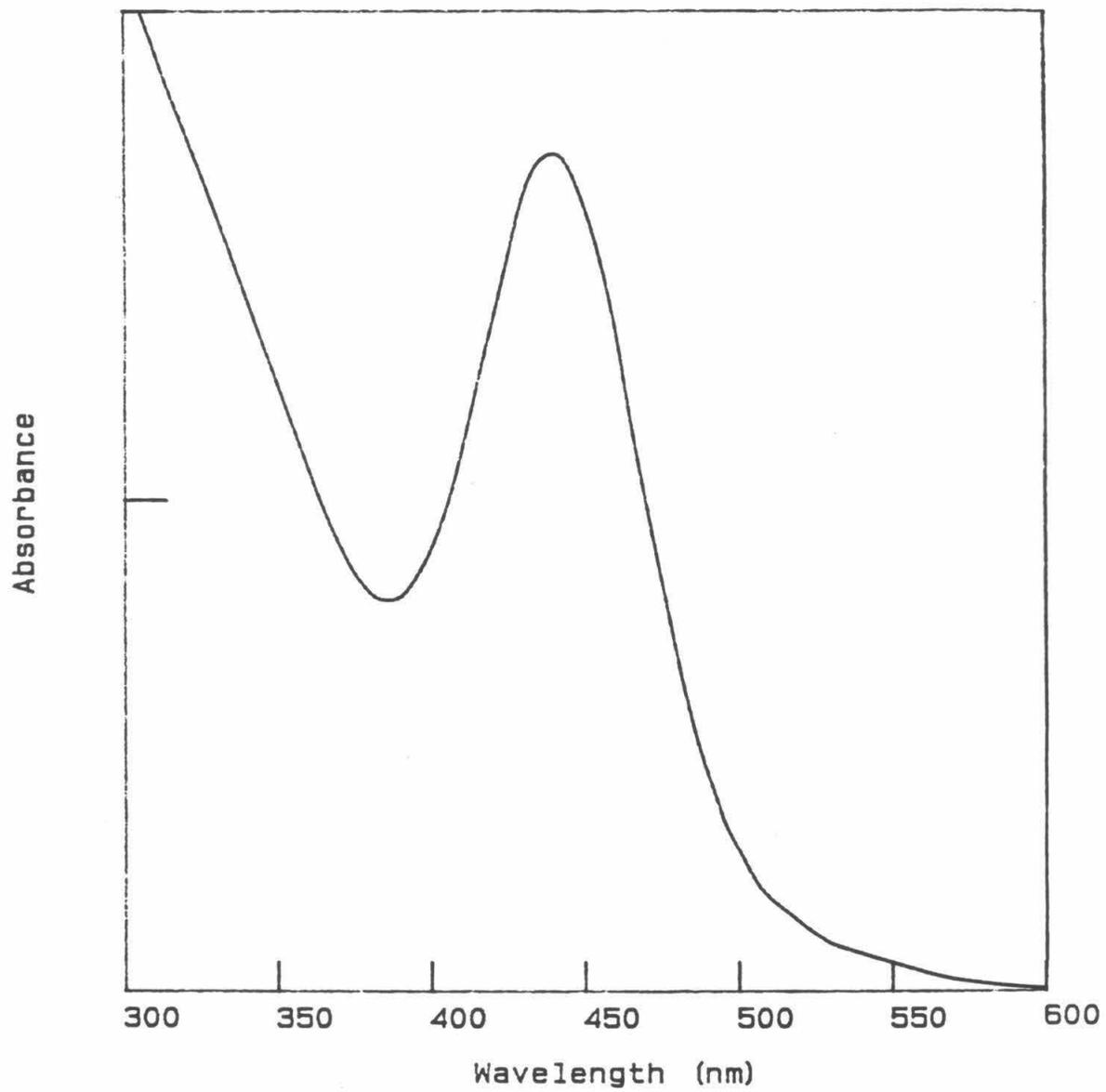
ilar platinum(II) monomers. The ^{31}P NMR revealed only one phosphorus signal at $\delta = 4.41$ ppm.

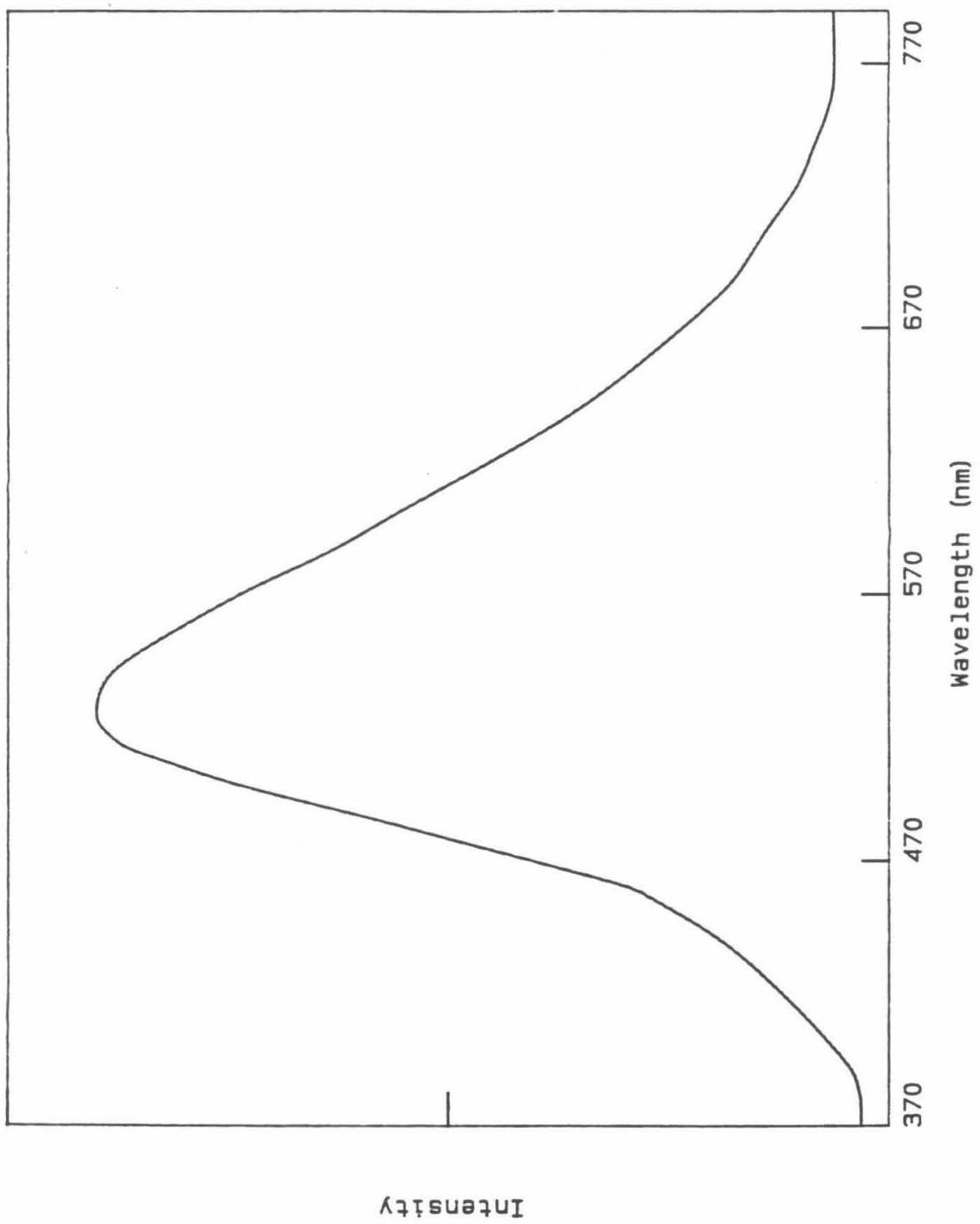
Similar reaction between $\text{Pd}_2(\text{CNCH}_3)_6^{2+}$ and dppm afforded $\text{Pd}_2(\mu\text{-dppm})_2(\mu\text{-CNCH}_3)(\text{CNCH}_3)_2^{2+}$ which has been crystallographically characterized.⁸ Based on the palladium reaction, the platinum complex obtained is presumed to be $\text{Pt}_2(\mu\text{-dppm})_2(\mu\text{-CNCH}_3)(\text{CNCH}_3)_2^{2+}$. Of course, additional characterization and possible structural evaluation is necessary to confirm this structural assignment. Attempts to grow crystals have been hampered by the slow decomposition of the complex when in solution. (Further loss of CNCH_3 is detected when the sample remains in solution for prolonged periods of time.)

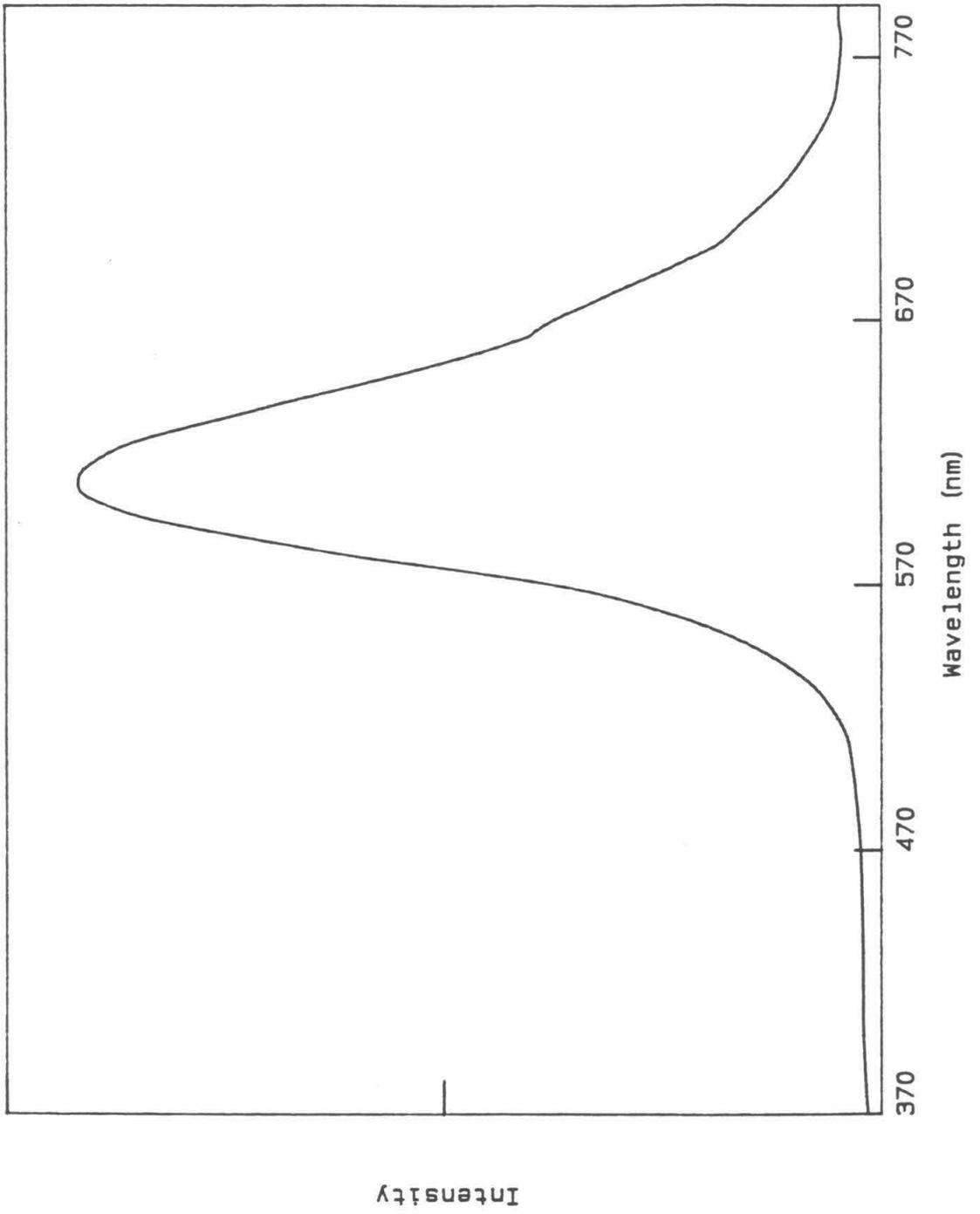
For this new dppm A-frame $d^8\text{-d}^8$ complex, as for the studied $d^9\text{-d}^9$ starting materials, no fluid solution or solid state luminescence at room temperature was observed. However at low temperature in the solid state, orange emission was detected ($\lambda_{\text{max}} = 604$ nm); see Figure 8.3b. Reaction of $\text{Pt}_2(\text{CNCH}_3)_6^{2+}$ with bis(dimethylphosphino)methane, dmpm, afforded a complex similar in absorption and emission properties as the complex formed from the $\text{Pt}_2(\text{CNCH}_3)_6^{2+} + \text{dppm}$ reaction.

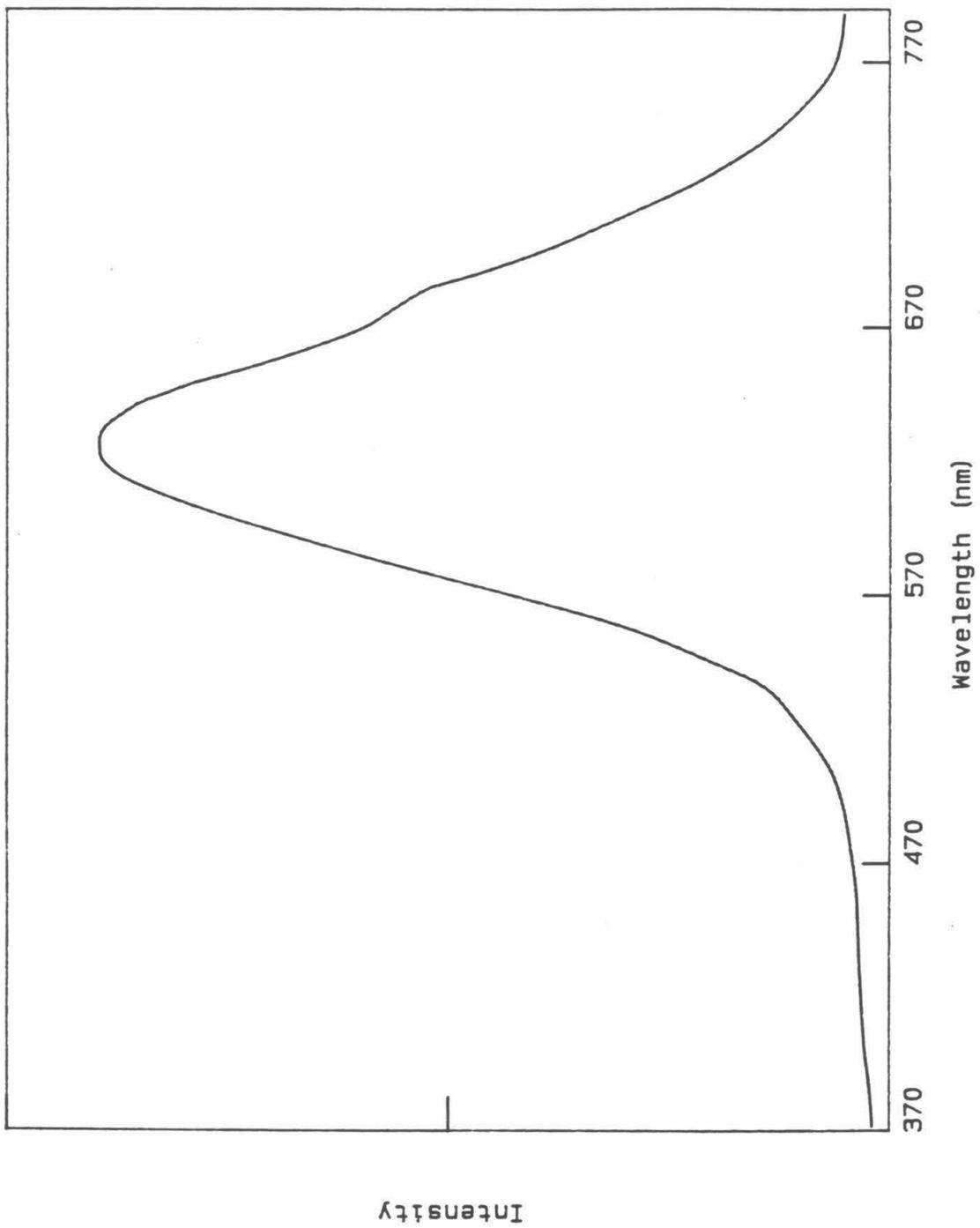
Additional reactions between $\text{Pt}_2(\text{CNCH}_3)_6^{2+}$ and other bridging ligands were investigated in hopes of unearthing a dimeric complex which did show fluid solution emission at ambient temperature. Reaction of $\text{Pt}_2(\text{CNCH}_3)_6^{2+}$ with 1,3-diisocyanopropane, $\text{CN}(\text{CH}_2)_3\text{NC}$, resulted in the expected orange solution ($\lambda_{\text{max}} = 443$ nm, see Figure 8.5a) and also produced a species which luminesced in fluid solution at room temperature (emission spectra are shown in Figure 8.5b-d). These complexes appear to be air and possibly even moisture sensitive. Eventual decomposition of the product is observed. The emission lifetime of

- Figure 8.5**
- (a) Electronic absorption spectrum of $[\text{Pt}_2(\mu\text{-CN}(\text{CH}_2)_3\text{NC})_2(\mu\text{-CNCH}_3)(\text{CNCH}_3)_2](\text{BF}_4)_2$ in acetonitrile (300 K). Absorption maximum at 443 nm.
- (b) Uncorrected emission spectrum of $[\text{Pt}_2(\mu\text{-CN}(\text{CH}_2)_3\text{NC})_2(\mu\text{-CNCH}_3)(\text{CNCH}_3)_2](\text{BF}_4)_2$ in acetonitrile, 300 K. Emission maximum at 529 nm. (A grating anomaly occurs at ~667 nm.)
- (c) Uncorrected emission spectrum of $[\text{Pt}_2(\mu\text{-CN}(\text{CH}_2)_3\text{NC})_2(\mu\text{-CNCH}_3)(\text{CNCH}_3)_2](\text{BF}_4)_2$ in dichloromethane, 77 K. Emission maximum at 613 nm. (A grating anomaly occurs at ~667 nm.)
- (d) Uncorrected emission spectrum of $[\text{Pt}_2(\mu\text{-CN}(\text{CH}_2)_3\text{NC})(\mu\text{-CNCH}_3)(\text{CNCH}_3)_2](\text{BF}_4)_2$, powder sample, 77 K. Emission maximum at 626 nm. (A grating anomaly occurs at ~667 nm.)









this complex in room temperature fluid solution is 220 ns, and increases to 7.4 μ s at 77 K in frozen solution. As in the dppm case, this complex is presumed to be $[\text{Pt}_2(\mu\text{-CN}(\text{CH}_2)_3\text{NC})_2(\mu\text{-CNCH}_3)(\text{CNCH}_3)](\text{BF}_4)_2$.

Why is fluid solution emission observed with the bridging isocyanide ligand and not the phosphine bridging ligand? Apparently, the bridging isocyanide ligand is decreasing the platinum dimer's non-radiative rate constant. One possible non-radiative deactivation pathway in the dppm complex is rapid cleavage and reassembly of the Pt-P(dppm) bond.¹⁸ Replacement of the neutral dppm bridging ligand with the bridging 1,3-diisocyanopropane ligand may hinder this deactivation pathway by inducing an electrostatic interaction between the Pt(+1) center and the formal negative charge on the carbon atom in the isocyanide ligand.

Obviously, further work in the areas of product characterization and photophysical analysis is required, and merited, on these reactions.

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APPENDIX

**Anisotropic Thermal Parameters and
Listings of Observed and Calculated Structure Factors
for $\text{K}_4[\text{Pt}_2(\text{P}_2\text{O}_5\text{H}_2)_4\text{Cl}]$ at 22 and 300 K**

Table A.1

Anisotropic Parameters Pt₂Cl at 22 K.

$$U \times 10^4$$

Atom	U_{11}	U_{22}	U_{33}	U_{12}	U_{13}	U_{23}
Cl	241(11)	241(11)	378(64)	0	0	0
K1	172(5)	172(5)	107(9)	61(7)	0	0
K2	103(5)	103(5)	197(10)	15(6)	0	0
P	86(4)	82(4)	103(4)	-18(4)	13(4)	0(4)
O1	82(13)	137(14)	148(16)	-19(11)	26(12)	7(13)
O2	152(14)	68(13)	166(17)	-25(11)	27(14)	10(13)
O3	204(24)	174(24)	90(19)	-90(18)	0	0

Table A.2

Anisotropic Parameters Pt₂Cl at 300 K. $U \times 10^4$

Atom	U_{11}	U_{22}	U_{33}	U_{12}	U_{13}	U_{23}
Cl	835(31)	835(31)	334(42)	0	0	0
K1	478(10)	478(10)	250(11)	202(12)	0	0
K2	338(8)	338(8)	490(18)	36(10)	0	0
P	202(5)	194(5)	241(5)	-29(4)	32(5)	8(5)
O1	284(18)	317(18)	418(23)	-48(15)	117(18)	2(19)
O2	352(20)	192(16)	544(28)	-52(14)	52(21)	82(19)
O3	491(37)	584(42)	235(22)	-335(30)	0	0

Table A3 (for deposit).

Structure Factors for Pt₂Cl at 22 K.

The columns contain, in order,
1, 10F_o, 10F_c and $10[(F_o^{**2} - F_c^{**2})/\sigma(F_o^{**2})]$;
A minus sign preceding 10F_o indicates that F_o^{**2} is
negative.

Structure Factors for Pt ₂ Cl at 22 K.										Page	2					
2	1567	1490	49	5	1753	1766	-7	8	289	277	6	8	4	1		
3	2497	2506	-4	6	2238	2216	10	9	373	355	10					
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6	1598	1591	3	9	1241	1245	-2		7	7	1	2	513	444	66	
7	124	88	13	10	226	172	18					3	1220	1200	14	
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9	1307	1325	-11						1	1346	1328	9	5	211	43	67
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3	32	37	-1	5	-63	11	-8	9	952	947	2					
4	611	613	-2	6	257	229	18	10	191	123	15					
5	77	65	4	7	451	424	19		8	0	1					
6	222	194	17	8	282	230	27									
7	67	102	-11	9	127	71	12		0	3897	3935	-11	4	143	135	4
8	480	447	24	10	44	32	0		1	1568	1529	25	5	326	313	10
9	64	36	4						2	1677	1634	25	6	95	90	1
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11	113	126	-3						4	283	245	31	8	331	302	18
	6	4	1		0	1606	1619	-8	5	2088	2072	7	9	134	18	21
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2	1050	1035	13	4	638	608	28	9	1416	1413	1					
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4	729	717	11	6	1034	990	33	11	1450	1419	15		1	642	629	13
5	1428	1417	7	7	327	328	0						2	1172	1147	19
6	1608	1608	0	8	943	942	0		8	1	1		3	2016	1986	15
7	379	378	0	9	889	883	4						4	687	670	14
8	1210	1214	-2	10	258	253	2		0	912	935	-21	5	1043	1020	17
9	981	987	-4	11	720	696	14		1	57	85	-12	6	1370	1342	17
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7	122	76	15	9	163	124	14		8	2	1		4	308	293	11
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Structure Factors for Pt2Cl at 22 K.										Page	3				
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0	95	95	0	4	367	353	8	4	-80	129	-28				
1	141	125	5	5	518	536	-11	5	186	184	0	18	1	1	
2	107	29	13	6	675	698	-14	6	136	98	8				
3	503	486	10									0	125	11	17
4	132	122	2	16	5	1		17	3	1		1	-138	31	-26
5	70	23	4	0	564	561	2					2	76	79	0
	15	11	1	1	296	298	-1	0	1245	1242	1	3	-149	25	-28
				2	352	329	13	1	508	540	-21	4	-102	7	-12
0	1133	1159	-16	3	112	125	-3	2	441	466	-15	5	-41	18	-2
1	619	624	-3	4	133	164	-10	3	985	986	0				
2	573	601	-19	5	410	374	20	4	185	215	-12	18	2	1	
3	948	963	-9	6	159	125	9	5	693	694	0				
4	391	367	13					6	992	975	10	0	1199	1230	-18
	15	12	1	16	6	1						1	695	686	6
				0	1515	1502	7	17	4	1		2	770	800	-19
0	358	335	13	1	409	417	-5	0	181	220	-17	3	1083	1086	-1
1	335	72	15	2	751	781	-21	1	104	136	-9	4	358	381	-12
2	208	137	25	3	1398	1372	15	2	-120	92	-31	5	938	942	-2
3	135	57	16	4	277	304	-14	3	89	115	-6				
	15	13	1	5	614	651	-24	4	-30	113	-15	0	258	260	-1
				6	1072	1062	6	5	75	154	-19	1	-137	35	-26
0	745	765	-13									2	200	226	-11
1	199	195	1	16	7	1		17	5	1		3	49	28	1
	16	0	1	0	322	331	-5	0	1281	1260	12	4	159	162	0
				1	208	218	-4	1	560	581	-13				
0	1679	1696	-9	2	94	122	-7	2	458	475	-10	18	4	1	
1	897	926	-20	3	218	179	16	3	956	960	-1				
2	1195	1188	4	4	-107	19	-15	4	257	303	-23	0	1130	1123	4
3	1855	1845	4	5	307	291	7	5	691	693	-1	1	150	208	-22
4	670	695	-17									2	669	691	-14
5	1460	1421	20	16	8	1		17	6	1		3	1534	1514	10
6	1354	1320	18	0	1664	1692	-14	0	529	515	8	4	479	504	-14
7	-113	96	-24	1	751	770	-12	1	-21	112	-16				
	16	1	1	2	870	880	-6	2	264	299	-18	18	5	1	
				3	1603	1605	0	3	404	380	13	0	126	72	11
0	309	320	-6	4	462	450	7	4	91	120	-6	1	-74	78	-14
1	248	262	-7	5	1161	1159	0	5	38	57	-1	2	-83	20	-8
2	254	253	0									3	-36	30	-2
3	302	279	12									4	-91	23	-9
4	-102	16	-14	0	393	375	10	0	1495	1490	2				
5	398	394	2	1	62	93	-6	1	579	601	-14	0	1244	1245	0
6	556	532	14	2	167	182	-5	2	747	776	-19	1	520	566	-30
7	-37	37	-2	3	186	220	-14	3	1572	1548	12	2	889	890	0
	16	2	1	4	-56	67	-9	4	386	430	-26	3	1374	1357	9
0	1700	1717	-8	16	10	1		17	8	1		18	7	1	
				0	1225	1209	9	0	88	157	-20	0	118	119	0

Structure Factors for Pt ₂ Cl at 22 K.										Page	7				
1	-103	49	-16	1	289	322	-16				19	5	1		
2	159	123	10	2	554	620	-42	0	1380	1357	12				
3	115	4	13	3	1056	1049	3	1	578	605	-16	0	1655	1605	24
								2	602	637	-21	1	795	769	15
	18	8	1		19	2	1	3	1329	1325	2	2	671	694	-14
0	1357	1372	-8	0	188	230	-17		19	4	1		19	6	1
1	585	607	-13	1	-88	85	-17								
				2	346	324	10	0	115	24	13	0	351	341	5
	19	1	1	3	296	271	10	1	132	171	-12				
0	926	940	-8		19	3	1	2	-109	4	-13				

Table A4 (for deposit).

Structure Factors for Pt₂Cl at 300 K.

The columns contain, in order, k , $10F_o$, $10F_c$ and $10|(F_o^{**2} - F_c^{**2})/\sigma(F_c^{**2})|$.
A minus sign preceding $10F_o$ indicates that F_o^{**2} is negative.

Structure Factors for Pt2Cl at 300 K.											Page	2			
5	924	925	0	2	572	573	-1	12	167	166	0	7	215	228	-7
				3	159	65	50	13	-58	45	-9	8	-47	21	-4
	6	k	1	4	900	922	-21	14	304	294	7				
				5	107	52	21	15	156	57	\$1		22	k	1
0	1088	1147	-59	6	609	635	-29	16	315	310	2				
1	550	564	-21	7	243	243	0					0	166	207	-20
2	529	467	87	8	268	240	25		17	k	1	1	75	9	8
3	-121	56	-67	9	163	108	30					2	270	283	-8
4	1121	1110	10	10	432	441	-9	1	455	462	-7	3	115	4	18
5	132	25	57	11	141	135	3	2	123	90	13	4	266	282	-10
6	948	887	48	12	585	593	-6	3	360	380	-18	5	-55	1	-4
								4	119	50	21				
	7	k	1					5	367	384	-14		0	k	2
					13	k	1	6	-84	4	-14				
1	1164	1140	22	1	637	638	-1	7	417	427	-9	0	2673	2556	42
2	197	174	28	2	173	137	26	8	-35	44	-6				
3	598	586	15	3	516	532	-19	9	397	383	11		1	k	2
4	241	180	65	4	110	42	24	10	112	72	12				
5	445	451	-8	5	595	608	-14	11	222	220	1	1	2570	2404	62
6	296	266	33	6	147	148	0	12	87	27	11				
7	1117	1122	-3	7	342	365	-24	13	295	310	-11		2	k	2
				8	41	5	3	14	68	10	7				
	8	k	1	9	557	553	3	15	276	297	-15	0	1451	1457	-4
				10	83	37	12					1	270	340	-114
0	1484	1483	1	11	444	476	-33		18	k	1	2	104	76	11
1	196	80	102	12	77	51	6								
2	960	971	-11	13	437	439	-1	0	418	436	-16		3	k	2
3	360	322	47					1	54	10	5				
4	154	3	66					2	413	435	-21	1	351	318	48
5	289	273	15		14	k	1	3	94	39	13	2	205	296	-118
6	699	679	22	0	325	319	6	4	210	251	-28	3	1807	1746	31
7	103	46	24	1	150	118	20	5	61	48	2				
8	999	1011	-9	2	526	540	-15	6	379	385	-5		4	k	2
				3	58	33	5	7	101	21	16				
	9	k	1	4	669	671	-2	8	388	406	-15	0	964	825	140
				5	116	29	27	9	67	1	7	1	385	449	-103
1	637	625	14	6	558	568	-10	10	259	282	-16	2	2310	2253	28
2	248	222	29	7	135	78	25	11	120	10	23	3	334	373	-58
3	748	729	20	8	324	346	-21	12	235	255	-13	4	622	546	74
4	68	11	14	9	108	67	14	13	-40	29	-4				
5	970	971	-1	10	377	365	10						5	k	2
6	148	114	23	11	71	9	10		19	k	1				
7	595	601	-6	12	461	477	-15					1	654	599	72
8	121	63	27	13	68	31	7	1	260	262	-1	2	-68	18	-20
9	660	694	-28	14	284	294	-5	2	84	29	10	3	1945	1910	19
								3	411	405	5	4	219	254	-44
	10	k	1					4	113	49	16	5	1434	1437	-1
					15	k	1	5	461	452	7				
0	813	831	-20	1	620	638	-19	6	54	11	4		6	k	2
1	193	162	30	2	142	23	41	7	322	315	5				
2	793	794	0	3	422	426	-4	8	-51	16	-5	0	479	513	-55
3	278	249	30	4	43	5	4	9	202	177	13	1	-89	68	-52
4	548	567	-24	5	308	310	-1	10	98	37	13	2	1433	1396	28
5	234	201	29	6	27	54	-4	11	283	265	12	3	414	429	-22
6	479	470	10	7	504	494	10	12	144	62	25	4	1113	1072	36
7	126	25	37	8	135	122	6					5	267	163	99
8	766	791	-26	9	493	512	-19		20	k	1	6	1017	999	14
9	-35	49	-9	10	83	33	11								
10	524	552	-23	11	437	429	6	0	260	275	-10		7	k	2
				12	111	46	18	1	104	107	0				
				13	194	179	9	2	261	299	-27	1	964	974	-10
				14	141	20	31	3	139	96	15	2	277	263	20
1	450	444	7	15	303	307	-1	4	426	412	11	3	1221	1199	18
2	158	77	51					5	115	70	13	4	270	155	103
3	920	912	7		16	k	1	6	267	286	-12	5	629	606	27
4	180	127	39					7	71	95	-6	6	395	403	-10
5	674	680	-6	0	700	709	-9	8	163	207	-23	7	1515	1536	-12
6	162	132	20	1	200	131	42	9	106	59	12				
7	649	656	-8	2	424	423	0	10	252	254	-1		8	k	2
8	104	85	8	3	225	210	11								
9	284	292	-7	4	301	307	-6		21	k	1	0	1497	1474	16
10	127	18	35	5	189	177	7					1	249	124	114
11	566	577	-8	6	361	370	-8	1	271	278	-5	2	1197	1210	-10
				7	113	98	5	2	-61	26	-7	3	307	309	-2
				8	543	551	-8	3	313	325	-8	4	458	378	91
				9	77	41	8	4	-44	2	-3	5	253	273	-22
0	365	353	14	10	403	428	-24	5	266	275	-5	6	1034	998	32
1	396	391	5	11	157	128	13	6	-75	39	-11	7	187	120	49

Structure Factors for Pt2Cl at 300 K.										Page	4				
5	1778	1753	14	0	1348	1348	0	21	k	3	3	205	175	26	
6	282	279	3	1	232	196	26				4	476	469	8	
7	974	966	7	2	1039	1022	13	1	537	534	2	5	161	97	37
8	107	66	16	3	103	66	12	2	146	172	-11	6	608	604	4
9	956	940	14	4	567	549	17	3	702	689	10	7	216	186	23
10	228	212	12	5	84	86	0	4	124	0	21	8	459	459	0
11	1110	1112	-1	6	884	882	1	5	651	630	17				
				7	191	134	30	6	88	114	-7	9	k	4	
	12	k	s	8	1091	1084	5	7	422	435	-9				
				9	152	131	10					1	556	564	-10
0	857	846	10	10	765	766	-1	22	k	3	2	180	144	28	
1	138	91	27	11	91	76	4				3	233	185	41	
2	1427	1414	8	12	505	529	-22	0	616	594	17	4	86	20	19
3	270	242	26	13	145	96	17	1	74	27	6	5	404	414	-11
4	1821	1796	14	14	555	556	-1	2	533	542	-7	6	75	17	13
5	209	207	1	15	82	30	9	3	32	19	0	7	563	571	-8
6	1136	1135	1									8	74	84	-4
7	121	67	22					0	k	4	9	493	510	-13	
8	729	742	-12												
9	95	73	8	1	1149	1136	9	0	1661	1568	76	10	k	4	
10	1053	1044	7	2	226	199	17								
11	144	96	21	3	660	647	11	1	k	4	0	230	235	-5	
12	1253	1243	6	4	-16	53	-5				1	126	97	17	
				5	624	625	-1	1	773	760	12	2	485	489	-4
	13	k	s	6	228	196	20				3	140	11	47	
				7	999	983	12	2	k	4	4	431	453	-26	
1	846	827	18	8	137	40	28				5	80	23	14	
2	302	305	-3	9	853	846	5	0	1254	1257	-2	6	521	532	-11
3	1687	1676	6	10	174	114	26	1	320	346	-39	7	169	140	18
4	119	11	32	11	562	565	-2	2	514	483	32	8	236	234	1
5	1497	1481	10	12	-64	72	-16				9	106	24	23	
6	181	150	19	13	465	474	-7	3	k	4	10	169	143	11	
7	816	805	10	14	179	94	33								
8	248	214	26					1	644	638	7	11	k	4	
9	713	713	0					2	89	111	-17				
10	245	222	17					3	510	496	14	1	455	472	-20
11	1164	1165	0	0	777	769	7				2	188	161	21	
12	166	106	28	1	146	24	34	4	k	4	3	507	516	-10	
13	873	880	-4	2	744	743	0				4	145	105	21	
				3	48	51	0	0	561	554	10	5	187	141	31
	14	k	s	4	886	861	21	1	134	128	6	6	114	22	27
				5	114	8	21	2	416	405	15	7	377	389	-11
0	1306	1283	17	6	827	839	-10	3	172	24	85	8	147	87	28
1	198	147	36	7	-70	32	-10	4	540	425	97	9	161	138	13
2	1126	1129	-2	8	668	671	-2					10	183	163	12
3	131	6	36	9	74	44	6	5	k	4	11	594	598	-2	
4	1128	1129	0	10	514	524	-8								
5	118	98	8	11	123	71	16	1	214	41	125	12	k	4	
6	909	914	-3	12	565	556	7	2	192	17	98				
7	208	196	8	13	57	4	5	3	555	597	-57	0	306	305	1
8	1019	1000	15					4	75	27	13	1	128	106	12
9	106	17	20					5	1382	1379	2	2	433	449	-18
10	1061	1066	-3									3	290	280	9
11	145	116	13	1	618	633	-13	6	k	4	4	326	326	0	
12	684	701	-16	2	148	150	0				5	147	102	23	
13	124	72	17	3	829	815	11	0	1153	1215	-56	6	325	344	-18
14	622	614	5	4	-44	28	-4	1	256	253	3	7	29	17	1
				5	807	788	15	2	294	250	50	8	243	246	-2
	15	k	s	6	202	169	17	3	455	469	-17	9	105	75	11
				7	643	632	9	4	726	722	4	10	443	462	-18
1	1236	1230	4	8	137	37	26	5	135	133	1	11	99	12	18
2	324	325	0	9	516	508	7	6	200	127	39	12	213	219	-3
3	882	866	14	10	93	70	5								
4	97	46	14	11	568	557	8	7	k	4		13	k	4	
5	895	894	1												
6	224	189	22					1	735	732	3	1	92	107	-7
7	1007	1006	0					2	395	415	-28	2	88	45	13
8	133	55	26	0	436	443	-6	3	498	512	-17	3	451	467	-17
9	1072	1056	13	1	102	90	3	4	113	29	31	4	148	94	25
10	240	207	21	2	738	703	29	5	122	44	35	5	603	599	4
11	627	632	-4	3	74	90	-4	6	150	63	43	6	122	72	19
12	95	23	14	4	793	795	-1	7	449	442	5	7	212	234	-17
13	608	607	1	5	90	103	-3					8	61	33	5
14	132	92	14	6	669	642	22	8	k	4	9	252	255	-2	
15	594	613	-12	7	115	50	16				10	-29	3	-1	
				8	427	430	-2	0	283	248	39	11	191	197	-3
	16	k	s	9	65	61	0	1	324	294	36	12	106	54	14
								2	444	459	-19	13	300	264	18

Structure Factors for Pt ₂ Cl at 300 K.										Page	5				
				5	51	9	4			3	126	127	0		
	14 k	4		6	293	288	3		6 k	5	4	1210	1192	13	
				7	103	17	16				5	114	94	8	
0	366	365	0	8	156	159	-1	0	1307	1353	-37	6	757	766	-8
1	115	63	20	9	98	59	9	1	476	478	-2	7	277	284	-5
2	389	387	2	10	118	109	3	2	908	877	30	8	322	325	-2
3	225	204	16	11	-45	2	-3	3	117	57	26	9	100	109	-3
4	240	225	11	12	194	213	-10	4	1288	1290	-1	10	581	594	-13
5	55	57	0					5	120	26	33	11	193	178	9
6	252	272	-16		19 k	4		6	1041	1036	3	12	785	804	-12
7	100	6	18												
8	332	347	-13	1	243	244	0		7 k	5		13 k	5		
9	104	90	4	2	69	68	0								
10	281	288	-5	3	237	241	-2	1	1522	1528	-4	1	669	668	1
11	66	37	5	4	-55	4	-5	2	90	5	23	2	252	245	6
12	195	177	10	5	149	177	-13	3	749	736	13	3	834	823	10
13	94	16	14	6	68	8	7	4	285	249	35	4	118	38	25
14	178	191	-4	7	157	145	5	5	542	534	8	5	877	889	-10
				8	120	56	17	6	189	182	5	6	175	187	-7
	15 k	4		9	157	180	-11	7	1306	1294	7	7	430	427	3
				10	78	52	5					8	108	72	11
1	387	391	-4						8 k	5		9	594	603	-8
2	156	131	13		20 k	4		0	1787	1812	-15	10	124	121	1
3	258	263	-4					1	174	77	59	11	647	660	-11
4	100	46	15	0	231	223	4	2	1188	1201	-10	12	151	80	26
5	212	230	-12	1	-86	15	-12	3	170	182	-10	13	594	595	-1
6	71	112	-14	2	170	198	-14	4	222	120	73				
7	291	276	11	3	119	91	8	5	238	226	10		14 k	5	
8	-33	51	-7	4	205	219	-8	6	901	888	12	0	508	515	-7
9	247	241	4	5	87	17	10	7	153	68	38	1	122	58	23
10	112	68	13	6	129	131	0	8	1357	1366	-5	2	679	680	0
11	243	227	11	7	-49	18	-4					3	61	36	5
12	-30	34	-3	8	177	167	4		9 k	5		4	864	856	7
13	214	224	-6					1	1070	1073	-1	5	113	11	23
14	84	13	11		21 k	4		2	258	238	19	6	675	678	-2
15	239	237	0	1	162	146	6	3	799	791	8	7	77	8	10
				2	87	22	-10	4	81	20	15	8	470	477	-6
	16 k	4		3	181	203	-11	5	909	899	9	9	121	60	19
0	470	486	-15	4	-104	9	-16	6	112	112	0	10	543	543	0
1	78	12	11	5	220	182	18	7	1008	999	7	11	83	51	7
2	215	209	4	6	-56	25	-4	8	104	114	-4	12	559	566	-6
3	36	2	2					9	910	925	-10	13	89	49	9
4	228	242	-10		0 k	5						14	399	374	13
5	99	53	12	0	2208	2196	5		10 k	5			15 k	5	
6	207	219	-8					0	897	909	-12	1	789	793	-3
7	86	17	12					1	197	161	29	2	85	83	0
8	307	292	11		1 k	5		2	964	973	-8	3	538	540	-1
9	120	50	19					3	143	80	31	4	84	15	12
10	241	245	-2	1	1668	1633	19	4	815	822	-7	5	456	463	-7
11	35	23	1					5	188	180	5	6	124	122	0
12	199	175	12		2 k	5		6	836	817	17	7	617	619	-2
13	112	24	19					7	128	139	-6	8	115	92	7
14	256	214	25	0	1705	1676	17	8	940	947	-6	9	673	675	-1
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	17 k	4		2	1088	1042	34		555	588	-25	11	483	507	-21
1	297	291	5						11 k	5		12	-26	33	-3
2	119	48	20		3 k	5						13	228	259	-20
3	179	177	1	1	879	905	-27		642	614	29	14	-58	51	-9
4	-53	37	-7	2	78	7	20		98	27	21	15	361	391	-16
5	216	252	-25	3	1771	1749	11								
6	140	47	28					1	1123	1126	-2		16 k	5	
7	284	285	0		4 k	5		2	203	157	31				
8	63	88	-6					3	916	914	0	0	995	967	23
9	306	296	7	0	321	272	61	4	156	81	35	1	197	193	2
10	71	76	-1	1	485	503	-23	5	770	774	-3	2	584	583	0
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12	106	35	15	3	112	38	27	7	396	405	-8	4	349	334	12
13	197	164	16	4	2165	2171	-2	8	47	50	0	5	205	202	1
								9	737	738	0	6	434	464	-27
	18 k	4			5 k	5		10				7	133	131	0
0	146	88	22	1	956	963	-6		12 k	5		8	744	725	16
1	-15	11	0	2	178	154	21					9	127	71	17
2	222	231	-5	3	1286	1300	-11	0	467	460	8	10	564	568	-3
3	84	79	1	4	113	64	21	1	428	450	-24	11	178	162	8
4	335	330	4	5	1571	1569	1	2	762	777	-15	12	214	229	-9

Structure Factors for Pt ₂ Cl at 300 K.										Page	6				
13	115	83	9	1	1287	1279	5	11	k	6	13	450	463	-11	
14	371	359	9	2	158	109	34	14	-60	10	-5				
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6	35	34	0	4	2656	2678	-8	9	758	746	10	5	87	61	6
7	590	590	0					10	102	68	10	6	629	640	-9
8	97	59	9		5	k	6	11	859	852	4	7	156	190	-17
9	556	538	16									8	862	861	0
10	156	116	16	1	1206	1210	-3		12	k	6	9	146	98	17
11	296	296	0	2	191	165	22					10	622	627	-4
12	95	22	13	3	1763	1731	19	0	521	518	2	11	130	58	20
	18	k	5	4	241	240	1	1	130	121	4	12	386	393	-6
				5	1507	1560	-30	2	1057	1059	-1				
0	506	517	-10		6	k	6	3	256	241	11		17	k	6
1	64	36	4					4	1320	1318	0				
2	532	536	-3	0	1554	1565	-7	5	234	243	-7	1	729	731	-1
3	-89	25	-14	1	158	58	50	6	1013	1009	3	2	118	92	9
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5	111	48	15	3	177	151	20	8	532	544	-11	4	43	16	2
6	522	535	-10	4	1321	1317	3	9	121	119	0	5	661	632	24
7	-56	39	-7	5	216	218	-1	10	711	710	0	6	102	94	2
8	508	503	4	6	1264	1261	1	11	143	146	-1	7	746	746	0
9	100	13	15					12	933	953	-12	8	-34	3	-1
10	312	324	-9		7	k	6					9	607	611	-3
11	140	11	28	1	1803	1840	-22	1	789	772	16	10	-88	43	-15
	19	k	5	2	135	34	41	2	196	189	4	11	363	392	-23
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3	523	526	-2	6	44	3	4	6	156	91	27	2	558	570	-11
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2	409	408	1	6	1100	1096	3	0	90	38	13	1	509	513	-3
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	1	k	6	0	1189	1186	2		15	k	6		0	k	7
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	3	k	7	9	35	36	0	11	187	144	19									
				10	202	241	-19	12	81	26	8		5	k	8					
1	193	193	0		11	k	7		16	k	7		1	590	577	13				
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2	144	138	4						17	k	7									
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	7	k	7		13	k	7		0	148	139	3	2	836	870	-32				
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