Interacting single atoms with nanophotonics for chip-integrated quantum networks

Thesis by

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In Partial Fulfillment of the Requirements for the Degree of

Doctor of Philosophy



California Institute of Technology

Pasadena, California

2013

(Defended May 22, 2013)

© 2013 Daniel James Alton All Rights Reserved To my parents, sister, and loved ones.

Acknowledgments

First and foremost, I would like to thank God, the Almighty, for all of His blessings.

I would like to thank my advisor, professor Jeff Kimble, a mentor who I greatly respect and admire. Not just as a great scientist with exceptionally high standards of rigor and integrity, but it never cease to amaze me, professor Kimble's depth and breadth of thinking. Looking into matters from the biggest horizon to the smallest details, connecting the many dots and people to the matter, and looking far into the future and down to the shortest time scales in the experiment. It has been a great privilege and I am deeply grateful.

I would like to thank our collaborators, professor Kerry Vahala and professor Oskar Painter, for sharing their pioneering expertise and leadership in the development of novel photonic devices, with very high, world-record quality, which have been critical in the work described in this thesis. My thesis committee members, professor Jeff Kimble, professor Kerry Vahala, professor Olexei Motrunich, and professor Michael Roukes, who have been extraordinary sources of inspirations, thank you very much.

My sincere gratitude to professor Peter Zoller and professor Jun Ye, visiting professors in our group who have shared their insights and advice for certain parts of the work described in this thesis. To professor Jun Ye, for inviting me to spend a week at his laboratory at JILA, Boulder, Colorado, to learn about the making of optical amplifiers, and to the members in the Ye group, especially Sebastian Blatt.

The work described in this thesis is a result of contributions and hard work of many individuals that I would like to acknowledge here. In the experiments with microtoroidal resonators: Takao Aoki, a very talented scientist and experimentalist that I admire and learn greatly from; Scott Parkins, who enlightened us with the fundamental theoretical formulations especially in our photon router work; Nate Stern, a multi-talented scientist whose work and drive I admire both in the experimental and theoretical parts of our work; Hansuek Lee and Eric Ostby from Vahala group, highly talented individuals who shared their expertise and fabricated the wonderful microtoroidal photonic devices critical in our experiments; last but not least, Cindy Regal and Barak Dayan, extraordinary scientists whose works are simply amazing. In the experiments with optical nanofibers: Kyung Soo Choi and Akihisa Goban, extraordinarily talented and accomplished young scientists who I especially commend and thankful to not just for our collaboration in the lab but also outside of the lab for our enduring friendship. Ding Ding, Clement Lacroute, Martin Pototschnig, and Tobias Thiele, who made important contributions in our nanofiber trap work. In the experiments with photonic crystal based structures: Firstly in Lab 11, Akhisa Goban and Chen-Lung Hung whose cold atom transport design formed the basis of the experimental setup in Lab 1 discussed in this thesis. Chen-Lung's extraordinary talent and exceptionally open-minded approach had been critical and greatly valuable in the projects; The photonic crystal based device design/characterization/fabrication team: Su-Peng Yu, Jonathan Hood (Kimble group), Richard Norte, Sean Meneehan, and Justin Cohen (Painter group), whose work have been critical in our experiments; In Lab 2: Jae Hoon Lee, Juan Muniz, and Ding Ding; Thanks especially to Juan Muniz for our close collaboration in investigating atom trapping schemes, where Juan's great talents and contributions have been critical; Last but not least, the team in Lab 1 who have directly worked in the photonic crystal based experiments discussed in this thesis: Andrew McClung, an exceptionally talented researcher and computer expert whose combination with out-of-the-box thinking have led to various major contributions in our lab; Pol Forn-Diaz, a very talented scientist with broad interests and at the same time extraordinary attention to details, who I greatly respect and am thankful for to have worked together inside and outside of the lab; Martin Pototschnig, an insightful and highly talented experimentalist, who is capable of achieving amazing results especially when the insights are unleashed; and Clement Lacroute, for his important and valuable contributions. I also want to acknowledge Taofiq Paraiso and Alex Krause from Painter group.

I would like to thank the members of the Kimble group (Caltech Quantum Optics group), Vahala and Painter groups, and the Institute for Quantum Information and Matter (IQIM) at Caltech, who have been great sources of inspirations, friendships and supports at various locations, times, and situations, I am greatly thankful for these years where our paths have crossed, and hopeful that they will again in the future. In addition to the above-mentioned, I would like to specifically thank: Dalziel Wilson, Yi Zhao, Kang-Kuen Ni, Scott Papp, Russ Miller, Hui Deng, Tracy Northup, Julien Laurat, Elizabeth Wilcut Connolly, Scott Kelber, Daniel Chao, Sarah Kaiser, Michael Martin, Jiang Li, Kiyoul Yang, Darrick Chang, Alexey Gorshkov, Liang Jiang, K C Fong, and Matt Eichenfield. I would like to thank Scott Curtis our group's administrator who always got the job done exceptionally well regardless of the number of lines, people, organizations, or challenges involved. I also want to thank several particular administrators in the Physics department, Donna Driscoll, Marcia Brown, Alan Rice, Louisa Fung, and Loly Ekmekjian.

I would like to take this opportunity to thank my undergraduate research advisors, Prof. Ping Koy Lam, Prof. Tim Ralph, and Dr. Thomas Symul, at Australian National University and University of Queensland, who have introduced me to the world of quantum optics, to whom I am greatly and always be thankful.

I would like to thank my parents, sister, and grandparents, who have given their endless support, inspirations, and love in my life. I am forever grateful for their sacrifices, guidance and faith in me. Despite the extend of the oceans, masses of land, and the four different time zones that I have lived in for extended periods of time, they have always been there. I am blessed to have met a very special person whom I hold close and dear, who has enriched my life in every aspect since the last two years. I thank her for all the love and support. It is to them, I wish to dedicate this thesis.

> Daniel Alton May 2013 Pasadena, CA

Abstract

Underlying matter and light are their building blocks of tiny atoms and photons. The ability to control and utilize matter-light interactions down to the elementary single atom and photon level at the nano-scale opens up exciting studies at the frontiers of science with applications in medicine, energy, and information technology. Of these, an intriguing front is the development of quantum networks where $N \gg 1$ single-atom nodes are coherently linked by single photons, forming a collective quantum entity potentially capable of performing quantum computations and simulations. Here, a promising approach is to use optical cavities within the setting of cavity quantum electrodynamics (QED). However, since its first realization in 1992 by Kimble et al., current proof-of-principle experiments have involved just one or two conventional cavities. To move beyond to $N \gg 1$ nodes, in this thesis we investigate a platform born from the marriage of cavity QED and nanophotonics, where single atoms at ~ 100 nm near the surfaces of lithographically fabricated dielectric photonic devices can strongly interact with single photons, on a chip. Particularly, we experimentally investigate three main types of devices: microtoroidal optical cavities, optical nanofibers, and nanophotonic crystal based structures. With a microtoroidal cavity, we realized a robust and efficient photon router where single photons are extracted from an incident coherent state of light and redirected to a separate output with high efficiency. We achieved strong single atom-photon coupling with atoms located ~ 100 nm near the surface of a microtoroid, which revealed important aspects in the atom dynamics and QED of these systems including atom-surface interaction effects. We present a method to achieve state-insensitive atom trapping near optical nanofibers, critical in nanophotonic systems where electromagnetic fields are tightly confined. We developed a system that fabricates high quality nanofibers with high controllability, with which we experimentally demonstrate a stateinsensitive atom trap. We present initial investigations on nanophotonic crystal based structures as a platform for strong atom-photon interactions. The experimental advances and theoretical investigations carried out in this thesis provide a framework for and open the door to strong single atom-photon interactions using nanophotonics for chip-integrated quantum networks.

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2.1Overview of atom-photon interaction. a) Two level atom interacting with a single photonic mode a_p at rate g. b) Dressed atom energy levels E_{jn} where j = g, efor ground, excited states (dashed lines: absent atom-photon coupling). c) Excited atom decay rate into the photonic mode $a_{\rm p}$ (e.g., waveguide mode, intracavity mode), $\Gamma_{\rm p}$, and decay (loss) rate to the environment, Γ_0 . The coupling rate between the mode $a_{\rm p}$ and detector is κ , which is equal to $\Gamma_{\rm p}$ in direct detection, but may be different than $\Gamma_{\rm p}$ for a cavity system. d) Probability $P_{\rm e}$ of an initially excited atom to be in the excited state after a time $g_{iv}t$ where $g_{iv} = 105MHz$. (i) Atom free-space decay rate $\Gamma_0/2\pi = 5.2$ MHz. (ii) Enhanced decay rate $\Gamma_p = 2\Gamma_0$. (iii) With $g/2\pi = 105$ MHz, $\kappa/2\pi = 20$ MHz, $\Gamma_0/2\pi = 5.2$ MHz (Cesium D2 line) [5]. e) Atom-photon interaction strengths parametrized by $\chi = \Gamma_{\rm p}/\Gamma_0$ for waveguides 1. to 3. and cavities 4. to 8. Limits are discussed in main text. Inset: Some data points showing χ realized in various experiments, 1a-1b: Nanofiber trap in [248] and [91], also with the corresponding cooperativity parameter C for cavity QED systems with Fabry-Perot 2.2Atom-photon interaction without a cavity. a,c,e,f) Atom interacting with photonic mode $a_{\rm p}$, $E(\omega_{\rm p})$ is the oscillating electric field at optical frequency $\omega_{\rm p}$; σ^- and $d_{\rm ge}$ are atomic lowering operator and electric dipole moment; $H_{\rm int}$: atom-photon interaction Hamiltonian; $P_{\rm in}, P_R, P_T$: input, reflected, transmitted optical power; $\Gamma_{\rm p}$ and Γ_0 are decay rate into photonic mode a_p and decay (loss) rate into the environment respectively; $A_{\rm eff}$ and σ are photonic effective area and atomic scattering cross-section respectively. c) f: focal length of the pair of lenses; $w_{\rm in}$: input Gaussian beam waist (radius) size. b) Cesium D2 line energy levels/manifolds. d,g,h) T, R: transmittance and reflectance; $R_{\rm Sc} = \sigma_0 / A_{\rm eff}$, atom scattering rate, where σ_0 is the atomic resonant scattering cross-section. d) Results for strongly focused light; χ : full model; χ' : paraxial approximation; $u = w_{in}/f$, focusing strength; (i): Experimental result for T of [233]. Top g) Comparison between our approximate model (solid curves) and full results of [126] (dashed curves). Bottom g) Results using our model for parameters in [248] (fiber radius 250 nm) and [91] (fiber radius 215 nm) with measurements of (1-T) shown by (i) and (ii) respectively. The variable d is the atom to fiber's surface distance. h) Contour plot of χ . Points (i) and (ii) correspond to parameters in [248] and [91] respectively. \mathbf{g},\mathbf{h}) As evident in \mathbf{g}), the prediction model agrees with [91], point (ii), but this is not the case for [248], point (i). This is discussed further in the

2.3Atom-photon interaction within a cavity. a,b,c,d) Q: cavity quality factor, $V_{\rm m}$: cavity mode volume; $P_{\rm in}, P_R, P_T$: input, reflected, transmitted optical powers; g: atom-photon coupling rate; κ_i : intrinsic cavity loss; κ_{ex} : extrinsic input/output coupling rate; $E_{\rm max}$ and $E(\vec{r}_{\rm a})$ are the maximum electric field and the electric field at atom's position \vec{r}_{a} ; Γ_{p} : atom's decay rate into cavity photonic mode; Γ_{0} : atom's decay rate into the environment. **e,f)** C: cooperativity parameter; $\chi = \Gamma_{\rm p}/\Gamma_0$; d =atom-to-surface distance. a) Fabry-Perot cavity; labels in e) and f): a1, experimental parameters of [33], a2, ultimate limit [33]. b) Microtoroidal cavity; labels in e) and f): b1 and b2, experimental parameters of [9] and [5], b3 and b4, projected limits [224, 131]. c) Atomic mirror cavity (formed by 2 $N_{\rm m}$ atoms): c1, prediction from [41] with $|E|/|E_{\text{max}}| = 0.33$, c2, with $|E|/|E_{\text{max}}| = 1$. d) Photonic crystal cavity: d1-d4 for currently realizable $Q/V_{\rm m}$ value to the projected limit [147], with $|E|/|E_{\rm max}| = 0.5$, d5-d8 for same range of $Q/V_{\rm m}$ but with $|E|/|E_{\rm max}| = 1$ and an enhancement factor of 10 in atom's decay rate into the photonic mode that may be gained by utilization of photonic crystal band structure effect. Note: a1, a2, d5-d8 indicate values of χ (with $\Gamma_0 = \gamma_0$, the free-space decay rate), they are not functions of d. The horizontal lines

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Microtoroid cavity QED schematic. (a) Optical input/output coupling enabled 3.1by a tapered fiber (diameter $D_{\rm f}$) positioned at a small fiber/toroid surface-to-surface gap $x_{\rm ft} < \lambda$. Spatial cylindrical coordinates $\{\rho, \phi, z\}$ with origin at toroid center. The toroid geometry can be described by its major $(D_{\rm M})$, minor $(D_{\rm m})$, and principal $(D_{\rm p})$ diameters. On the toroid's cross-sectional minor circle plane, ψ describes the latitudinal angle, and $d = d(\rho, z)$ is the atom-to-surface distance. (b) Tapered fiber optical input/output fields $\{a_{in}, a_{out}, b_{in}, b_{out}\}$ coupled at rate κ_{ex} to toroid counterpropagating intracavity fields $\{a, b\}$, coupled by internal scatterers at a rate h, suffering intrinsic loss at rate κ_i . A nearby atom located at \vec{r} is coupled to the cavity at rate $g(\vec{r})$, and has a free space spontaneous emission rate γ . (c) i-iv): normalized electric field |E|profiles and the components $\{E_z^{\theta}, E_{\rho}^{\theta}, E_{\phi}^{\theta}\}$ (where θ indicates the optical phase), with $E_z^{\max} = 1.00 |E|^{\max}, E_{\rho}^{\max} = 0.086 |E|^{\max}, E_{\phi}^{\max} = 0.118 |E|^{\max}; v-vii)$ shows the lowestorder mode function $f(\rho, z)$ for a toroid with $\{D_p, D_m\} = \{24, 3\} \mu m, m = 118$ and $\lambda = 852$ nm, and the cross-sections along d and z. (d) SEM images of two fabricated mictoroids with $D_{\rm p} \sim 18 \ \mu {\rm m}$ and $D_{\rm p} \sim 24 \ \mu {\rm m}$.

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- 3.10 Tapered nanofiber fabrication and experiment overview. a) Schematic showing a nanofiber mounted on an aluminium holder inside an UHV chamber, with three pairs of counter-propagating magneto-optical trapping and cooling beams forming cold atom cloud overlapped with the nanofiber. b) Photograph of the vacuum chamber, with arrow pointing towards the red-glowing nanofiber. c, e) Close-up and environment pictures of our old taper-pulling setup. d, f) Close-up and clean-hood environment pictures of our improved taper-pulling setup used to fabricate tapered nanofibers for our nanofiber atom trap experiment.

- 3.11 Nanophotonic beam and mirror. a) Schematic of a nanobeam device showing optical fiber to silicon nitride waveguide butt-coupling, adiabatic adapter to nanobeam mode (z1-z6), a nanobeam waveguide with width w and height h, followed by a photonic crystal mirror (z8-z9). The dimensions are discussed in the text. b) SEM images of a fabricated device (courtesy of Painter group), showing a sample structure with \sim mm size thru-hole in (i), fiber butt coupling (ii), nanobeam waveguide with electric field profile (iii), and photonic crystal mirror at the end (iv). c) Normalized electric field |E| profiles and the components $\{E_x^{\theta}, E_y^{\theta}, E_z^{\theta}\}$ (where θ indicates the optical phase and location along the z-axis). E.g., $\theta = 0 \leftrightarrow z = 0, \theta = \pi/2 \leftrightarrow z = \lambda/4$) for the nanobeam fundamental HE₁₁ mode polarized along x, with $E_x^{\max} = 0.840|E|^{\max}, E_y^{\max} = 0.340|E|^{\max}, E_z^{\max} = 0.560|E|^{\max}$.

- 3.13Double nanophotonic beam and mirror. a, c) Schematic of a double nanobeam device showing optical fiber to silicon nitride waveguide butt-coupling, adiabatic adapter to nanobeam mode, a Y-junction single-to-double beam mode converter, a double nanobeam waveguide with width w, height h, separated by a gap, followed by a photonic crystal mirror. The dimensions are discussed in the text. **b**) Dispersion curves showing effective refractive index $n_{\rm eff} = \beta/k$ where $\beta =$ propagation constant of the guided mode, $k = 2\pi/\lambda$ = free-space wave number, $\alpha = 200$ nm and n =2.0, of the first lowest order supermodes, for symmetric (even) modes: x-polarized (i) and y-polarized (ii), and anti-symmetric (odd) modes: x-polarized (iii) and ypolarized (iv). Higher-order modes start to appear beyond $V \simeq 3$ in the shaded region. c) Double nanobeam waveguide. d) Normalized electric field |E| profiles and the components $\{E_x^{\theta}, E_y^{\theta}, E_z^{\theta}\}$ (where θ indicates the optical phase and location along the z-axis. E.g., $\theta = 0 \leftrightarrow z = 0, \theta = \pi/2 \leftrightarrow z = \lambda/4$ for the double nanobeam (w = 300 nm, h = 200 nm, gap = 200 nm) four lowest order modes polarized along x. For (i), $E_x^{\text{max}} = 0.920 |E|^{\text{max}}, E_y^{\text{max}} = 0.440 |E|^{\text{max}}, E_z^{\text{max}} = 0.429 |E|^{\text{max}}.$ For (ii), $E_x^{\max} = 0.454 |E|^{\max}, E_y^{\max} = 0.895 |E|^{\max}, E_z^{\max} = 0.469 |E|^{\max}$. For (iii), $E_x^{\max} \ = \ 0.863 |E|^{\max}, E_y^{\max} \ = \ 0.445 |E|^{\max}, E_z^{\max} \ = \ 0.572 |E|^{\max}. \ \ \text{For \ (iv)}, \ E_x^{\max} \ = \ 0.572 |E|^{\max}.$ $0.498|E|^{\max}, E_y^{\max} = 0.910|E|^{\max}, E_z^{\max} = 0.446|E|^{\max}, \dots, \dots, \dots, \dots, \dots$
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- 3.15Nanophotonic beams and cavities experiment setup. a) Schematic of experimental setup showing two chambers separated by 70 cm connected by a differential pumping tube, where a magneto-optically trapped atom cloud is formed in the first chamber (i), pierced through by a near-resonant push beam (green arrow) that forms a jet of atoms, to be captured by a second magneto-optical trap in the science chamber (vi) formed by three pairs of counter-propagating beams shown by the red arrows, and in b). Following this stage, the cloud of atoms in the science chamber is transported and recaptured by a mini-magneto-optical trap inside the chip's thru-hole over the nanophotonic devices, formed by three pairs of small counter-propagating cooling and trapping beams shown in b). The setup is designed with multiple vacuum valves (ii), (iii), (iv) allowing frequent loading/unloading of nanophotonic device chip mounted on a multiplexer (vii) and translation stage (viii). c) Fluorescence image showing atom cloud transport from science chamber large MOT to mini-MOT inside the chip, taken with CCD camera with viewing direction shown by the cyan arrow in a) and b), also shown on the right panel of b). d) Setup built for our experiment.

- 4.3 (a) False detection ratio F, (b) transmitted signal $\mathcal{T}_0(t=0)$ at the center of an atomic transit, and (c,d) intensity correlation functions $g_{T,R}^{(2)}(\tau = 0)$ at zero time delay for the transmitted T and reflected R light as functions of the threshold C_{th} for the selection of atom transits. In all cases, $\Delta t_{atom} = 4\mu s$ and $\bar{n} = 0.093$ 100
- Radiative interactions and optical potentials for an atom near the surface 5.1of a toroidal resonator. (a) Simple overview of the experiment showing a cloud of cold cesium atoms released so that a few atoms fall within the evanescent field of a microtoroidal resonator. Light in a tapered optical fiber excites the resonator with input power $P_{\rm in}$ at frequency $\omega_{\rm p}$, leading to transmitted and reflected outputs $P_{\rm T}, P_{\rm R}$. (b) Cross section of the microtoroid at $\phi = 0$ showing the coherent coupling coefficient $|g(\vec{r}) = g(\rho, z, \phi)|$ for a TE polarized whispering-gallery mode. The microtoroid has principal and minor diameters $(D_{\rm p}, D_{\rm m}) = (24, 3) \ \mu {\rm m}$, respectively. (c) (i) Coherent coupling $|g(d, z, \phi)|$ for the external evanescent field as a function of distance d = $\rho - D_{\rm p}/2$ from the toroid's surface for $(z, \phi) = (0, 0)$. (ii) The effective dipole potentials $U_{\rm d}$ for resonant $\omega_{\rm p} = \omega_{\rm a}^{(0)}$, red $\omega_{\rm p} < \omega_{\rm a}^{(0)}$ and blue $\omega_{\rm p} > \omega_{\rm a}^{(0)}$ free-space detunings of the probe $P_{\rm in}$ (intracavity photon number ~ 0.1, circulating power ~ 100 nW, circulating field intensity at surface ~ 0.01 $\mu W/\mu m^2$). The Casimir-Polder surface potential U_s for the ground state of atomic Cs is also shown. (iii) The atomic decay rate $\gamma(d)$ as a function of distance d from the toroid's surface for TE (γ_{\parallel}) and TM (γ_{\perp}) modes. All rates in this figure are scaled to the decay rate in free space for the amplitude of the Cs $6P_{3/2} \rightarrow 6S_{1/2}$ transition, $\gamma_0/2\pi = 2.6$ MHz. The approximate distance scale probed 106

5.2Observation (a) and simulation (b-e) of atomic transits within the evanescent field of the micro-toroidal resonator for $\Delta_{ca} = \Delta_{pa} = 0$. a) Observed cavity transmission $T_B(t)$ versus time t following a triggering event at t = 0, with approximately 5×10^4 triggered transits included. The data are fit to the sum of an exponential (I) and a Gaussian (II) (green curve), with time constants $\delta t_{\rm I} = 0.78 \pm 0.02$ μ s and $\delta t_{\rm II} = 3.75 \pm 0.09 \ \mu$ s, with each component shown by the dotted lines. (b) Simulation result for 1000 triggered atoms for the cavity transmission $T_B^{(s)}(t)$ versus time t (points) from an ensemble of triggered trajectories. The green curve is a fit to the sum of an exponential and Gaussian with time constants $\delta t_{\rm I}^{(s)} = 0.69 \ \mu {\rm s}, \ \delta t_{\rm II}^{(s)} = 4.0$ μ s while the dotted lines represent the individual fit components. **c-e** Probability densities $p_i(d), p_i(g), p_i(\delta_a)$ for the distance d, coupling g, and transition frequency shift $\delta_{\rm a} = \omega_{\rm a}(d) - \omega_{\rm a}^{(0)}$ from the same simulation set as for (b). $\{d, g, \delta_{\rm a}\}$ are averaged over the first 500 ns following the trigger. For these results, the trajectories are divided into two classes based on simulated detection events for photon transition, $i = \{I, II\}$ corresponding to the two time constants $\delta t_{\rm I}^{(s)}$ (blue shaded curve) and $\delta t_{\rm II}^{(s)}$ (red shaded curve) in (b). This is a stochastic division and hence the distributions and trajectory characteristics show some overlap between sets I and II. Note: Intracavity photon number ~ 0.1 , circulating power ~ 100 nW, circulating field intensity at surface ~ 0.01 $\mu W/\mu m^2$ 109 5.3Dynamics and trajectories for strongly coupled atoms moving in surface and dipole potentials $\{U_s, U_d\}$. (a) Transmission T(t) for $\Delta_{ca}/2\pi = -40$ MHz (left) and +40 MHz (right) measured after an atom trigger at t = 0. In each panel, the circles are data for 2×10^3 trigger events; the lines are simulations of T(t) for the full model (blue), for $U_{\rm s} = 0$ (magenta), and for $U_{\rm s} = U_{\rm d} = 0$ (green). Exponential fits to the data give time constants $\delta t_{\rm red} = 0.11 \pm 0.01$ and $\delta t_{\rm blue} = 0.53 \pm 0.03 \ \mu s$, while fits to the full simulation yield time constants $\delta t_{\rm red}^{(s)} = 0.19 \pm 0.02 \ \mu s$ and $\delta t_{\rm blue}^{(s)} =$ $0.59 \pm 0.06 \ \mu$ s, where quantitative differences are attributed to simplifications inherent in the simulation model (see SI). (b) Representative atomic trajectories projected onto the $\rho - z$ plane for simulations in panel (a), with the TE mode intensity plotted on a gray scale. The upper panels are for $\Delta_{ca}/2\pi = -40$ MHz while the lower panels are for $\Delta_{ca}/2\pi = +40$ MHz. The color bars at the top of the panels match the colors of the curves in (a). For each panel, orange lines are untriggered trajectories, while triggered trajectories are represented by blue lines which turn red after a trigger at t = 0. (c) Simulations showing trajectories from a full 3D simulation with $U_{\rm s}$, $U_{\rm d}$, as well as a two-color dipole potential (FORT) triggered "on" by atom detection at t = 0. $\Delta_{\rm ca}/2\pi = +40$ MHz in correspondence to (a), (b). Blue lines represent falling atoms with the FORT beams "off" (t < 0), while red lines are trajectories after the FORT is triggered "on" and an atom begins to orbit the toroid. To illustrate the timescale, the trajectories are colored pink for $t > 50 \ \mu s$. Note: intracavity photon number ~ 0.1 , circulating power ~ 100 nW, circulating field intensity at surface ~ 0.01 μ W/ μ m². . 111

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- 5.7**Real time detection of single atom transits.** (a) Normalized transmission spectra $T(\Delta_{\rm pa})$ as a function of probe detuning $\Delta_{\rm pa}$ for g = 0 and $g/2\pi = 50$ MHz ($\theta = 0$ and $\theta = \pi/4$) at critical coupling. The spectrum for $\theta = \pi/2$ is the mirror image of the $\theta = 0$ case about the $\Delta_{pa} = 0$ axis. (b) Transmitted photon flux as a function of g for $\Delta_{\rm pa} = 0$. An atom trajectory with increasing g (say from g = 0 to $g/2\pi = 50$ MHz) results in increased $P_{\rm T}$ illustrated by the cyan arrow. (c) Experimental counts $C_1(t) + C_2(t)$ for 1501 transits from 596 atom drops with 4% false detection rate where the triggers are aligned at t = 0. (d) The same data aligned by redefining t = 0 to be the mean photon arrival time for each individual transit (blue). This alignment removes selection biasing seen in panel (a) and allows plotting of the distribution of trigger times relative to the transit center (red). Most triggers occur just prior to the peak of transmission of atom transits. The data in (c) and (d) have been smoothed for clarity, which artificially broadens the selection biasing effects in (c). In (b), (c) and (d) the maximum off-resonant transmitted photon flux is $P_{\rm T} \approx 18$ MCts/s ~ 4 pW. 122

Calculated atom-surface potential U_s^g for a Cesium atom at distance d from 5.9a SiO₂ surface with radius of curvature $R = D_m/2 = 1.5 \ \mu m$ (red) and $R \to \infty$ (blue). The limiting cases for $R \to \infty$ are shown as dotted lines. In the region where surface forces are important, the cylindrical correction provides an accurate expression for the CP potentials. For d > R, the cylindrical correction formula is no longer valid. 1275.10Experimental spectral data for various cavity detuning cases: (a) $\Delta_{ca}/2\pi =$ +40 MHz. (b) $\Delta_{ca}/2\pi = -40$ MHz. (c) $\Delta_{ca}/2\pi = +60$ MHz. In each difference spectrum, we plot the simulation for the full model (blue), $U_{\rm d} = 0$ (cyan), and $U_{\rm s} = 0$ (magenta), and $U_{\rm d} = U_{\rm s} = 0$ (green). The full simulation and $U_{\rm d} = U_{\rm s} = 0$ cases also 6.1Variations of the dipole decay rate $\gamma_{s}(d)$ for a dipole oriented parallel (||) and perpendicular (\perp) to the surface normal as a function of distance d from a semi-infinite region of SiO_2 . The decay rate is in units of the vacuum decay rate γ_0 and the wavelength of the transition is $\lambda = 852 \text{ nm.} \ldots \ldots \ldots \ldots \ldots$ 1376.2Dispersive response functions for SiO_2 and cesium atoms. (a) The dielectric function $\epsilon(i\xi)$ for SiO₂ evaluated for frequency ξ along the imaginary axis. (b) Total atomic polarizability $\alpha(i\xi)$ evaluated for frequency ξ along the imaginary axis for the $6S_{1/2}$ ground state (red) and the $6P_{3/2}$ excited state (blue) of cesium calculated as 6.3Atom-surface potentials U_s^g (red) and U_s^{ex} (blue) for a cesium atom at distance d from an SiO₂ surface. The solid lines are for a planar surface whereas the dashed lines are for a curved surface with radius of curvature $R = D_{\rm m}/2 = 1.5$ μ m. The limiting regimes for U_s^g with a planar surface are shown as dotted lines, each calculated from analytic expressions not using the Lifshitz formalism. The cylindrical surface correction weakens the potential, which is noticeable in the retarded and thermal regimes. 140Plots of $T(g_{tw}, \theta, P)$ for (a) $\Delta_{ca}/2\pi = 0$ MHz, and (b) $\Delta_{ca}/2\pi = 60$ MHz, 6.4calculated numerically from (3.13). Atoms with higher g_{tw} generally have higher T and a larger probability for detection. The variation of T with θ is evident, with a different periodicity for the two cavity detunings. 143

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An atom nearby this silicon nitride nanobeam has an initial velocity v associated with the thermal cloud velocity distribution for a given temperature $T_{\rm cloud}$, several forces acting on the atom including gravity $(F_{\rm g})$, surface force $(F_{\rm s})$, and optical forces $(F_{\rm opt})$ from the magneto-optical trapping and cooling beams. **b)** Number of atoms (N_{atoms}) that are on average in the vicinity of a single nanobeam's evanescent field region as a function of atom cloud density (ρ), located within a volume of $V = V_0 = 2 \times x0 \times y0 \times y0$ L, where x0 = 100 nm, y0 = 100 nm, and $L = 500 \ \mu m$ for the red curve (labeled (ii)). 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Chamber (e) is continuously kept at ultra-high vacuum by the ion pump (f), located at $L2 \approx 70$ cm. Atoms transported from the source chamber are collected at the science chamber by a second MOT with three orthogonal large (≈ 2 cm diameter) retroreflected beams. c) After a compression stage, the atom cloud (i) is transported (ii-iii) into the through-hole of the chip where the device is (iv) by moving the quadrupole magnetic field zero from (i) to (iv). The atoms are then collected by a third MOT, a mini-MOT, formed by three orthogonal small (≈ 1 mm diameter) retroreflected beams intersecting at the center of the chip's through-hole. Here, about $\sim 10^6$ atoms are collected (density $\sim 10^{10}$ atoms/cm³). d) Fluorescence image showing the second MOT in chamber (e). e) Absorption images showing atom cloud transport starting from the second MOT location (i), during transport (ii-iii), and the mini-MOT location inside the chip's through-hole (iv). . . . 216 8.5 Atom trapping schemes with a single nanobeam (part 1). a-b) Trapping potential U for magic-compensated scheme (similar to Sec. 7.2). a) and b) show the same curves over different plot ranges. Single nanobeam width w = 400 nm, height h =200 nm (Fig. 3.11). Red-detuned standing-wave ($\lambda_{\rm red} = 937.1$ nm, power $P_{\rm red} = 2 \times 0.4$ mW) and counter-propagating blue-detuned beams ($\lambda_{\text{blue}} = 686.1 \text{ nm}$ and 686.7 nm). dy: atom-to-surface gap along y-axis (x=0). Blue, red, magenta, green curves for P_{blue} $= 2 \times \{1, 2, 3, 10\}$ mW respectively. 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8.6 Atom trapping schemes with a single nanobeam (part 2). Normal incident incoherent counter-propagating red-detuned external illumination light trapping scheme. The single silicon nitride nanobeam has a width of w = 300 nm and height h = 200 nm. The external counter-propagating illumination beams have a wavelength $\lambda_{\rm red} = 852.55$ nm with a small \approx GHz relative detuning, and total intensity $I_{\rm red} = 2 \times 32 \ \mu W/\mu m^2$ (i.e., total power of $P_{\rm red} = 2 \times 10$ mW for a beam waist radius $\approx 10 \ \mu$ m). a) Contour plot of trap potential U as a function of the x-coordinate and atom-to-surface distance along the y-axis (x = 0), dy. b) Line cut along the y-axis (x = 0) showing the trapping potential. The red curves correspond to cesium excited state $(6P_{3/2}, F'=4)$, the blue curves correspond to the ground state $(6S_{1/2}, F=4)$, and the black-dashed curves that correspond to the ground state ($6S_{1/2}$, F=3). c-d) Contour plots of electric field amplitude |E| on the x - y cross-sectional plane, where the nanobeam's waveguide axis is along the z-axis (out of page). The white arrows indicate the external illumination beams. e-f) Line cuts of |E| of the corresponding countour plots, along the y-axis (x = 0) shown by the blue curves (i), and along the x-axis (y = 0) shown by the red 226

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- 8.10 Atom trapping schemes with a double nanobeam (part 3). Two-color RF switching scheme. Double Si nanobeam parameters: (each: width w = 300 nm, height h = 200 nm), gap along the x-axis is 200 nm. Switching between red- and blue-detuning for an x-polarized even mode forms a pseudo trap potential analogous to a quadrupole Paul trap in ion trap systems. a) Illustration of sinusoidal oscillation between maximum blue-detuned beam intensity at time $t_{\rm blue}$, (i), with a trap along x-axis and an antitrap along y-axis, transition time through $t_{\rm both}$, (ii) to (iv), to maximum red-detuned beam intensity (zero blue-detuned beam intensity) at time $t_{\rm red}$, (v), followed by the reverse cycle (arrows). In part a), the effective potential δU corresponds to the trapping potential using the parameters described below, with offset applied to the potential U, i.e., $\delta U = U(x,y) - U(x = y = 0)$, such that $\delta U(x = y = 0) = 0$ at all time. The time-dependent offset U(x = y = 0) ranges from -0.86 mK (at time t_{red} , (v) in a)) to +6.25 mK (at time t_{blue} , (i) in a)). On the x - y plane, the origin (x = y = 0) is at the symmetry center between the two double beam. \mathbf{b},\mathbf{d}) Trap potential U formed by a standing wave blue-detuned beam (wavelength $\lambda_{\rm blue} = 851$ nm, power $P_{\rm blue} = 2 \times$ 20 μ W). d) Red curve (i): line cut along x-axis (y = 0). Blue curve (ii): line cut along y-axis (x = 0). c,e) Trap potential U formed by a standing wave red-detuned beam (wavelength $\lambda_{\text{blue}} = 853 \text{ nm}$, power $P_{\text{blue}} = 2 \times 1 \mu \text{W}$). e) Red curve (i): line cut along x-axis (y = 0). Blue curve (ii): line cut along y-axis (x = 0). 238A.1 (a) x-Attocube. (b) z-Attocube. (c) Attocube piezo scan controller ANC200 for fine motion (top) and piezo step controller ANC150 for coarse motion (bottom).... 239A.2(a) Schematic of setup. (b) z-Attocube (stacked) setup. (c) x-Attocube (single) setup. Shown in (b) and (c) are: mirror (A), aluminium block (B) and block (C). 240A.3 **Calibration factor** (piezo displacement per applied voltage) distributions for a single x-Attocube for 0.5 Hz and 5 Hz sweep frequencies. 241A.4 Attocube transfer functions highlighting intrinsic mechanical resonances associated with the four configurations: (a)/(b) Single/stacked x-Attocube; (c)/(d) Single/stacked

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